

CONVERGENCE OF VORTEX METHODS FOR 3-D EULER EQUATIONS*

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Abstract

In this paper we apply an approach introduced in [6] [7], where continuous norms and high order estimates and extension are used, to study the convergence of vortex methods for the 3-D Euler equations in bounded domains as the initial vorticity ω_0 and the curl of the body force f are non-compactly supported functions. Convergence results are proved.

Key words: Euler equations, Vortex methods, Convergence, Initial-boundary value problem.

1. Introduction

The convergence problem of vortex methods for the Euler equations has been studied by many authors. Hald and DelPrete proved the convergence for two-dimensional initial value problems [3]. Three-dimensional initial value problems were studied by Beale and Majda [2] and Beale [1]. Ying [4] and Ying and Zhang [5], [6] proved the convergence of vortex methods for two-dimensional initial-boundary value problems of the Euler equations. Ying [7] proved the convergence of vortex methods for three-dimensional initial-boundary value problems of the Euler equations under the assumption that the initial vorticity ω_0 and the curl of the body force f are compactly supported.

In this paper, we will prove the convergence of the vortex method for three-dimensional initial-boundary value problems without assuming that the ω_0 and $\nabla \times f$ are compactly supported. In contrast to [7], there are two new difficulties. One is how to extend the physical quantities such as velocity and the force function outside Ω . The other one is that the approximate velocity g^ϵ (ie. eqn. (20)) is no longer divergence-free outside Ω . We will use the approach in [4] [5] and [7] to perform the extension of the physical quantities outside Ω . Although $\nabla \cdot g^\epsilon \neq 0$ outside Ω , we will show that $\nabla \cdot g^\epsilon$ is small (see eqn. (48) (49)). This is sufficient for our convergence analysis.

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2. Formulation of Vortex Methods

We consider the initial-boundary value problems of inviscid incompressible flow as follows:

$$\frac{\partial u}{\partial t} + (u \cdot \nabla)u + \frac{\nabla p}{\rho} = f, \quad (1)$$

$$\nabla \cdot u = 0, \quad (2)$$

$$u \cdot n|_{x \in \partial\Omega} = 0, \quad (3)$$

$$u|_{t=0} = u_0(x), \quad (4)$$

where $u = (u_1, u_2, u_3) \in \mathbf{R}^3$ is velocity, $p \in \mathbf{R}$ is pressure, ρ is a constant standing for density, f is body force, u_0 is the initial distribution of the velocity satisfying $\nabla \cdot u_0 = 0$ and $u_0 \cdot n|_{\partial\Omega} = 0$ and ∇ is the gradient operator $(\frac{\partial}{\partial x_1}, \frac{\partial}{\partial x_2}, \frac{\partial}{\partial x_3})$. The domain Ω is assumed bounded with a sufficiently smooth boundary $\partial\Omega$ and n is the unit outward normal vector along the boundary. For simplicity we assume Ω is simply connected and convex. Let $\omega = \nabla \times u$ be the vorticity and F be the curl of the body force f , then applying the operator curl to the equation (1) and the initial condition (4), we obtain

$$\frac{\partial \omega}{\partial t} + (u \cdot \nabla)\omega - (\omega \cdot \nabla)u = F, \quad (5)$$

$$\omega|_{t=0} = \omega_0 \equiv \nabla \times u_0. \quad (6)$$

To get velocity u from the vorticity ω , we need the stream function, which is not unique for three-dimensional problems. We accept the formulation in [7]

$$-\Delta\psi + \nabla z = \omega, \quad u = \nabla \times \psi, \quad (7)$$

$$\nabla \cdot \psi = 0, \quad (8)$$

$$\psi \times n|_{x \in \partial\Omega} = 0, \quad z|_{x \in \partial\Omega} = 0, \quad (9)$$

where ψ is the stream function. Problems (7)-(9) admit a unique smooth solution for any smooth ω (see [7], §4).

We assume that the problem (1)-(4) admits a sufficiently smooth solution (u, p) on $\overline{\Omega} \times [0, T]$. Under this assumption, we consider the vortex method formulation.

Let positive constants h and ϵ be mesh sizes. $j = (j_1, j_2, j_3) \in \mathbf{Z}^3$, $B_j = \{x; j_i h < x_i < (j_i + 1)h, i = 1, 2, 3\}$, $X_j = ((j_1 + \frac{1}{2})h, (j_2 + \frac{1}{2})h, (j_3 + \frac{1}{2})h)$. We define a "vortex blob" function $\zeta(x)$ with a support in ball $\{x; |x| \leq 1\}$, which satisfies

$$\int \zeta(x) dx = 1. \quad (10)$$

Consider the following scheme: Set $\Omega_d = \{x; \text{dist}(x, \overline{\Omega}) < d\}$, where $d > 0$ is a parameter. We solve the problem in $\overline{\Omega}_d \times [0, T]$. u_0, F are not defined outside Ω and $\Omega \times [0, T]$ and we extend u_0, F in the following way. By (7)-(9), we get $\psi(x, 0)$, then extend it smoothly to \mathbf{R}^3 . $\psi(x, 0)$ needn't satisfy the divergence free condition outside Ω and we assume $\psi(x, 0)$ is compactly supported. Set $u_0 = \nabla \times \psi(x, 0)$, then u_0 is extended

too. u_0 is compactly supported and $\nabla \cdot u_0 = 0$ in \mathbf{R}^3 . F can be easily extended to $\mathbf{R}^3 \times [0, T]$ with a compact support.

The approximation of ω_0 is

$$\omega_0 \approx \sum_{j \in J} \alpha_j \zeta_\epsilon(x - X_j), \tag{11}$$

where $\alpha_j = \omega_0(X_j)h^3$, $\zeta_\epsilon(x) = \frac{1}{\epsilon^3} \zeta(\frac{x}{\epsilon})$, and $J = \{j; X_j \in \Omega_d\}$.

The approximate solution of (5)-(9) is solved as follows:

$$\omega^\epsilon(x, t) = \sum_{j \in J} \alpha_j^\epsilon(t) \zeta_\epsilon(x - X_j^\epsilon(t)), \tag{12}$$

$$\frac{d\alpha_j^\epsilon}{dt} = (\alpha_j^\epsilon(t) \cdot \nabla) g^\epsilon(X_j^\epsilon(t), t) + h^3 F(X_j^\epsilon(t), t), \tag{13}$$

$$\alpha_j^\epsilon(0) = \alpha_j, \tag{14}$$

$$\frac{dX_j^\epsilon}{dt} = g^\epsilon(X_j^\epsilon(t), t), \tag{15}$$

$$X_j^\epsilon(0) = X_j, \tag{16}$$

$$-\Delta \psi^\epsilon + \nabla z^\epsilon = \omega^\epsilon, \quad u^\epsilon = \nabla \times \psi^\epsilon, \tag{17}$$

$$\nabla \cdot \psi^\epsilon = 0, \tag{18}$$

$$\psi^\epsilon \times n |_{x \in \partial \Omega} = 0, \quad z^\epsilon |_{x \in \partial \Omega} = 0, \tag{19}$$

$$g^\epsilon(x, t) = \sum_{i=1}^M a_i u^\epsilon(x^{(i)}, t), \tag{20}$$

where a_i satisfies the following algebraic system

$$\sum_{i=1}^M (-i)^j a_i = 1, \quad (j = 0, 1, \dots, M - 1). \tag{21}$$

And the definition of $x^{(i)}$ is as follows:

$$\text{if } x \in \bar{\Omega}, \text{ then } x^{(i)} = x; \text{ otherwise } x^{(i)} = (i + 1)y - ix, \tag{22}$$

where y is the nearest point on $\partial \Omega$ to x . To solve (12)-(22), further discretization is needed. For instance (17)-(19) may be solved by a finite element scheme [8].

To prove the convergence of the scheme (12)-(22). We extend (u, p) in the following way. Let (u, p) be the solution of (1)-(4). By (7)-(9) we can determine ψ , then extend ψ smoothly to $\mathbf{R}^3 \times [0, T]$ and let ψ be compactly supported. Similarly, p is extended to $\mathbf{R}^3 \times [0, T]$. By (7) we can determine u , then by (1), we get an extension of f . It is different from the above extension. We write it as \tilde{f} , and $\tilde{F} = \nabla \times \tilde{f}$. We require the extension of ψ at $t = 0$ is corresponding with the extension of u_0 above. According to this extension, (u, p) is the solution of initial value problems. Convergence of vortex

method for three-dimensional initial value problems has been proved in [7]. We can use these results for initial value problems.

3. Convergence for Initial-Boundary Value Problems

First, we define the characteristic curves $\xi(t; \eta, l)$:

$$\frac{d}{dt}\xi(t; \eta, l) = u(\xi(t; \eta, l), t), \quad (23)$$

$$\xi(l; \eta, l) = \eta. \quad (24)$$

Similarly we can define $\xi^\epsilon(t; \eta, l)$ if the function u in (23) is replaced by g^ϵ . Set $x(\eta, t) = \xi(t; \eta, 0)$, $x^\epsilon(\eta, t) = \xi^\epsilon(t; \eta, 0)$, and $e(t) = x(\cdot, t) - x^\epsilon(\cdot, t)$ in $L^p(\Omega_d)$. The Sobolev norm $\|e(t)\|_{l,p,\Omega_d}$ is defined in the usual way. Furthermore, let $\alpha(\eta, t)$ and $\alpha^\epsilon(\eta, t)$ be the solutions of the following:

$$\frac{d\alpha}{dt} = (\alpha \cdot \nabla)u(x(\eta, t), t) + h^3 \tilde{F}(x(\eta, t), t), \quad (25)$$

$$\alpha(\eta, 0) = \omega_0(\eta)h^3, \quad (26)$$

$$\frac{d\alpha^\epsilon}{dt} = (\alpha^\epsilon \cdot \nabla)g^\epsilon(x^\epsilon(\eta, t), t) + h^3 F(x^\epsilon(\eta, t), t), \quad (27)$$

$$\alpha^\epsilon(\eta, 0) = \omega_0(\eta)h^3. \quad (28)$$

By (5)-(6), (13)-(16), (23)-(28) it is easy to see that $x^\epsilon(X_j, t) = X_j^\epsilon(t)$, $\alpha^\epsilon(X_j, t) = \alpha_j^\epsilon(t)$ and $\alpha(\eta, t) = \omega(x(\eta, t), t)h^3$. We define $\bar{\omega}(t) = (\alpha(\cdot, t) - \alpha^\epsilon(\cdot, t))/h^3$ in $L^p(\Omega_d)$. In the sequel, we will always assume that C is a generic constant which may not be the same in the different expressions and M_1, C_0, C_1, C_2 are some special constants.

Lemma 3.1. *If $p > 3$, and if there is a constant $C_1 > 0$, such that $\|u^\epsilon\|_{2,p,\Omega} \leq C_1$, then there exists a constant $C_0 > 0$, such that*

$$|\xi^\epsilon(t; \eta_1, l) - \xi^\epsilon(t; \eta_2, l)| \leq C_0 |\eta_1 - \eta_2|, \quad \forall \eta_1, \eta_2 \in \Omega_d, \quad t, l \in [0, T]. \quad (29)$$

Proof. Since Ω is convex, by (22), we can get $|\eta_1^{(i)} - \eta_2^{(i)}| \leq (i+2)|\eta_1 - \eta_2|$. By the assumption of the lemma $\|u^\epsilon\|_{2,p,\Omega} \leq C_1$. Using the embedding theorem, we have $\|u^\epsilon\|_{1,\infty,\Omega} \leq C$. By (23)-(24) (20), we have

$$\begin{aligned} \left| \frac{d}{dt}(\xi^\epsilon(t; \eta_1, l) - \xi^\epsilon(t; \eta_2, l)) \right| &= |g^\epsilon(\xi^\epsilon(t; \eta_1, l), t) - g^\epsilon(\xi^\epsilon(t; \eta_2, l), t)| \\ &\leq C |\xi^\epsilon(t; \eta_1, l) - \xi^\epsilon(t; \eta_2, l)|, \\ \xi^\epsilon(l; \eta_1, l) - \xi^\epsilon(l; \eta_2, l) &= \eta_1 - \eta_2. \end{aligned}$$

Using the Gronwall inequality, we obtain (29).

Lemma 3.2. *Under the assumption of Lemma 3.1, we have*

$$\|\bar{\omega}(t)\|_{0,p,\Omega_d} \leq C \int_0^t (\|e(t)\|_{0,p,\Omega_d} + \|u - u^\epsilon\|_{1,p,\Omega} + \|\bar{\omega}(t)\|_{0,p,\Omega_d} + d^{M-1}) dt, \quad (30)$$

$$\|\bar{\omega}(t)\|_{1,p,\Omega_d} \leq C \int_0^t (\|e(t)\|_{1,p,\Omega_d} + \|u - u^\epsilon\|_{2,p,\Omega} + \|\bar{\omega}(t)\|_{1,p,\Omega_d} + d^{M-2}) dt, \quad (31)$$

where the constant C is independent of d .

Proof. The equations (25) and (27) give

$$\begin{aligned} \frac{d(\alpha - \alpha^\epsilon)}{dt} &= (\alpha \cdot \nabla)(u(x(t), t) - u(x^\epsilon(t), t)) \\ &+ (\alpha \cdot \nabla)(u(x^\epsilon(t), t) - g^\epsilon(x^\epsilon(t), t)) \\ &+ ((\alpha - \alpha^\epsilon) \cdot \nabla)g^\epsilon(x^\epsilon(t), t) \\ &+ h^3 \tilde{F}(x(t), t) - h^3 F(x^\epsilon(t), t), \end{aligned} \quad (32)$$

where for simplicity we omit the independent variable η . Obviously $|\alpha/h^3| \leq C$. By assumption and the embedding theorem, $|\nabla u^\epsilon|_{0,\infty,\Omega} \leq C$. By (20)-(21), we have

$$\begin{aligned} &|\nabla \cdot (u(x^\epsilon(t), t) - g^\epsilon(x^\epsilon(t), t))| = \left| \sum_{i=1}^M a_i \nabla \cdot (u(x^\epsilon(t), t) - u^\epsilon(x^\epsilon(t)^{(i)}, t)) \right| \\ &\leq \left| \sum_{i=1}^M a_i \nabla \cdot (u(x^\epsilon(t), t) - u(x^\epsilon(t)^{(i)}, t)) \right| + \left| \sum_{i=1}^M a_i \nabla \cdot (u(x^\epsilon(t)^{(i)}, t) - u^\epsilon(x^\epsilon(t)^{(i)}, t)) \right| \\ &\leq C d^{M-1} + C \sum_{i=1}^M |D(u - u^\epsilon)(x^\epsilon(t)^{(i)}, t)|, \end{aligned}$$

where we have used Taylor expansion and (21) to get the last inequality. We define a mapping $\Phi^{(i)} : x \mapsto x^{(i)}$, then the second term of equation (32) can be estimated as the following:

$$|(\alpha \cdot \nabla)(u(x^\epsilon(t), t) - g^\epsilon(x^\epsilon(t), t))| \leq C h^3 \left(d^{M-1} + \sum_{i=1}^M |D(u - u^\epsilon) \circ \Phi^{(i)}(x^\epsilon(t))| \right).$$

The fourth term also can be estimated as follows.

$$\begin{aligned} &|h^3 \tilde{F}(x(t), t) - h^3 F(x^\epsilon(t), t)| \\ &\leq |h^3 \tilde{F}(x(t), t) - h^3 F(x(t), t)| + |h^3 F(x(t), t) - h^3 F(x^\epsilon(t), t)| \\ &\leq C h^3 d^M + C h^3 |(x - x^\epsilon)(t)|, \end{aligned}$$

where we also have used Taylor expansion and $F \equiv \tilde{F}$ in $\bar{\Omega} \times [0, T]$. Thus, by (32), we obtain

$$\begin{aligned} \|\bar{\omega}(t)\|_{0,p,\Omega_d} &\leq C \int_0^t (\|e(t)\|_{0,p,\Omega_d} + \sum_{i=1}^M \|D(u - u^\epsilon) \circ \Phi^{(i)}(x^\epsilon(t))\|_{0,p,\Omega_d} \\ &+ \|\bar{\omega}(t)\|_{0,p,\Omega_d} + \|e(t)\|_{0,p,\Omega_d} + d^{M-1}) dt. \end{aligned} \quad (33)$$

The equation (17) implies $\nabla \cdot u^\epsilon = 0$, hence the mapping $\eta \mapsto x^\epsilon(\eta, t)$ in Ω is measure preserving and by Lemma 3.1 the Jacobi matrix of $\eta \mapsto x^\epsilon(\eta, t)$ in $\Omega_d \setminus \Omega$ is bounded by C ; consequently;

$$\begin{aligned} & \|D(u - u^\epsilon) \circ \Phi^{(i)}(x^\epsilon(t))\|_{0,p,\Omega_d} \\ & \leq C \|D(u - u^\epsilon) \circ \Phi^{(i)}(x^\epsilon(t))\|_{0,p,\Omega} + C \|D(u - u^\epsilon) \circ \Phi^{(i)}(x^\epsilon(t))\|_{0,p,\Omega_d \setminus \Omega} \\ & \leq C (\|u - u^\epsilon\|_{1,p,\Omega} + \|D(u - u^\epsilon)\|_{0,p,\Omega}) = C \|u - u^\epsilon\|_{1,p,\Omega}. \end{aligned} \quad (34)$$

By substituting (34) into (33), we get (30).

Next, we apply the operator D_η to the equation (32). For notational convenience, we consider one component of α , u , x , and α^ϵ , u^ϵ , x^ϵ and denote by D the derivative with respect to one special variable.

$$\begin{aligned} & \frac{d}{dt} D(\alpha - \alpha^\epsilon) = D\alpha D(u(x(t), t) - u(x^\epsilon(t), t)) \\ & + \alpha(D^2 u(x(t), t) Dx(t) - D^2 u(x^\epsilon(t), t) Dx^\epsilon(t)) \\ & + D\alpha D(u(x^\epsilon(t), t) - g^\epsilon(x^\epsilon(t), t)) \\ & + \alpha(D^2 u(x^\epsilon(t), t) Dx^\epsilon(t) - D^2 g^\epsilon(x^\epsilon(t), t) Dx^\epsilon(t)) \\ & + D(\alpha - \alpha^\epsilon) Dg^\epsilon(x^\epsilon(t), t) \\ & + (\alpha - \alpha^\epsilon) D^2 g^\epsilon(x^\epsilon(t), t) Dx^\epsilon(t) \\ & + h^3 D\tilde{F}(x(t), t) Dx(t) - h^3 DF(x^\epsilon(t), t) Dx^\epsilon(t). \end{aligned} \quad (35)$$

To estimate the L^p - norm of the right-hand side, we need to prove $|Dx^\epsilon(t)| \leq C$. Indeed $Dx^\epsilon(t)$ satisfies

$$\frac{d}{dt} Dx^\epsilon(t) = Dg^\epsilon(x^\epsilon(t), t) Dx^\epsilon(t).$$

Using the Gronwall inequality, we can estimate $Dx^\epsilon(t)$. To estimate the sixth term of (35) we notice that

$$\begin{aligned} & \|(\alpha - \alpha^\epsilon) D^2 g^\epsilon(x^\epsilon(t), t) Dx^\epsilon(t)\|_{0,p,\Omega_d} \\ & \leq C \max(|\alpha - \alpha^\epsilon|) |g^\epsilon|_{2,p,\Omega_d} \leq C \|\alpha - \alpha^\epsilon\|_{1,p,\Omega_d} |u^\epsilon|_{2,p,\Omega}, \end{aligned} \quad (36)$$

where we have used the embedding theorem and (20). By the assumption of the lemma, the right-hand side of (36) is bounded by $C \|\alpha - \alpha^\epsilon\|_{1,p,\Omega_d}$. The other terms of (35) are estimated in a straightforward way, then we obtain (31).

Lemma 3.3. *Under the assumption of Lemma 3.1, we have*

$$\|e(t)\|_{1,p,\Omega_d} \leq C \int_0^t (\|u - u^\epsilon\|_{1,p,\Omega} + \|e(t)\|_{1,p,\Omega_d} + d^{M-1}) dt, \quad (37)$$

$$|e(t)|_{2,p,\Omega_d} \leq C \int_0^t (\|u - u^\epsilon\|_{2,p,\Omega} + \|e(t)\|_{2,p,\Omega_d} + d^{M-2}) dt. \quad (38)$$

Proof. From equation (23), we obtain

$$\begin{aligned} \frac{d}{dt}(x(t) - x^\epsilon(t)) &= u(x(t), t) - g^\epsilon(x^\epsilon(t), t) \\ &= u(x(t), t) - u(x^\epsilon(t), t) + \sum_{i=1}^M a_i(u(x^\epsilon(t), t) - u(x^\epsilon(t)^{(i)}, t)) \\ &\quad + \sum_{i=1}^M a_i(u(x^\epsilon(t)^{(i)}, t) - u^\epsilon(x^\epsilon(t)^{(i)}, t)), \end{aligned}$$

thus

$$\|e(t)\|_{0,p,\Omega_d} \leq C \int_0^t (\|e(t)\|_{0,p,\Omega_d} + d^M + \|u - u^\epsilon\|_{0,p,\Omega}) dt, \quad (39)$$

where we have also used the abbreviations in Lemma 3.2. By (23) we obtain

$$\begin{aligned} \frac{d}{dt}D(x(t) - x^\epsilon(t)) &= D(u - g^\epsilon)(x^\epsilon(t), t)Dx^\epsilon(t) \\ &\quad + Du(x(t), t)Dx(t) - Du(x^\epsilon(t), t)Dx^\epsilon(t), \\ \frac{d}{dt}D^2(x(t) - x^\epsilon(t)) &= D^2u(x(t), t)((Dx(t))^2 - (Dx^\epsilon(t))^2) \\ &\quad + (D^2u(x(t), t) - D^2g^\epsilon(x^\epsilon(t), t))(Dx^\epsilon(t))^2 + Dg^\epsilon(x^\epsilon(t), t)(D^2x(t) - D^2x^\epsilon(t)) \\ &\quad + (D(u(x(t), t) - Dg^\epsilon(x^\epsilon(t), t))D^2x(t)). \end{aligned}$$

Using the same deduction as (39), we obtain (37) and (38).

Set $C_2 > 0$, such that $u \equiv 0$ as $|x| \geq C_2$. We consider the following set: $J_1 = \{j; |jh| \leq C_2\}$, then J_1 is a finite set. We introduce an operator G as the following: For given $\omega(x)$, there is a unique $u(x)$ satisfies (7)-(9), then we write $u = G\omega$, then

$$\|G\omega\|_{m+1,p,\Omega} \leq C\|\omega\|_{m,p,\Omega} \quad m \geq 0. \quad (\text{see}[7], \text{\S}4 \text{ Lemma 4.1}) \quad (40)$$

Now we start to estimate $u - u^\epsilon$. As in [7], we make the following decomposition:

$$\begin{aligned} u - u^\epsilon &= v_1 + v_2 + v_3, \\ v_1 &= u - G(\omega(\cdot, t) * \zeta_\epsilon), \\ v_2 &= G \left(\omega(\cdot, t) * \zeta_\epsilon - \sum_{j \in J_1} \alpha_j(t) \zeta_\epsilon(\cdot - X_j(t)) \right), \\ v_3 &= G \left(\sum_{j \in J_1} \alpha_j(t) \zeta_\epsilon(\cdot - X_j(t)) - \sum_{j \in J} \alpha_j^\epsilon(t) \zeta_\epsilon(\cdot - X_j^\epsilon(t)) \right), \end{aligned}$$

where $\alpha_j(t) = \alpha(\eta, X_j, 0)$, $X_j(t) = x(X_j, t)$.

Lemma 3.4. *If there is an integer $k \geq 1$ such that*

$$\int_{\mathbf{R}^3} x^r \zeta(x) dx = 0, \quad 1 \leq |r| \leq k - 1, \quad (41)$$

then

$$|v_1(\cdot, t)|_{l,p,\Omega} \leq C\epsilon^k, 1 \leq p \leq +\infty,$$

for any integer $l \geq 0$. (see[7], §4 Lemma 4.2)

Lemma 3.5. If $\zeta \in W^{m+l-1,\infty}(\mathbf{R}^3)$, $m \geq 1$, $l \geq 1$, then for any $r \in [1, 3/2]$

$$\|v_2(\cdot, t)\|_{l,p,\Omega} \leq C \left(\left(1 + \frac{h}{\epsilon}\right)^{3/r} \frac{h^m}{\epsilon^{m+l-1}} \right), \quad 1 \leq p < +\infty, \tag{42}$$

as $t \in [0, T]$. (see [7], §4 Lemma 4.3)

We are now in a position to estimate v_3 . Set $J_2 = \{j; X_j \in \Omega_{C_0\epsilon} \cap \Omega_d\}$, where C_0 is the constant in Lemma 3.1.

Lemma 3.6. Under the assumption of Lemma 3.1, if $p > 3$, $l \geq 1$, $\zeta \in W^{l+2,\infty}(\mathbf{R}^3)$, $\|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}} \leq M_1\epsilon$, $h \leq C_2\epsilon$, then

$$\begin{aligned} \|v_3\|_{l,p,\Omega} \leq & \frac{C}{\epsilon^{l-1}} \left\{ \left(1 + \frac{h + \|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}}}{\epsilon^2}\right) \|e(t)\|_{0,p,\Omega_{C_0\epsilon}} + \frac{h}{\epsilon} \|e(t)\|_{2,p,\Omega_{C_0\epsilon}} \right. \\ & \left. + \|\bar{\omega}(t)\|_{0,p,\Omega_{C_0\epsilon}} + h|\bar{\omega}(t)|_{1,p,\Omega_{C_0\epsilon}} + \int_0^t |u - u^\epsilon|_{1,p,\Omega} dt + \epsilon^{M-1} \right\}. \tag{43} \end{aligned}$$

Proof. First, we prove, if $C_0\epsilon < d$, then

$$v_3 = G \left(\sum_{j \in J_2} \alpha_j(t)\zeta_\epsilon(\cdot - X_j(t)) - \sum_{j \in J_2} \alpha_j^\epsilon(t)\zeta_\epsilon(\cdot - X_j^\epsilon(t)) \right). \tag{44}$$

Indeed, if $\text{supp}\zeta_\epsilon(\cdot - X_j^\epsilon(t)) \cap \Omega \neq \emptyset$, then there exists a point $x \in \bar{\Omega}$, such that $|x - X_j^\epsilon(t)| < \epsilon$. By Lemma 3.1 $|\xi^\epsilon(0; x, t) - X_j| < C_0\epsilon$ and $\xi^\epsilon(0; x, t) \in \bar{\Omega}$, thus $X_j \in \Omega_{C_0\epsilon}$. The same reason can be stated for $\zeta_\epsilon(\cdot - X_j(t))$. So we know if $X_j \notin \Omega_{C_0\epsilon}$, then $\zeta_\epsilon(x - X_j(t)) = 0$, $\zeta_\epsilon(x - X_j^\epsilon(t)) = 0, \forall x \in \Omega$, which proves (44).

Let $\omega_3 = \sum_{j \in J_2} (\alpha_j(t)\zeta_\epsilon(x - X_j(t)) - \alpha_j^\epsilon(t)\zeta_\epsilon(x - X_j^\epsilon(t)))$, then $v_3 = G\omega_3$. Let us estimate $\|\partial^\gamma \omega_3\|_{0,p,\Omega}$, where $|\gamma| = l - 1$. The function ω_3 can be further decomposed into

$$\begin{aligned} \omega^{(1)} &= \sum_{j \in J_2} (\zeta_\epsilon(x - X_j(t)) - \zeta_\epsilon(x - X_j^\epsilon(t)))\alpha_j(t), \\ \omega^{(2)} &= \sum_{j \in J_2} \zeta_\epsilon(x - X_j(t))(\alpha_j(t) - \alpha_j^\epsilon(t)), \\ \omega^{(3)} &= \sum_{j \in J_2} (\zeta_\epsilon(x - X_j(t)) - \zeta_\epsilon(x - X_j^\epsilon(t)))(\alpha_j^\epsilon(t) - \alpha_j(t)). \end{aligned}$$

Let us first estimate $\omega^{(1)}$. $\omega^{(1)}$ also can be written as

$$\begin{aligned} \omega^{(1)} &= \int_{\Omega_{C_0\epsilon}} (\zeta_\epsilon(x - \xi(t; X_j, 0)) - \zeta_\epsilon(x - \xi^\epsilon(t; X_j, 0)))\omega(\xi(t; X_j, 0), t) d\eta = I_1 + I_2 + I_3, \\ I_1 &= \int_{\Omega_{C_0\epsilon}} (\zeta_\epsilon(x - \xi(t; X_j, 0)) - \zeta_\epsilon(x - \xi^\epsilon(t; X_j, 0)))\omega(\xi(t; X_j, 0), t) \end{aligned}$$

$$-(\zeta_\epsilon(x - \xi(t; \eta, 0)) - \zeta_\epsilon(x - \xi^\epsilon(t; \eta, 0)))\omega(\xi(t; X_j, 0), t)d\eta,$$

$$I_2 = \int_{\Omega_{C_0\epsilon}} (\zeta_\epsilon(x - \xi(t; \eta, 0)) - \zeta_\epsilon(x - \xi^\epsilon(t; \eta, 0)))(\omega(\xi(t; X_j, 0), t) - \omega(\xi(t; \eta, 0), t))d\eta,$$

$$I_3 = \int_{\Omega_{C_0\epsilon}} (\zeta_\epsilon(x - \xi(t; \eta, 0)) - \zeta_\epsilon(x - \xi^\epsilon(t; \eta, 0)))\omega(\xi(t; \eta, 0), t)d\eta.$$

The following is the same as the estimate of $\omega^{(1)}$ in Lemma 3.6 in [7],

$$\|\partial^\gamma I_1\|_{0,p,\Omega} \leq \frac{C}{\epsilon^{l-1}} \left\{ \left(\frac{h + \|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}}}{\epsilon^2} \right) \|e(t)\|_{0,p,\Omega_{C_0\epsilon}} + \frac{h}{\epsilon} \|e(t)\|_{2,p,\Omega_{C_0\epsilon}} \right\} \quad (45)$$

$$|\partial^\gamma I_2|_{0,p,\Omega} \leq \frac{C}{\epsilon^{l-1}} \|e(t)\|_{0,p,\Omega_{C_0\epsilon}}. \quad (46)$$

As for the third term, since the mapping $\eta \mapsto \xi(t; \eta, 0)$ is measure preserving, we have

$$\begin{aligned} I_3 &= \int_{\Omega_{C_0\epsilon}} \zeta_\epsilon(x - \xi(t; \eta, 0))\omega(\xi(t; \eta, 0), t)d\eta - \int_{\Omega_{C_0\epsilon}} \zeta_\epsilon(x - \xi^\epsilon(t; \eta, 0))\omega(\xi(t; \eta, 0), t)d\eta \\ &= \int_{\Omega_{C_0\epsilon}} \zeta_\epsilon(x - \xi)\omega(\xi, t)d\xi - \int_{\Omega_{C_0\epsilon}} \zeta_\epsilon(x - \xi)\omega(\xi(t; \xi^\epsilon(0; \xi, t), 0), t)|J|d\xi \\ &= A + B, \end{aligned}$$

where

$$A = \int_{\Omega_{C_0\epsilon}} \zeta_\epsilon(x - \xi)(\omega(\xi, t) - \omega(\xi(t; \xi^\epsilon(0; \xi, t), 0), t))d\xi,$$

$$B = \int_{\Omega_{C_0\epsilon}} \zeta_\epsilon(x - \xi)\omega(\xi(t; \xi^\epsilon(0; \xi, t), 0), t)(|J| - 1)d\xi.$$

J is the Jacobi matrix of $\eta \mapsto \xi^\epsilon(t; \eta, 0)$. Analogous to the estimate of $\omega^{(1)}$ in Lemma 3.6 in [7], we have

$$\|\partial^\gamma A\|_{0,p,\Omega} \leq \frac{C}{\epsilon^{l-1}} \|e(t)\|_{0,p,\Omega_{C_0\epsilon}}. \quad (47)$$

Applying Lemma 5.1 in [6], Chap2, §5, we obtain

$$|J| = \exp \left(\int_0^t \nabla \cdot g^\epsilon(x^\epsilon(\eta, \tau), \tau)d\tau \right),$$

and

$$\|\nabla g^\epsilon\|_{0,\infty,\Omega_{C_0\epsilon}} \leq C \|Du^\epsilon\|_{0,\infty,\Omega} \leq C \|u^\epsilon\|_{2,p,\Omega} \leq C,$$

where we have used (20), embedding theorem and the assumption of the lemma. Thus

$$\begin{aligned} ||J| - 1| &= \left| \exp \left(\int_0^t \nabla \cdot g^\epsilon(x^\epsilon(\eta, \tau), \tau)d\tau \right) - 1 \right| \\ &\leq C \left| \int_0^t \nabla \cdot g^\epsilon(x^\epsilon(\eta, \tau), \tau)d\tau \right| \\ &\leq C \int_0^t |\nabla \cdot g^\epsilon(x^\epsilon(\eta, \tau), \tau) - \nabla \cdot g(x^\epsilon(\eta, \tau), \tau)|d\tau + C \int_0^t |\nabla \cdot g(x^\epsilon(\eta, \tau), \tau)|d\tau, \end{aligned} \quad (48)$$

where $g(x, t) = \sum_{i=1}^M a_i u(x^{(i)}, t)$, as $x \in \Omega_{C_0\epsilon}$, so $g(x, t) \equiv u(x, t)$ as $x \in \bar{\Omega}$. Noting that $\nabla \cdot u = 0$, $(x, t) \in \mathbf{R}^3 \times [0, T]$ and applying Taylor expansion, we obtain

$$|\nabla \cdot g(x, t)| = |\nabla \cdot g(x, t) - \nabla \cdot u(x, t)| = \left| \nabla \cdot \sum_{i=1}^M a_i (u(x^{(i)}, t) - u(x, t)) \right| \leq C\epsilon^{M-1}, \quad (49)$$

$$\forall x \in \Omega_{C_0\epsilon}.$$

Noting that $|J| = 1$ as $\eta \in \Omega$, by (48)-(49), we have

$$\begin{aligned} |\partial^\gamma B| &= \left| \int_{\Omega_{C_0\epsilon}} \partial^\gamma \zeta_\epsilon(x - \xi) \omega(\xi(t; \xi^\epsilon(0; \xi, t), 0), t) (|J| - 1) d\xi \right| \\ &= \left| \int_{\Omega_{C_0\epsilon} \setminus \Omega} \partial^\gamma \zeta_\epsilon(x - \xi) \omega(\xi(t; \xi^\epsilon(0; \xi, t), 0), t) (|J| - 1) d\xi \right| \\ &\leq C \int_{\Omega_{C_0\epsilon} \setminus \Omega} \left| \partial^\gamma \zeta_\epsilon(x - \xi) \int_0^t (\nabla \cdot g^\epsilon(x^\epsilon(\eta, \tau), \tau) - \nabla \cdot g(x^\epsilon(\eta, \tau), \tau)) d\tau \right| d\xi \\ &\quad + C \int_{\Omega_{C_0\epsilon}} |\partial^\gamma \zeta_\epsilon(x - \xi)| \epsilon^{M-1} d\xi. \end{aligned}$$

Applying Lemma 3.5 in [7], we obtain

$$\|\partial^\gamma B\|_{0,p,\Omega} \leq \frac{C}{\epsilon^{l-1}} \left(\int_0^t \|u - u^\epsilon\|_{1,p,\Omega} dt + \epsilon^{M-1} \right). \quad (50)$$

Combining (45)-(47) and (50), we get

$$\begin{aligned} \|\partial^\gamma \omega^{(1)}\|_{0,p,\Omega} &\leq \frac{C}{\epsilon^{l-1}} \left\{ \left(1 + \frac{h + \|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}}}{\epsilon^2} \right) \|e(t)\|_{0,p,\Omega_{C_0\epsilon}} + \frac{h}{\epsilon} \|e(t)\|_{2,p,\Omega_{C_0\epsilon}} \right. \\ &\quad \left. + \int_0^t \|u - u^\epsilon\|_{1,p,\Omega} dt + \epsilon^{M-1} \right\}. \end{aligned}$$

Analogous to the estimate of $\omega^{(2)}$, $\omega^{(3)}$ in Lemma 3.6 in [7], we have

$$\begin{aligned} \|\partial^\gamma \omega^{(2)}\|_{0,p,\Omega} &\leq \frac{C}{\epsilon^{l-1}} \left(\|\bar{\omega}(t)\|_{0,p,\Omega_{C_0\epsilon}} + h|\bar{\omega}(t)|_{1,p,\Omega_{C_0\epsilon}} \right), \\ \|\partial^\gamma \omega^{(3)}\|_{0,p,\Omega} &\leq \frac{C}{\epsilon^{l-1}} \left(\|\bar{\omega}(t)\|_{0,p,\Omega_{C_0\epsilon}} + h|\bar{\omega}(t)|_{1,p,\Omega_{C_0\epsilon}} \right). \end{aligned}$$

Thus we obtain the estimate for ω_3

$$\begin{aligned} \|\partial^\gamma \omega_3\|_{0,p,\Omega} &\leq \frac{C}{\epsilon^{l-1}} \left\{ \left(1 + \frac{h + \|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}}}{\epsilon^2} \right) \|e(t)\|_{0,p,\Omega_{C_0\epsilon}} + \frac{h}{\epsilon} \|e(t)\|_{2,p,\Omega_{C_0\epsilon}} \right. \\ &\quad \left. + \|\bar{\omega}(t)\|_{0,p,\Omega_{C_0\epsilon}} + h|\bar{\omega}(t)|_{1,p,\Omega_{C_0\epsilon}} + \int_0^t \|u - u^\epsilon\|_{1,p,\Omega} dt + \epsilon^{M-1} \right\}. \end{aligned}$$

By (40), we get

$$\|v_3\|_{l,p,\Omega} = \|G\omega_3\|_{l,p,\Omega} \leq C\|\omega_3\|_{l-1,p,\Omega},$$

which proves the inequality (43).

Finally, let us prove the convergence theorem.

Theorem 3.1. *If $p > 3$, (41) holds for $k = M - 1 \geq 3$, $\zeta \in W^{m+1,\infty}(\mathbf{R}^3)$, $m \geq 3$ and $h \leq C_2\epsilon^2$, then there is a constant $\epsilon_0 > 0$, such that if $\epsilon \leq \epsilon_0$, then*

$$\|u - u^\epsilon\|_{1,p,\Omega} + \|e(t)\|_{1,p,\Omega_{C_0\epsilon}} + \|\bar{w}(t)\|_{0,p,\Omega_{C_0\epsilon}} \leq C \left(\epsilon^k + \frac{h^m}{\epsilon^m} \right), \quad (51)$$

$$\|u - u^\epsilon\|_{2,p,\Omega} + \|e(t)\|_{2,p,\Omega_{C_0\epsilon}} + \|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}} \leq C \left(\epsilon^{k-1} + \frac{h^m}{\epsilon^{m+1}} \right). \quad (52)$$

Proof. Let $t = 0$, then $X_j(0) = X_j^\epsilon(0) = X_j$, $\alpha_j(0) = \alpha_j^\epsilon(0) = \alpha_j$, so $v_3(x, 0) = 0$, $u(x, 0) - u^\epsilon(x, 0) = v_1(x, 0) + v_2(x, 0)$. By Lemma 3.4 and 3.5, $\|u^\epsilon(\cdot, 0)\|_{2,p,\Omega} \leq \|u(\cdot, 0)\|_{2,p,\Omega} + C$. Set a constant $C_3 > \|u(\cdot, 0)\|_{2,p,\Omega} + C$, then, by continuity, $\|u^\epsilon(\cdot, t)\|_{2,p,\Omega} < C_3$ in a neighbourhood of $t = 0$. We again fix a constant $M_2 > 0$. Since $\|e(0)\|_{0,\infty,\Omega_{C_0\epsilon}} = 0$, $\|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}} < M_2\epsilon^2$ in a neighbourhood, too. Let the intersection of these two intervals be $[0, T_*]$. We notice that T_* depends on h and ϵ .

We define the following norms:

$$\begin{aligned} \|u - u^\epsilon\|_{2,p,\Omega}^* &= \|u - u^\epsilon\|_{1,p,\Omega} + \epsilon|u - u^\epsilon|_{2,p,\Omega}, \\ \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* &= \|e(t)\|_{1,p,\Omega_{C_0\epsilon}} + \epsilon|e(t)|_{2,p,\Omega_{C_0\epsilon}}, \\ \|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* &= \|\bar{w}(t)\|_{0,p,\Omega_{C_0\epsilon}} + \epsilon|\bar{w}(t)|_{1,p,\Omega_{C_0\epsilon}}. \end{aligned}$$

We may assume $\epsilon \leq 1$, then Lemma 3.2 gives

$$\|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* \leq C \left(\int_0^t \|e(t)\|_{1,p,\Omega_{C_0\epsilon}} + \|u - u^\epsilon\|_{2,p,\Omega}^* + \|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* dt + \epsilon^{M-1} \right). \quad (53)$$

Lemma 3.3 gives

$$\|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* \leq C \left(\int_0^t \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* + \|u - u^\epsilon\|_{2,p,\Omega}^* dt + \epsilon^{M-1} \right). \quad (54)$$

And Lemma 3.4, 3.5 and 3.6 give

$$\|u - u^\epsilon\|_{2,p,\Omega}^* \leq C \left\{ \epsilon^k + \frac{h^m}{\epsilon^m} + \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* + \|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* + \int_0^t \|u - u^\epsilon\|_{2,p,\Omega}^* dt \right\}, \quad (55)$$

for $t \in [0, T_*]$. Applying Gronwall inequality, we obtain

$$\|u - u^\epsilon\|_{2,p,\Omega}^* \leq C \left\{ \epsilon^k + \frac{h^m}{\epsilon^m} + \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* + \|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* \right\}. \quad (56)$$

By substituting (56) into (53) and (54), we get

$$\|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* \leq C \int_0^t \left(\epsilon^k + \frac{h^m}{\epsilon^m} + \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* + \|\bar{w}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* \right) dt,$$

$$\|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* \leq C \int_0^t \left(\epsilon^k + \frac{h^m}{\epsilon^m} + \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* + \|\bar{\omega}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* \right) dt.$$

Applying Gronwall inequality again, we obtain

$$\|\bar{\omega}(t)\|_{1,p,\Omega_{C_0\epsilon}}^* + \|e(t)\|_{2,p,\Omega_{C_0\epsilon}}^* \leq C \left(\epsilon^k + \frac{h^m}{\epsilon^m} \right).$$

Equation (56) gives

$$\|u - u^\epsilon\|_{2,p,\Omega}^* \leq C \left(\epsilon^k + \frac{h^m}{\epsilon^m} \right),$$

thus (51) and (52) hold for $t \in [0, T_*]$. We notice that the constant C is independent of h and ϵ .

By virtue of the embedding theorem, we have

$$\|e(t)\|_{0,\infty,\Omega_{C_0\epsilon}} \leq C \left(\epsilon^k + \frac{h^m}{\epsilon^m} \right) = C \left(\epsilon^{k-2} + \frac{h^m}{\epsilon^{m+2}} \right) \epsilon^2.$$

We take $\epsilon_0 \leq 1$ small enough such that $C \left(\epsilon^{k-2} + \frac{h^m}{\epsilon^{m+2}} \right) < M_2$ and $\|u\|_{2,p,\Omega} + C \left(\epsilon^{k-1} + \frac{h^m}{\epsilon^{m+1}} \right) < C_3$ if $\epsilon \leq \epsilon_0$, then it is easy to show that T_* is indeed independent of ϵ and h and equal to T . Now the proof is complete.

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