A CRITICAL VALUE FOR GLOBAL NONEXISTENCE OF SOLUTION OF A WAVE EQUATION*

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Abstract Consider the Cauchy problem for a wave equation on R^2 : $u_{tt} - \Delta u = |u|^{p-1}u$. In 1981 Glassey gave a guess to a critical value $p(2) = \frac{1}{2}(3 + \sqrt{17})$: when p > p(2) there may exist a global solution and when 1 the solution may blow up. By our main result in this paper a counter example to the guess is given that the solution may also blow up in finite time even if <math>p(2) .

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Consider the Cauchy problem for a wave equation on R^2 :

$$\begin{cases} \frac{\partial^2 u(x,t)}{\partial t^2} - \Delta u = |u|^{p-1}u, & x \in R^2, \quad 0 < t < T, \\ u(x,0) = f(x), & u_t(x,0) = g(x), & x \in R^2, \end{cases}$$
(1)

where we assume that

(H1)
$$f(x), g(x) \in C_0^{\infty}(\mathbb{R}^2), \operatorname{supp}\{f, g\} \subset \{ \|x\| \le L \},$$

$$\int_{R^2} f \mathrm{d}x > 0, \qquad \int_{R^2} g \mathrm{d}x > 0.$$

Theorem(Glassey[1]) When $1 , <math>T < +\infty$, *i.e* the solution of (1) may blow up in finite time $T < +\infty$.

In Case \mathbb{R}^3 , John[2] gave the critical value $p(3) = 1 + \sqrt{2}$: when 1 the solution may blow up and when <math>p > p(3) there may exist global solution. Thus, Glassey gave a guess that p(2) may also be a critical value for blow up: there may exist a global solution if p > p(2). But a counter example can be given to Glassey's guess by our following main result.

Theorem Let $u(x,t) \in C^2(\mathbb{R}^2 \times [0,T])$ be a nontrivial solution of (1) with finite speed of propagation. Assume that

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- The same as Glassey Theorem, and $\int_{\mathbb{R}^2} fg dx > 0$, $f(x) \neq 0$, (H1)
- 3 ,(H2)
- (H3) $I_0 = \frac{2}{p+1} \int_{\mathbb{R}^2} |f|^{p+1} dx \left[\int_{\mathbb{R}^2} \left(|\nabla f|^2 + |g|^2 \right) dx \right] \ge 0.$ Then $T < +\infty$, i.e the solution of (1) may blow up in finite time $T < +\infty$.

Remark It is well known that (H1) implies the existence of a unique classical solution to (1).

Proof We will estimate $F(t) = \int_{\mathbb{R}^2} u^2(x, t) dx$ by using the method similar to[3]. First, multiplying the equation (1) by u(x,t) and integrating over R^2 , we have

$$\frac{1}{2}F''(t) = \frac{p-1}{p+1}\int_{R^2} |u|^{p+1} \mathrm{d}x + \frac{2}{p+1}\int_{R^2} |u|^{p+1} \mathrm{d}x + \int_{R^2} |u_t|^2 \mathrm{d}x - \int_{R^2} |\nabla u|^2 \mathrm{d}x.$$
 (2)

Next, multiplying the equation (1) by u_t and integrating over $R^2 \times [0, t]$ we have

$$\int_{R^2} |u_t|^2 \mathrm{d}x = \frac{2}{p+1} \int_{R^2} |u_t|^{p+1} \mathrm{d}x - \int_{R^2} |\nabla u|^2 \mathrm{d}x - I_0.$$
(3)

By (H3), (2) and (3) yield

$$\frac{1}{2}F''(t) = \frac{p-1}{p+1}\int_{R^2} |u|^{p+1} \mathrm{d}x + 2\int_{R^2} |u_t|^2 \mathrm{d}x + I_0.$$
(4)

Thus $F''(t) \ge 0$ and F'(t) is monotone nondecreasing. Therefore $F'(t) \ge F'(0) > 0$ by (H1), and F(t) is also monotone nondecreasing , thus $F(t) \ge F(0) = \int_{B^2} f^2 dx > 0$.

Now, by finite speed of propagation and by (H1), we have

$$F(t) = \int_{R^2} u^2 \mathrm{d}x = \int_{\|x\| \le t+L} u^2 \mathrm{d}x \le \left\{ \int_{\|x\| \le t+L} |u|^{p+1} \mathrm{d}x \right\}^{\frac{2}{p+1}} \left\{ \int_{\|x\| \le t+L} 1 \mathrm{d}x \right\}^{\frac{p-1}{p+1}}$$

i.e.

$$F(t)^{\frac{p+1}{2}} \le \pi^{\frac{p-1}{2}} (t+L)^{p-1} \int_{\mathbb{R}^2} |u|^{p+1} \mathrm{d}x.$$
 (5)

Combining (5) with (4), we obtain

$$F''(t) \ge C_0(t+L)^{1-p}F(t)^{\frac{p+1}{2}}$$
(6)

but $F(t) \ge F(0) = \int_{B^2} f^2 dx > 0$, thus

$$F''(t) \ge k_0(t+L)^{1-p},$$

 \mathbf{SO}

$$F'(t) \ge F'(0) + \frac{k_0}{2-p}(t+L)^{2-p} - \frac{k_0}{2-p}L^{2-p}.$$