

SUPERCONVERGENCE OF FEM FOR SINGULAR SOLUTION*

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Superconvergence of the finite element method (FEM) has been discussed extensively for the problem having smooth solution (See Krizek and Neittaanmaki [8]). A typical result in this direction is the following (see Lin and Xie [4] for details). Consider the model problem

$$-\Delta u = f \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega,$$

where $\Omega \subset \mathbf{R}^2$ is a bounded domain with a smooth boundary $\partial\Omega$ and f is a smooth function. In order to keep the mesh varying regularly we impose on Ω a kind of "piecewise almost uniform triangulation" which can be constructed piecewisely by the vertices of a smoothly transformed uniform mesh. For any node z in the interior of each piece there exist two triangles e and e' such that $e \cap e' = \{z\}$. Then, the average gradient

$$\bar{\nabla} u^h(z) = \frac{1}{2}(\nabla u^h|_e + \nabla u^h|_{e'})$$

has not only the usual type of superconvergence

$$(\bar{\nabla} u^h - \nabla u)(z) = O(h^2)$$

but also an extrapolation type of superconvergence

$$\frac{1}{3}\bar{\nabla}(4u^{h/2} - u^h)(z) - \nabla u(z) = O(h^4 \log \frac{1}{h}).$$

We are concerned in this paper with the superconvergence for the singular solution due to re-entrant corners or changing the boundary conditions.

For simplicity we suppose that Ω is composed of rectangles and the boundary $\partial\Omega$ is parallel to the x - and y -axis and has only one re-entrant corner at the origin 0. Let α be the interior angle at 0 and $\beta = \pi/\alpha$.

It is easy to see that

$$u \in H_{(\tau+1)}^3 \quad \text{for } \tau > 1 - \beta,$$

where the Sobolev space $H_{(\tau+1)}^3$ is defined using the weighted norm

$$\|u\|_{3,(\tau+1)} = \left[\sum_{|j| \leq 3} \int_{\Omega} (|X|^{\tau-2+|j|} |\partial^j u|)^2 dX \right]^{1/2}$$

with $X = (x, y)$.

We now introduce a rectangular mesh $T^h = \{e\}$, where (x_e, y_e) denotes the center of the element e and $2h_e$ and $2k_e$ are its widths in the x - and y -direction, respectively. Further, we set

$$d_e = \max(h_e, k_e), \quad h = \max\{d_e, e \in T^h\},$$

$$d_0 = \max\{d_e, e \in T^h, 0 \in e\}, \quad r_e = \min\{|X|, X \in e\}.$$

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Let T^h be split into two parts,

$$\Omega_0 = \{e \in T^h, r_e < d_0\}, \quad \Omega_1 = \{e \in T^h, d_0 \leq r_e\},$$

where the local meshes are assumed to satisfy the grading conditions

$$d_0 \leq ch^q, \quad q > \frac{t}{\beta}, \quad t \leq 2;$$

$$c_1 hr_e^p \leq d_e \leq chr_e^p, \quad \forall d_e \leq r_e, \quad p = 1 - \frac{1}{q},$$

where q is the grading parameter and t is the superconvergence parameter. For example, if $\Omega = (-1, 1) \times (-1, 1) \setminus [0, 1] \times \{0\}$ (a slit domain), such meshes can be constructed by taking nodes

$$(\pm(i/n)^q, \pm(j/n)^q) (1 \leq i, j \leq n), \quad q > 2t.$$

Since a larger t will lead to a larger q , the user has to make up his choice between a higher accuracy and a less graded mesh. We note that the total number of nodes of the graded meshes is the same as for a uniform mesh of size h , and that the size of the largest element is of the order h .

Let

$$\Omega_2 = \{X \in \Omega, |X| \geq \rho > 0\},$$

z be the interior node of Ω_2 and N the number of all interior nodes of Ω_2 :

$$N = O(h^{-2}).$$

For such z there exist two elements e and e' such that $e \cap e' = \{z\}$ and we can define, for $v \in S^h$ the piecewise bilinear finite element space, the average gradient

$$\bar{\partial}_x v(z) = \frac{h_e}{h_e + h_{e'}} \partial_x v|_{e'}(z) + \frac{h_{e'}}{h_e + h_{e'}} \partial_x v|_e(z),$$

$$\bar{\partial}_y v(z) = \frac{k_e}{k_e + k_{e'}} \partial_y v|_{e'}(z) + \frac{k_{e'}}{k_e + k_{e'}} \partial_y v|_e(z).$$

Let $u^I \in S^h$ be the interpolation of u and $u^R \in S^h$ the Ritz projection of u . It is easy to see from Taylor expansion the superconvergence of u^I after averaging:

Lemma 1.

$$|(\bar{\partial} u^I - \partial u)(z)| \leq ch^t \|u\|_{3, \infty, \Omega_2},$$

where the notation $\bar{\partial}$ means $\bar{\partial}_x$ or $\bar{\partial}_y$.

Our purpose is to prove the superconvergence of u^R after averaging:

Theorem. *The grading parameter q increases the gradient accuracy from β -order to nearly $q\beta$ -order:*

$$\left[\frac{1}{N} \sum_{z \in \Omega_2} |(\bar{\partial} u^R - \partial u)(z)|^2 \right]^{1/2} \leq ch^t, \quad t < q\beta.$$

The proof of our theorem is based on the lemmas as follows (c.f. [1]-[2]).

Lemma 2. *For the function $F(x)$ satisfying $F(x_e \pm h_e) = 0$, we have*

$$\int_{x_e - h_e}^{x_e + h_e} F dx = \frac{1}{2} \int_{x_e - h_e}^{x_e + h_e} P F'' dx,$$

where $P(x) = (x - x_e + h_e)(x - x_e - h_e)$.

Proof. Note that

$$P(x_e \pm h_e) = 0, \quad P''(x) = 2.$$

Integrating by parts leads to Lemma 2:

$$\int_{x_e-h_e}^{x_e+h_e} PF'' dx = - \int_{x_e-h_e}^{x_e+h_e} P' F' dx = \int_{x_e-h_e}^{x_e+h_e} 2F dx.$$

Lemma 3. For $u \in H^3(e)$ and $\nu \in S^h$, there holds

$$a_e(u^I - u, \nu) = -\frac{1}{2} \int_e (P \partial_y \partial_x^2 u \partial_y \nu + Q \partial_x \partial_y^2 u \partial_x \nu) + \int_e (P + Q) \partial_y \partial_x (u^I - u) \partial_y \partial_x \nu,$$

where $P(x)$ is defined as in Lemma 2 and $Q(y) = (y - y_e + k_e)(y - y_e - k_e)$.

Proof. Since

$$a_e(u^I - u, \nu) = \int_e \partial_x (u^I - u) \partial_x \nu + \int_e \partial_y (u^I - u) \partial_y \nu$$

we need only to expand, say, the second integral. Let l_1 and l_2 stand for the two edges of element e parallel to the x -direction. Integrating by parts we obtain

$$\int_e \partial_y (u^I - u) \partial_y \nu = \left[\int_{l_1} - \int_{l_2} \right] (u^I - u) \partial_y \nu dx.$$

Setting $F = (u^I - u) \partial_y \nu$ in Lemma 2 we have

$$\partial_x^2 F = -\partial_x^2 u \partial_y \nu + 2 \partial_x (u^I - u) \partial_y \partial_x \nu$$

and hence, by Lemma 2,

$$\int_{l_i} (u^I - u) \partial_y \nu dx = -\frac{1}{2} \int_{l_i} P(x) \partial_x^2 u \partial_y \nu dx + \int_{l_i} P(x) \partial_x (u^I - u) \partial_y \partial_x \nu.$$

Thus

$$\int_e \partial_y (u^I - u) \partial_y \nu = -\frac{1}{2} \int_e P(x) \partial_y (\partial_x^2 u \partial_y \nu) + \int_e P(x) \partial_y (\partial_x (u^I - u) \partial_y \partial_x \nu)$$

and Lemma 3 follows.

Since u becomes singular in Ω_0 we need the following

Lemma 4. For $\tau = 1 - t/q$,

$$|a_{\Omega_0}(u^I - u, \nu)| \leq ch^t \|u\|_{2,(\tau)} \|\nu\|_1.$$

Proof. Since $u \in W^{1,a} \cap W^{2,b}$ for

$$2 < a = \frac{2b}{2-b} < \frac{2}{\tau}, \quad b < \frac{2}{1+\tau}$$

we have, by the Holder inequality,

$$|u - u^I|_{1,\Omega_0} \leq cd_0^{1-2/a} |u - u^I|_{1,a,\Omega_0} \leq cd_0^{1-2/a} |u|_{1,a,\Omega_0}.$$

An imbedding theorem implies that

$$|u|_{1,a,\Omega_0} \leq c \|u\|_{2,b,\Omega_0} + cd_0^{-1} |u|_{1,b,\Omega_0}$$

and the Holder inequality implies that

$$\|u\|_{2,b,\Omega_0} \leq cd_0^{2/b-1-\tau} \|u\|_{2,(\tau)}, \quad |u|_{1,b,\Omega_0} \leq cd_0^{2/b-1-(\tau-1)} |u|_{1,(\tau-1)}$$

and hence,

$$|u - u^I|_{1,\Omega_0} \leq cd_0^{1-\tau} \|u\|_{2,(\tau)} \leq ch^t \|u\|_{2,(\tau)}.$$

Lemma 5. $\|u^I - u^R\|_1 \leq ch^t \|u\|_{3,(\tau+1)}$.

Proof. For $\nu \in S^h$

$$a(u^I - u^R, \nu) = a(u^I - u, \nu) = \sum_{e \in \Omega_1} a_e(u^I - u, \nu) + a_{\Omega_0}(u^I - u, \nu).$$

For $e \in \Omega_1$ we use Lemma 3 and the following estimates:

$$|\partial_x \partial_y (u^I - u)|_{0,e} \leq cd_e |u|_{3,e}, \quad |\partial_x \partial_y \nu|_{0,e} \leq cd_e^{-1} \|\nu\|_{1,e}.$$

Since

$$d_e^2 \leq cd_e^t r_e^{2-t} \leq ch^t r_e^{1+\tau},$$

we obtain

$$|a_e(u^I - u, \nu)| \leq ch^t \|u\|_{3,(\tau+1),e} \|\nu\|_{1,e}.$$

Thus, combining with Lemma 4 we obtain Lemma 5.

We now prove our theorem as follows.

Using an inverse inequality and noting $r_e \geq \rho$ we have

$$|\bar{\partial}(u^I - u^R)(z)| \leq cd_e^{-1} \|u^I - u^R\|_{1,e} \leq ch^{-1} r_e^{-\nu} \|u^I - u^R\|_{1,e} \leq ch^{-1} \rho^{-\nu} \|u^I - u^R\|_{1,e}.$$

Thus,

$$\left[\frac{1}{N} \sum |\bar{\partial}(u^I - u^R)(z)|^2 \right]^{1/2} \leq c \|u^I - u^R\|_1.$$

Combining with Lemma 1 we obtain our theorem.

A $2t$ -order accuracy can be achieved for the finite element solution after extrapolation (see [2], [5]-[6]). A similar result holds true also for the finite element gradient (see [3]).

References

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