SHORT NOTE

A Remark on "An Efficient Real Space Method for Orbital-Free Density-Functional Theory"

Carlos J. García-Cervera*

Mathematics Department, University of California, Santa Barbara, CA 93106, USA.

Received 14 November 2007; Accepted (in revised version) 6 December 2007

Communicated by Weinan E

Available online 31 December 2007

Abstract. In this short note we clarify some issues regarding the existence of minimizers for the Thomas-Fermi-von Weiszacker energy functional in orbital-free density functional theory, when the Wang-Teter corrections are included.

AMS subject classifications: 65M05, 74G65, 78M50

Key words: Density functional theory, Thomas-Fermi, constrained optimization.

In [1] it was claimed that there always exists a minimizer; however, the statement of Theorem 2.1 is incomplete. In this note we present the full statement, with a detailed proof.

The theorem stated in [1] holds as long as the number of electrons is below a certain critical value. The correct statement for the theorem in [1] is:

Theorem 1 (Existence of minimizers). *Given* $v \in C^{\infty}(\overline{\Omega})$, and $K_{WT} \in L^2_{loc}(\mathbb{R}^3)$, consider the problem

$$\inf_{u\in\mathcal{B}}F[u],\tag{1}$$

where F and \mathcal{B} are

$$F[u] = \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \frac{7C_{TF}N^{2/3}}{25} \int_{\Omega} u^{10/3} + \frac{4C_{TF}N^{2/3}}{5} \int_{\Omega} |u|^{5/3} \left(K_{WT} * |u|^{5/3} \right) + \frac{N}{2} \int_{\Omega} u^2 \left(\frac{1}{|\mathbf{x}|} * u^2 \right) - \frac{3}{4} \left(\frac{3N}{\pi} \right)^{1/3} \int_{\Omega} u^{8/3} + \int_{\Omega} u^2 \varepsilon (Nu^2) + \int_{\Omega} v(\mathbf{x}) u^2(\mathbf{x}) d\mathbf{x},$$
(2)

*Corresponding author. Email address: cgarcia@math.ucsb.edu (C. J. García-Cervera)

http://www.global-sci.com/

©2008 Global-Science Press

C. J. García-Cervera / Commun. Comput. Phys., 3 (2008), pp. 968-972

and

$$\mathcal{B} = \left\{ u \in H_0^1(\Omega) \middle| u \ge 0, \ \int_{\Omega} u^2 = 1 \right\}.$$
(3)

In (2), the set Ω is open and bounded, and star-shaped with respect to 0; ε is defined as

$$\varepsilon(Nu^2) = \begin{cases} \frac{\gamma}{1+\beta_1\sqrt{r_s}+\beta_2r_s}, & r_s \ge 1, \\ A\ln(r_s)+B+Cr_s\ln(r_s)+Dr_s, & r_s \le 1, \end{cases}$$
(4)

where $r_s = (4\pi Nu^2/3)^{-\frac{1}{3}}$; the parameters used are $\gamma = -0.1423$, $\beta_1 = 1.0529$, $\beta_2 = 0.3334$, A = 0.0311, B = -0.048, and $C = 2.019151940622 \times 10^{-3}$ and $D = -1.163206637891 \times 10^{-2}$ are chosen so that $\varepsilon(r)$ and $\varepsilon'(r)$ are continuous at r = 1 [6].

Then, there exists $N_0 > 0$ *such that:*

1. If $N < N_0$ then $\exists u^* \in \mathcal{B}$ such that

$$F[u^*] = \min_{u \in \mathcal{B}} F[u].$$
(5)

2. *If* $N > N_0$ *then*

$$\inf_{u\in\mathcal{B}}F[u] = -\infty.$$
(6)

Proof. The second part of the theorem was proved in [2, 3]. We outline the proof here for completeness. Since $0 \in \Omega$, $\exists \delta_0 > 0$ such that $B(0, \delta_0) \subset \Omega$. Consider a compactly supported function $u_0 \in C_0^{\infty}(B(0, 1))$, such that

$$\int_{\mathbb{R}^3} u_0^2 = 1,$$
(7)

and consider the rescaling

$$u_{\delta}(\mathbf{x}) = \frac{1}{\delta^{3/2}} u_0\left(\frac{\mathbf{x}}{\delta}\right), \quad 0 < \delta < \delta_0.$$
(8)

Then $u_{\delta} \in \mathcal{B}$, and

$$F[u_{\delta}] = \frac{1}{\delta^2} \left(\frac{1}{2} \int_{\Omega} |\nabla u_0|^2 - \frac{7C_{TF} N^{2/3}}{25} \int_{\Omega} u_0^{10/3} \right) + \mathcal{O}\left(\frac{1}{\delta}\right).$$
(9)

Define

$$A_{0} = \inf_{u \in H_{0}^{1}(\Omega), \|u\|_{2}=1} \frac{\int_{\Omega} |\nabla u|^{2}}{\int_{\Omega} u^{10/3}} > 0.$$
(10)

Then if $A_0/2 < 7C_{TF}N^{2/3}/25$, we can choose u_0 so that the leading term in (9) is negative, and when $\delta \rightarrow 0$, the desired result follows.

969