

## $L^\infty$ CONVERGENCE OF TRUNC ELEMENT FOR THE BIHARMONIC EQUATION<sup>\*1)</sup>

M. Wang

(Department of Mathematics, Beijing University, Beijing, China)

### Abstract

The paper considers the  $L^\infty$  convergence for TRUNC finite elements solving the boundary value problems of the biharmonic equation. The nearly optimal  $L^\infty$  estimates for the error of first order derivatives are given.

The TRUNC plate element is proposed and developed by Argyris et al.. The numerical experiences show that the element has very good results<sup>[1,2]</sup>. The mathematical proof of convergence of the TRUNC element is given by Shi Zong-ci in paper [3]. This paper will consider the  $L^\infty$  convergence for the TRUNC plate element.

### 1. The TRUNC Plate Element

Given a triangle  $T$  with vertices  $a_i = (x_i, y_i), i = 1, 2, 3$ , we denote by  $\lambda_i$  the area coordinates for the triangle  $T$  and put

$$\begin{aligned} \xi_1 &= x_2 - x_3, & \xi_2 &= x_3 - x_1, & \xi_3 &= x_1 - x_2, \\ \eta_1 &= y_2 - y_3, & \eta_2 &= y_3 - y_1, & \eta_3 &= y_1 - y_2. \end{aligned}$$

The nodal parameters of the element are the function values and the values of the two first derivatives at the vertices of the triangle  $T$ . According to paper [3], on the triangle  $T$  the shape function is an incomplete cubic polynomial,

$$\begin{aligned} w &= b_1\lambda_1 + b_2\lambda_2 + b_3\lambda_3 + b_4\lambda_1\lambda_2 + b_5\lambda_2\lambda_3 + b_6\lambda_3\lambda_1 \\ &+ b_7(\lambda_1^2\lambda_2 - \lambda_1\lambda_2^2) + b_8(\lambda_2^2\lambda_3 - \lambda_2\lambda_3^2) + b_9(\lambda_3^2\lambda_1 - \lambda_3\lambda_1^2), \end{aligned} \quad (1.1)$$

which is uniquely determined by the nine nodal parameters  $w_i, w_x(i), w_y(i), i = 1, 2, 3$ . The coefficients  $b_i$  are determined as follows,

$$\left\{ \begin{aligned} b_i &= w_i, & i &= 1, 2, 3, \\ b_4 &= -\frac{1}{2}\{(w_x(1) - w_x(2))\xi_3 + (w_y(1) - w_y(2))\eta_3\}, \\ b_5 &= -\frac{1}{2}\{(w_x(2) - w_x(3))\xi_1 + (w_y(2) - w_y(3))\eta_1\}, \\ b_6 &= -\frac{1}{2}\{(w_x(3) - w_x(1))\xi_2 + (w_y(3) - w_y(1))\eta_2\}, \\ b_7 &= w_1 - w_2 - \frac{1}{2}(w_x(1) + w_x(2))\xi_3 - \frac{1}{2}(w_y(1) + w_y(2))\eta_3, \\ b_8 &= w_2 - w_3 - \frac{1}{2}(w_x(2) + w_x(3))\xi_1 - \frac{1}{2}(w_y(2) + w_y(3))\eta_1, \\ b_9 &= w_3 - w_1 - \frac{1}{2}(w_x(3) + w_x(1))\xi_2 - \frac{1}{2}(w_y(3) + w_y(1))\eta_2. \end{aligned} \right. \quad (1.2)$$

---

\* Received December 24, 1993.

<sup>1)</sup> The project was supported by the National Natural Sciences Foundation of China.

The shape form (1.1) with (1.2) is another one of Zienkiewicz's element. This element is a  $C^0$  element, nonconforming for plate bending problems, which converges to the true solution only for very special meshes. The TRUNC element is obtained by modifying the variational formulation.

Let  $\Omega$  be a convex polygon in  $R^2$ ,  $f \in L^2(\Omega)$ . Consider the plate bending problem with the clamped boundary conditions,

$$\begin{cases} \Delta^2 u = f, & \text{in } \Omega \\ u|_{\partial\Omega} = \frac{\partial u}{\partial N}|_{\partial\Omega} = 0 \end{cases} \quad (1.3)$$

The weak form of the problem (1.3) is to find  $u \in H_0^2(\Omega)$  such that,

$$a(u, v) = (f, v), \quad \forall v \in H_0^2(\Omega), \quad (1.4)$$

where

$$\begin{aligned} a(u, v) &= \int_{\Omega} (\Delta u \Delta v + (1 - \sigma)(2u_{xy}v_{xy} - u_{xx}v_{yy} - u_{yy}v_{xx})) dx dy, \\ (f, v) &= \int_{\Omega} f v dx dy, \end{aligned} \quad (1.5)$$

and  $0 < \sigma < \frac{1}{2}$  is the Poisson ratio.

Dividing  $\Omega$  into a regular family  $\mathbb{T}_h$  of triangular elements  $T$  with diameters  $h_T \leq h$ , and defining on each triangle  $T$  the shape function in the form (1.1) and (1.2), we obtain the finite element space  $V_h$ . Then, the standard finite element approximation of problem (1.4) is to find  $u_h \in V_h$  such that

$$a_h(u_h, v_h) = (f, v_h), \quad \forall v_h \in V_h, \quad (1.5)$$

where

$$a_h(u, v) = \sum_T \int_T (\Delta u \Delta v + (1 - \sigma)(2u_{xy}v_{xy} - u_{xx}v_{yy} - u_{yy}v_{xx})) dx dy. \quad (1.6)$$

The modification of variational formulation (1.6) is carried out as follows. Every function  $v_h \in V_h$  can be split into two parts,

$$v_h = \bar{v}_h + v'_h, \quad (1.7)$$

where

$$\bar{v}_h|_T = a_1\lambda_1 + a_2\lambda_2 + a_3\lambda_3 + a_4\lambda_1\lambda_2 + a_5\lambda_2\lambda_3 + a_6\lambda_3\lambda_1, \quad (1.8)$$

representing a full quadratic polynomial on  $T$ , and

$$v'_h = a_7(\lambda_1^2\lambda_2 - \lambda_1\lambda_2^2) + a_8(\lambda_2^2\lambda_3 - \lambda_2\lambda_3^2) + a_9(\lambda_3^2\lambda_1 - \lambda_3\lambda_1^2), \quad (1.9)$$

being a cubic polynomial. Define a new discrete bilinear form,

$$b_h(v_h, w_h) = a_h(\bar{v}_h, \bar{w}_h) + a_h(v'_h, w'_h), \quad \forall v_h, w_h \in V_h. \quad (1.10)$$

Obviously,

$$b_h(v_h, w_h) = a_h(v_h, w_h) - a_h(\bar{v}_h, w'_h) - a_h(v'_h, \bar{w}_h). \quad (1.11)$$

The solution, produced by TRUNC element, is exactly the one of a variational problem to find  $u_h \in V_h$  such that

$$b_h(u_h, v_h) = (f, v_h), \quad \forall v_h \in V_h. \quad (1.12)$$

## 2. Some Estimates

Before going to discuss the  $L^\infty$  convergence of the TRUNC element, we need some estimates about it. For function  $v_h \in L^2(\Omega)$  and  $v_h|_T \in H^m(T), \forall T \in \mathbb{T}_h$ , define a seminorm as follows,

$$|v_h|_{m,h} = \left( \sum_{T \in \mathbb{T}_h} |v_h|_{m,T}^2 \right)^{1/2}. \quad (2.1)$$

Throughout the paper,  $C$  is a generic constant independent of  $h$ .

From paper [3], we have the following lemmas.

**Lemma 1.** *Let  $u \in H^3(\Omega) \cap H_0^2(\Omega)$  be the solution of problem (1.4) and  $u_h \in V_h$  be the solution of problem (1.12). Then*

$$\|u - u_h\|_{1,\Omega} + h^{-1}|u - u_h|_{2,h} \leq Ch^2|u|_{3,\Omega}. \quad (2.2)$$

**Lemma 2.** *According to representations (1.7) to (1.9), the shape function is  $w_h = \bar{w}_h + w'_h$ . Then*

$$|\bar{w}_h|_{2,T} + |w'_h|_{2,T} \leq C|w_h|_{2,T}, \quad (2.3)$$

for each  $T \in \mathbb{T}_h$ .

For  $T \in \mathbb{T}_h$ , we define the interpolation operator  $\Pi_T$  such that for  $\forall v \in H^3(T)$ ,  $\Pi_T v$  is in the form of shape function (1.1) and  $\Pi_T v, (\Pi_T v)_x$  and  $(\Pi_T v)_y$  are equal to  $v, v_x$  and  $v_y$  at the vertices of  $T$  respectively. For  $w \in L^2(\Omega)$  and  $w|_T \in H^3(T), \forall T \in \mathbb{T}_h$ , we define  $\Pi_h w \in L^2(\Omega)$  such that  $\Pi_h w|_T = \Pi_T w$ .

Let  $(x_0, y_0) \in \bar{\Omega}$  be a fixed point, define the weight function  $\rho$  as follows,

$$\rho(x, y) = (x - x_0)^2 + (y - y_0)^2 + h^2.$$

For integer  $\beta$  and  $v \in H^m(T)$  and  $T \in \mathbb{T}_h$ , define

$$|v|_{m,(\beta),T} = \left( \sum_{i+j=m} \int_T \rho^{-\beta} \left| \frac{\partial^m v}{\partial x^i \partial y^j} \right|^2 dx dy \right)^{1/2}. \quad (2.4)$$

When  $v \in L^2(\Omega)$  and  $v|_T \in H^m(T)$  for  $\forall T \in \mathbb{T}_h$ , define

$$|v|_{m,(\beta)} = \left( \sum_{T \in \mathbb{T}_h} |v|_{m,(\beta),T}^2 \right)^{1/2}. \quad (2.5)$$

In the estimates involving the weight function  $\rho$ , the constant  $C$  is also independent of  $(x_0, y_0)$ . For the weight function, the following inequalities,

$$|v|_{m,(\gamma)} \leq h^{-(\gamma-\beta)} |v|_{m,(\beta)}, \quad \gamma > \beta, \quad (2.6)$$

are true for  $v \in L^2(\Omega)$  and  $v|_T \in H^m(T), \forall T \in \mathbb{T}_h$ . And

$$|v|_{0,(1)} \leq C |\ln h|^{1/2} \|v\|_{0,\infty,\Omega}, \quad \forall v \in L^\infty(\Omega), \quad (2.7)$$

$$|(v, w)| \leq |v|_{0,(\beta)} |v|_{0,(-\beta)}, \quad v, w \in L^2(\Omega). \quad (2.8)$$

For the interpolation operator  $\Pi_T$ ,

$$|v - \Pi_T v|_{k,(\beta),T} + |v - \overline{\Pi_T v}|_{k,(\beta),T} + |(\Pi_T v)'|_{k,(\beta),T} \leq Ch^{3-k} |v|_{3,(\beta),T}, \quad (2.9)$$

hold for  $0 \leq k \leq 3, v \in H^3(T), T \in \mathbb{T}_h$ .

For  $v, w \in L^2(\Omega)$  and  $v|_T, w|_T \in H^3(T)$ , we define

$$E(v, w) = \sum_{T \in \mathbb{T}_h} \int_{\partial T} [\Delta v - (1 - \sigma) \frac{\partial^2 v}{\partial^2 s}] \frac{\partial w}{\partial n} ds. \quad (2.10)$$

**Lemma 3.** For each integer  $\beta$ , the inequalities

$$|(\Delta^2 v, w_h) - b_h(\Pi_h v, w_h)| \leq Ch |v|_{3,(\beta)} |w_h|_{2,(-\beta)}, \quad (2.11)$$

$$|E(v, \bar{w}_h)| \leq Ch |v|_{3,(\beta)} |w_h|_{2,(-\beta)}, \quad (2.12)$$

$$|E(v - \bar{v}_h, w'_h)| \leq Ch (|v - \bar{v}_h|_{2,(\beta)} + h |v|_{3,(\beta)}) |w_h|_{3,(-\beta)}, \quad (2.13)$$

$$|E(v, \overline{\Pi_h w} - w)| \leq Ch^2 |v|_{3,(\beta)} |w|_{3,(-\beta)}, \quad (2.14)$$

are true for  $\forall v, w \in H_0^2(\Omega) \cap H^3(\Omega)$  and  $\forall v_h, w_h \in V_h$ .

*Proof.* By the way used in [3] directly or alternatively, we get

$$|(\Delta^2 v, w_h) - b_h(\Pi_h v, w_h)| \leq Ch \sum_{T \in \mathbb{T}_h} |v|_{3,T} |w_h|_{2,T}$$

$$|E(v, \bar{w}_h)| \leq Ch \sum_{T \in \mathbb{T}_h} |v|_{3,T} |w_h|_{2,T},$$

$$|E(v - \bar{v}_h, w'_h)| \leq Ch \sum_{T \in \mathbb{T}_h} (|v - \bar{v}_h|_{2,T} + h |v|_{3,T}) |w_h|_{3,T},$$

$$|E(v, \overline{\Pi_h w} - w)| \leq Ch^2 \sum_{T \in \mathbb{T}_h} |v|_{3,T} |w|_{3,T}.$$

On the other hand, for  $\forall T \in \mathbb{T}_h$ ,

$$\max_{(x,y) \in T} \rho(x,y) \leq C \min_{(x,y) \in T} \rho(x,y). \quad (2.15)$$

Then the lemma follows.

### 3. The $L^\infty$ Estimates

In this section, we will prove the following theorem.

**Theorem 1.** *Let  $u$  be the solution of problem (1.4) and  $u_h$  the solution of problem (1.12). Then*

$$|u - u_h|_{1,\infty,\Omega} \leq Ch^2 |\ln h|^{5/4} |u|_{3,\infty,\Omega}, \quad (3.1)$$

when  $u \in W^{3,\infty}(\Omega)$ .

By the interpolation result<sup>[6]</sup>, we have

$$\begin{aligned} |u - u_h|_{1,\infty,\Omega} &\leq |u - \Pi_h u|_{1,\infty,\Omega} + |\Pi_h u - u_h|_{1,\infty,\Omega} \\ &\leq Ch^2 |u|_{3,\infty,\Omega} + |\Pi_h u - u_h|_{1,\infty,\Omega}. \end{aligned} \quad (3.2)$$

So we must estimate  $|\Pi_h u - u_h|_{1,\infty,\Omega}$ . Let  $T' \in \mathbb{T}_h$  be the element such that  $|\Pi_h u - u_h|_{1,\infty,\Omega} = |\Pi_h u - u_h|_{1,\infty,T'}$ . Without loss the generality, suppose that

$$|\Pi_h u - u_h|_{1,\infty,T'} = \left| \frac{\partial(\Pi_h u - u_h)}{\partial x} \right|_{0,\infty,T'}.$$

Let  $(x_0, y_0) \in T'$  be the point such that

$$\left| \frac{\partial(\Pi_h u - u_h)}{\partial x} \right|_{0,\infty,T'} = \left| \frac{\partial(\Pi_h u - u_h)}{\partial x} (x_0, y_0) \right|.$$

To prove (3.1), we need some results about the regular Green function. Let  $q \in P_3(T')$  satisfy

$$\int_{T'} qp \, dx dy = \frac{\partial}{\partial x} p(x_0, y_0), \quad \forall p \in P_3(T'). \quad (3.3)$$

Define  $\delta_h \in L^2(\Omega)$  such that,

$$\delta_h(x, y) = \begin{cases} q(x, y) & (x, y) \in T' \\ 0 & \text{otherwise} \end{cases} \quad (3.4)$$

From Lemma 3 in [5], we have

$$h \|\delta_h\|_{0,\Omega} + \|\delta_h\|_{-1,\Omega} + \|\delta_h\|_{0,(-1)} \leq Ch^{-1}. \quad (3.5)$$

Let  $g$  be the regular Green function determined by

$$\begin{cases} \Delta^2 g = \delta_h, & \text{in } \Omega \\ g|_{\partial\Omega} = \frac{\partial g}{\partial N}|_{\partial\Omega} = 0 \end{cases} \quad (3.6)$$

and  $g_h$  be its finite element solution by the TRUNC element, i.e.,

$$b_h(g_h, v_h) = (\delta_h, v_h). \quad \forall v_h \in V_h. \quad (3.7)$$

From (2.2) and (3.5), we get

$$\|g - g_h\|_{1,\Omega} + h|g - g_h|_{2,h} + h^2|g|_{3,\Omega} \leq Ch. \quad (3.8)$$

By the way used in [5], we have

$$\|g\|_{2,\Omega} + |g|_{3,(-1)} \leq C|\ln h|^{1/2}. \quad (3.9)$$

**Lemma 4.**

$$|g - \Pi_h g|_{2,(-1)} \leq Ch|\ln h|^{1/2} \quad (3.10)$$

$$|\Pi_h g - g_h|_{2,(-1)} \leq Ch|\ln h|^{3/4}. \quad (3.11)$$

*Proof.* Inequality (??) follows from (??) and (2.9). Now we prove (??). For  $v, w \in H^2(T)$  and integer  $\beta$ , we define

$$a_T(v, w) = \int_T (\Delta v \Delta w + (1 - \sigma)(2v_{xy}w_{xy} - v_{xx}w_{yy} - v_{yy}w_{xx})) dx dy$$

$$a_{T,\beta}(v, w) = \int_T \rho^{-\beta} (\Delta v \Delta w + (1 - \sigma)(2v_{xy}w_{xy} - v_{xx}w_{yy} - v_{yy}w_{xx})) dx dy.$$

By Leibniz' formula, (2.8) and the inequality

$$\left| \frac{\partial \rho}{\partial x} \right|^2 + \left| \frac{\partial \rho}{\partial y} \right|^2 \leq 4\rho, \quad (3.12)$$

we obtain

$$a_{T,-1}(v, v) \leq a_T(v, \rho v) + C|v|_{2,T,(-1)}|v|_{1,T} - 2(1 + \sigma) \int_T \Delta v v dx dy \quad (3.13)$$

It is easy to prove that for  $v_h \in V_h$  the following inequality is true,

$$|v_h|_{2,T,(\beta)}^2 \leq C[a_{T,\beta}(\bar{v}_h, \bar{v}_h) + a_{T,\beta}(v'_h, v'_h)], \quad \forall T \in \mathcal{T}_h. \quad (3.14)$$

From (??), (??), (2.3) and (2.15), we derive that

$$\begin{aligned} |v_h|_{2,(-1)}^2 &\leq C \sum_{T \in \mathcal{T}_h} [a_{T,-1}(\bar{v}_h, \bar{v}_h) + a_{T,-1}(v'_h, v'_h)] \\ &\leq C \left\{ a_h(\bar{v}_h, \rho \bar{v}_h) + a_h(v'_h, \rho v'_h) + |\bar{v}_h|_{2,(-1)}|\bar{v}_h|_{1,\Omega} + |v'_h|_{2,(-1)}|v'_h|_{1,\Omega} \right. \\ &\quad \left. + \left| \sum_{T \in \mathcal{T}_h} \int_T \Delta \bar{v}_h \bar{v}_h dx dy + \sum_{T \in \mathcal{T}_h} \int_T \Delta v'_h v'_h dx dy \right| \right\} \\ &\leq C \left\{ a_h(\bar{v}_h, \rho \bar{v}_h) + a_h(v'_h, \rho v'_h) + |v_h|_{2,(-1)}|v_h|_{1,\Omega} \right. \\ &\quad \left. + \left| \sum_{T \in \mathcal{T}_h} \int_T \Delta \bar{v}_h v_h dx dy \right| + \left| \sum_{T \in \mathcal{T}_h} \int_T (\Delta v'_h - \Delta \bar{v}_h) v'_h dx dy \right| \right\}. \end{aligned} \quad (3.15)$$

Then we get, from (2.3) and (2.9),

$$\begin{aligned} |v_h|_{2,(-1)}^2 &\leq C \left\{ a_h(\bar{v}_h, \rho\bar{v}_h) + a_h(v'_h, \rho v'_h) + |v_h|_{2,(-1)} |v_h|_{1,\Omega} \right. \\ &\quad \left. + h^2 |v_h|_{2,h}^2 + \left| \sum_{T \in \mathbb{T}_h} \int_T \Delta \bar{v}_h v_h dx dy \right| \right\}. \end{aligned} \quad (3.16)$$

By similar way used in the proof of lemma 3.5 in [3] and the fact that  $v_h \in H_0^1(\Omega)$ , we have

$$\left| \sum_{T \in \mathbb{T}_h} \int_T \Delta \bar{v}_h v_h dx dy \right| \leq Ch |v_h|_{1,\Omega} |v_h|_{2,h}, \quad (3.17)$$

and then

$$\begin{aligned} |v_h|_{2,(-1)}^2 &\leq C \left\{ a_h(\bar{v}_h, \rho\bar{v}_h) + a_h(v'_h, \rho v'_h) + |v_h|_{2,(-1)} |v_h|_{1,\Omega} \right. \\ &\quad \left. + h^2 |v_h|_{2,h}^2 + h |v_h|_{1,\Omega} |v_h|_{2,h} \right\}. \end{aligned} \quad (3.18)$$

For  $v_h \in V_h$ ,  $\Pi_h(\rho v_h) \in V_h$ . Then from (2.3), (2.8) and (2.15), we get

$$\begin{aligned} |a_h(\bar{v}_h, \rho\bar{v}_h) + a_h(v'_h, \rho v'_h)| &= |b_h(v_h, \Pi_h(\rho v_h)) + a_h(\bar{v}_h, \rho\bar{v}_h - \overline{\Pi_h(\rho v_h)}) \\ &\quad + a_h(v'_h, \rho v'_h - (\Pi_h(\rho v_h))')| \\ &\leq |b_h(v_h, \Pi_h(\rho v_h))| + |\bar{v}_h|_{2,(-1)} |\rho\bar{v}_h - \overline{\Pi_h(\rho v_h)}|_{2,(1)} \\ &\quad + |v'_h|_{2,(-1)} |\rho v'_h - (\Pi_h(\rho v_h))'|_{2,(1)} \\ &\leq |b_h(v_h, \Pi_h(\rho v_h))| + |v_h|_{2,(-1)} (|\rho\bar{v}_h - \overline{\Pi_h(\rho v_h)}|_{2,(1)} \\ &\quad + |\rho v'_h - (\Pi_h(\rho v_h))'|_{2,(1)}). \end{aligned}$$

Set  $P_T^0 : L^2(T) \rightarrow P_0(T)$  be the  $L^2(T)$  orthogonal projection operator and  $P_h : L^2(\Omega) \rightarrow L^2(\Omega)$  defined such that for  $v \in L^2(\Omega)$ ,  $P_h v|_T = P_T^0(v|_T)$ ,  $\forall T \in \mathbb{T}_h$ . Then from (2.?) and the inverse inequality of polynomials,

$$\begin{aligned} |\rho\bar{v}_h - \overline{\Pi_h(\rho v_h)}|_{2,(1)} &= |(\rho - P_h \rho)\bar{v}_h - \overline{\Pi_h((\rho - P_h \rho)v_h)}|_{2,(1)} \\ &\leq |(\rho - P_h \rho)(\bar{v}_h - v_h)|_{2,(1)} + |(\rho - P_h \rho)v_h - \overline{\Pi_h((\rho - P_h \rho)v_h)}|_{2,(1)} \\ &\leq C(h|v_h|_{2,h} + |v_h|_{1,\Omega}), \end{aligned}$$

$$\begin{aligned} |\rho v'_h - (\Pi_h(\rho v_h))'|_{2,(1)} &= |(\rho - P_h \rho)v'_h - (\Pi_h((\rho - P_h \rho)v_h))'|_{2,(1)} \\ &\leq |(\rho - P_h \rho)v'_h|_{2,(1)} + |(\Pi_h((\rho - P_h \rho)v_h))'|_{2,(1)} \\ &\leq |(\rho - P_h \rho)v'_h|_{2,(1)} + Ch |(\rho - P_h \rho)v_h|_{3,(1)} \\ &\leq C(h|v_h|_{2,h} + |v_h|_{1,\Omega}). \end{aligned}$$

Thus we get

$$|a_h(\bar{v}_h, \rho\bar{v}_h) + a_h(v'_h, \rho v'_h)| \leq |b_h(v_h, \Pi_h(\rho v_h))| + C |v_h|_{2,(-1)} (h|v_h|_{2,h} + |v_h|_{1,\Omega}). \quad (3.19)$$

By (2.11), we have

$$\begin{aligned} |b_h(\Pi_h g - g_h, \Pi_h(\rho v_h))| &= |b_h(\Pi_h g, \Pi_h(\rho v_h)) - (\Delta^2 g, \Pi_h(\rho v_h))| \\ &\leq Ch|g|_{3,(-1)}|\Pi_h(\rho v_h)|_{2,(1)}. \end{aligned}$$

On the other hand,

$$\begin{aligned} |\Pi_h(\rho v_h)|_{2,(1)} &\leq C|\rho v_h|_{2,(1)} \\ &\leq C(|v_h|_{2,(-1)} + |v_h|_{1,\Omega} + |v_h|_{0,(1)}) \\ &\leq C(|v_h|_{2,(-1)} + |v_h|_{1,\Omega} + |\ln h|^{1/2}|v_h|_{0,\infty,\Omega}). \end{aligned}$$

From [4], we know that

$$|v_h|_{0,\infty,\Omega} \leq C|\ln h|^{1/2}|v_h|_{1,\Omega},$$

therefore,

$$|b_h(\Pi_h g - g_h, \Pi_h(\rho v_h))| \leq Ch|g|_{3,(-1)}(|v_h|_{2,(-1)} + (1 + |\ln h|)|v_h|_{1,\Omega}). \quad (3.20)$$

Combining (??) to (??) with  $v_h = \Pi_h g - g_h$ , we get

$$\begin{aligned} |\Pi_h g - g_h|_{2,(-1)}^2 &\leq C \left\{ h^2 |\Pi_h g - g_h|_{2,h}^2 + h |\Pi_h g - g_h|_{1,\Omega} |\Pi_h g - g_h|_{2,h} \right. \\ &\quad \left. + h |g|_{3,(-1)} (1 + |\ln h|) |\Pi_h g - g_h|_{1,\Omega} \right. \\ &\quad \left. + |\Pi_h g - g_h|_{2,(-1)} (h |v_h|_{2,h} + |\Pi_h g - g_h|_{1,\Omega} + h |g|_{3,(-1)}) \right\}. \end{aligned}$$

By (2.7), (2.9), (3.8) and (3.9), we have

$$|\Pi_h g - g_h|_{2,(-1)}^2 \leq Ch^2(1 + |\ln h|(1 + |\ln h|^{1/2})). \quad (3.21)$$

Inequality (??) is proved.

From (3.3) and (3.7), we see that

$$|\Pi_h u - u_h|_{1,\infty,\Omega} = |b_h(g_h, \Pi_h u - u_h)|. \quad (3.22)$$

From (1.11) and (1.12), we have

$$b_h(g_h, \Pi_h u - u_h) = a_h(g_h, \Pi_h u) - a_h(\bar{g}_h, (\Pi_h u)') - a_h(\overline{\Pi_h u}, g_h') - (\Delta^2 u, g_h).$$

From [3] we have

$$\begin{aligned} a_h(\bar{g}_h, (\Pi_h u)') &= E(\bar{g}_h, (\Pi_h u)') \\ a_h(\overline{\Pi_h u}, g_h') &= E(\overline{\Pi_h u}, g_h') \\ (\Delta^2 u, g_h) &= a_h(g_h, u) - E(u, g_h). \end{aligned}$$



Then

$$\begin{aligned} b_h(g_h, \Pi_h u - u_h) &= a_h(g - g_h, u - \Pi_h u) - (\Delta^2 g, u - \Pi_h u) + E(u - \overline{\Pi_h u}, g'_h) \\ &\quad + E(g - \bar{g}_h, (\Pi_h u)') + E(u, \overline{g_h - \Pi_h g}) + E(u, \overline{\Pi_h g} - g) \\ &\quad + E(g, \overline{\Pi_h u} - u) + E(u, g) + E(g, u), \end{aligned}$$

Noticing  $E(u, g) = E(g, u) = 0$ , by (2.8) and (2.12) to (2.14) we have,

$$\begin{aligned} |b_h(g_h, \Pi_h u - u_h)| &\leq C \left\{ |g - g_h|_{2,(-1)} |u - \Pi_h u|_{2,(1)} + |\delta_h|_{0,(-1)} |u - \Pi_h u|_{0,(1)} \right. \\ &\quad \left. + h^2 |u|_{3,(1)} |g|_{3,(-1)} + h(|u - \overline{\Pi_h u}|_{2,(1)} + h|u|_{3,(1)}) |g'_h|_{3,(-1)} \right. \\ &\quad \left. + h(|g - \bar{g}_h|_{2,(-1)} + h|g|_{3,(-1)}) |(\Pi_h u)'|_{3,(1)} \right\}. \end{aligned}$$

From (2.9), (3.5) and (3.9) to (3.11), we derive that

$$|b_h(g_h, \Pi_h u - u_h)| \leq Ch^2 |u|_{3,(1)} (1 + |\ln h|^{3/4} + |g'_h|_{3,(-1)}). \quad (3.23)$$

Using the inverse inequality of polynomials and (3.9) to (3.11), we have

$$\begin{aligned} |g'_h|_{3,(-1)} &\leq C |g_h|_{3,(-1)} \leq C (|\Pi_h g - g_h|_{3,(-1)} + |\Pi_h g|_{3,(-1)}) \\ &\leq C (h^{-1} |\Pi_h g - g_h|_{2,(-1)} + |\ln h|^{1/2}) \\ &\leq C |\ln h|^{3/4}. \end{aligned}$$

From (3.23), (2.7), we derive that

$$|b_h(g_h, \Pi_h u - u_h)| \leq Ch^2 |u|_{3,\infty,\Omega} |\ln h|^{5/4} \quad (3.24)$$

Inequality (3.1) follows from (3.24), (3.22) and (3.2).

## References

- [1] J.H. Argyris, M. Haase and H.P. Mlejnek, On an unconventional but natural formation of a stiffness matrix, *Comput. Meths. Appl.Mech. Eng.*, 22(1980), 1-22.
- [2] J.H. Argyris, M. Hasse and H.P. Mlejnek, Some considerations on the natural approach, *Comput. Meths. Appl. Mech. Engrg.*, 30(1982), 335-346.
- [3] Z.C. Shi, Convergence of the TRUNC plate element, *Comput. Meths. Appl. Mech. Eng.*, 62(1987), 71-88.
- [4] Wang M, On the inequalities for the maximum norm of nonconforming finite element spaces, *Mathematica Numerica Sinica*, 12(1990), in Chinese.
- [5] Wang M,  $L^\infty$  error estimates of nonconforming finite elements for the biharmonic equation, *J. Comput. Math.*, 11(1993), 276-288.
- [6] Zhang H.Q. and Wang M, The Mathematical Theory of Finite Element Methods, Science Press, Beijing, 1991, in chinese.