# HSS METHOD WITH A COMPLEX PARAMETER FOR THE SOLUTION OF COMPLEX LINEAR SYSTEM* 

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#### Abstract

In this paper, a complex parameter is employed in the Hermitian and skew-Hermitian splitting (HSS) method (Bai, Golub and Ng: SIAM J. Matrix Anal. Appl., 24(2003), $603-626$ ) for solving the complex linear system $A x=f$. The convergence of the resulting method is proved when the spectrum of the matrix $A$ lie in the right upper (or lower) part of the complex plane. We also derive an upper bound of the spectral radius of the HSS iteration matrix, and a estimated optimal parameter $\alpha$ (denoted by $\alpha_{\text {est }}$ ) of this upper bound is presented. Numerical experiments on two modified model problems show that the HSS method with $\alpha_{\text {est }}$ has a smaller spectral radius than that with the real parameter which minimizes the corresponding upper bound. In particular, for the 'dominant' imaginary part of the matrix $A$, this improvement is considerable. We also test the GMRES method preconditioned by the HSS preconditioning matrix with our parameter $\alpha_{\text {est }}$.


Mathematics subject classification: 65F10,65Y20.
Key words: Hermitian matrix, Skew-Hermitian matrix, Splitting iteration method, Complex linear system, Complex parameter.

## 1. Introduction

We are interested in the iterative solution of the following complex linear system

$$
\begin{equation*}
A x=f \tag{1.1}
\end{equation*}
$$

We consider the case in which $A \in C^{n \times n}$ is large, sparse, non-Hermitian and positive definite and $f \in C^{n}$; see several applications in $[12,17,21]$.

Bai, Golub and $\mathrm{Ng}[6]$ proposed the Hermitian/skew-Hermitian splitting (HSS) method based on the fact that the matrix $A$ naturally possesses the Hermitian/skew-Hermitian splitting

$$
A=H+S
$$

where $H=\frac{1}{2}\left(A+A^{H}\right)$ is the Hermitian matrix, $S=\frac{1}{2}\left(A-A^{H}\right)$ is the skew-Hermitian matrix, and $A^{H}$ is the conjugate transpose of the matrix $A$. The HSS method has the following form:

$$
\left\{\begin{array}{l}
(\alpha I+H) x^{\left(k+\frac{1}{2}\right)}=(\alpha I-S) x^{(k)}+f  \tag{1.2}\\
(\alpha I+S) x^{(k+1)}=(\alpha I-H) x^{\left(k+\frac{1}{2}\right)}+f
\end{array}\right.
$$

where the parameter $\alpha>0$ can be chosen. The above form can be equivalently rewritten as

$$
\begin{equation*}
x^{(k+1)}=T(\alpha) x^{(k)}+G(\alpha) f, \quad k=0,1,2, \cdots, \tag{1.3}
\end{equation*}
$$

[^0]where $T(\alpha)=(\alpha I+S)^{-1}(\alpha I-H)(\alpha I+H)^{-1}(\alpha I-S)$ is the iteration matrix, and $G(\alpha)=$ $2 \alpha(\alpha I+S)^{-1}(\alpha I+H)^{-1}$.

The following theorem [6] gives the convergence property of the HSS iteration.
Theorem 1.1. Suppose that $A \in C^{n \times n}$ is a positive definite matrix, $H=\frac{1}{2}\left(A+A^{H}\right), S=$ $\frac{1}{2}\left(A-A^{H}\right)$ are the Hermitian and Skew-Hermitian parts of $A$ respectively, and the parameter $\alpha>0$. Then the spectral radius $\rho(T(\alpha))$ of the iteration matrix $T(\alpha)$ of the HSS iteration is bounded by

$$
\begin{equation*}
\rho(T(\alpha)) \leq \sigma(\alpha)=\max _{\lambda_{j} \in \Lambda(H)}\left|\frac{\alpha-\lambda_{j}}{\alpha+\lambda_{j}}\right|, \tag{1.4}
\end{equation*}
$$

where $\Lambda(\cdot)$ represents the spectrum of the corresponding matrix. Since $A$ is positive definite ( $\lambda_{j}>0$ ), we have

$$
\rho(T(\alpha)) \leq \sigma(\alpha)<1, \quad \text { for all } \alpha>0
$$

i.e., the HSS iteration converges.

Furthermore, let $\lambda_{1} \geq \cdots \geq \lambda_{n}>0$ be the eigenvalues of $H$. Then the upper bound $\sigma(\alpha)$ has the optimal parameter

$$
\begin{equation*}
\tilde{\alpha}=\sqrt{\lambda_{1} \lambda_{n}} \tag{1.5}
\end{equation*}
$$

and

$$
\sigma(\tilde{\alpha})=\min _{\alpha>0} \sigma(\alpha)=\frac{\sqrt{\kappa(H)}-1}{\sqrt{\kappa(H)}+1},
$$

where $\kappa(H)=\frac{\lambda_{1}}{\lambda_{n}}$ is the spectral condition number of $H$.
However, we have the following observations:
(1) $\tilde{\alpha}$ is usually different from the optimal parameter

$$
\alpha^{*}=\arg \min _{\alpha>0} \rho(T(\alpha))
$$

(2) Numerical experiments in $[5,6,10,11]$ have shown that in most situations,

$$
\rho\left(T\left(\alpha^{*}\right)\right) \ll \rho(T(\tilde{\alpha})) .
$$

(3) $\tilde{\alpha}$ and $\sigma(\tilde{\alpha})$ do not include any information of $S$.

To further improve the efficiency of the HSS method, it is desirable to determine or find a good estimate for the optimal parameter $\alpha^{*}$. For some special constructed matrices, in particular, for saddle-point problems, the optimal parameter, or the quasi-optimal parameter [2], has been extensively discussed $[2,4,5,8,10]$, and the results show that the optimal parameter does include the information of $S$.

The matrix

$$
\begin{equation*}
P=\frac{1}{2 \alpha}(H+\alpha I)(S+\alpha I) \tag{1.6}
\end{equation*}
$$

can also be employed as a preconditioner [2,8,11,23], where $\alpha$ is referred to as the preconditioning parameter. The idea of HSS preconditioner is motivated from the HSS method.

More generally, the coefficient matrix $A \in C^{n \times n}$ can be splitted into

$$
A=N+S_{0}
$$

where $N$ is a normal matrix and $S_{0}$ is a skew-Hermitian matrix. Similarly to the HSS method, normal/skew-Hermitian splitting (NSS) method with a real parameter could be formed [7].

When $A$ is a positive definite matrix and $\alpha$ is a positive parameter, the NSS method converges, and the spectral radius $\rho\left(M_{0}(\alpha)\right)$ of the iteration matrix $M_{0}(\alpha)$ of the NSS iteration is bounded by

$$
\rho\left(M_{0}(\alpha)\right) \leq \sigma_{0}(\alpha)=\max _{\gamma_{j}+i \eta_{j} \in \Lambda(N)} \sqrt{\frac{\left(\alpha-\gamma_{j}\right)^{2}+\eta_{j}^{2}}{\left(\alpha+\gamma_{j}\right)^{2}+\eta_{j}^{2}}}
$$

In [7], the optimal real parameter of the upper bound $\sigma_{0}(\alpha)$ is also discussed, and the optimal upper bound of the contraction factor of the HSS iteration is the smallest among all NSS iterations.

In this paper, we employ a complex parameter $\alpha$ in the HSS iteration (1.2) for solving the complex linear system (1.1). This idea is natural as it does not increase the computational complexity of the HSS method for the complex linear systems. We show that the resulting method converges when the spectrum of the matrix $A$ lie in the right upper (or lower) part of the complex plane. An upper bound of the spectral radius $\rho(T(\alpha))$ of the HSS iteration matrix $T(\alpha)$ is given. This upper bound includes the spectral information of the matrix $S$. Moreover, a estimated optimal parameter $\alpha_{\text {est }}$ of this upper bound is presented. Numerical experiments on two modified model problems show that the HSS method with $\alpha_{\text {est }}$ has a smaller spectral radius than that with the real parameter that minimizes the upper bound (1.4). In particular, for the 'dominant' imaginary part of the matrix $A$ (see Experiment 2), this improvement is considerable. In Experiment 2, we also test the GMRES method preconditioned by the HSS preconditioner with $\alpha_{e s t}$, and investigate how sensitive are the estimated parameter $\alpha_{\text {est }}$ and $\rho\left(T\left(\alpha_{e s t}\right)\right)$ with respect to the spectral information of $H$ and $S$.

## 2. HSS Method with Complex Parameter

We still consider the HSS iteration (1.2) for the solution of the complex linear system (1.1), but now the parameter $\alpha$ is complex.

Since the matrix $S$ is skew-Hermitian, its eigenvalues are imaginary numbers or zero [22], denoted by $\lambda(S)=i \tau_{j}, i=\sqrt{-1}, \tau_{j} \in R, j=1,2, \cdots, n$. Suppose that the spectrum of the $n \times n$ complex matrix $A$ lie in the right upper (or lower) part of the complex plane, i.e., $A$ is a positive definite matrix, and all $\tau_{j} \geq 0$ (or $\tau_{j} \leq 0$ ). This assumption is needed for the convergence of the HSS method with a complex parameter $\alpha$.

Theorem 2.1. Suppose that the spectrum of the $n \times n$ complex matrix $A$ lie in the right upper (or lower) part of the complex plane. The parameter $\alpha=a+i b, a, b \in R$ is chosen such that $a>0$ and $b \cdot \tau_{j} \geq 0, j=1, \cdots, n$. Then the HSS iteration (1.2) with the complex parameter $\alpha$ converges. Furthermore, it holds that

$$
\begin{equation*}
\rho(T(\alpha)) \leq \omega(\alpha) \equiv \max _{\lambda_{j} \in \Lambda(H)}\left|\frac{\alpha-\lambda_{j}}{\alpha+\lambda_{j}}\right| \max _{i \tau_{j} \in \Lambda(S)}\left|\frac{\alpha-i \tau_{j}}{\alpha+i \tau_{j}}\right|<1 . \tag{2.1}
\end{equation*}
$$

Proof. Let $\tilde{T}(\alpha)=(\alpha I+S) T(\alpha)(\alpha I+S)^{-1}=(\alpha I-H)(\alpha I+H)^{-1}(\alpha I-S)(\alpha I+S)^{-1}$. Then

$$
\begin{aligned}
\rho(T(\alpha)) & =\rho\left((\alpha I-H)(\alpha I+H)^{-1}(\alpha I-S)(\alpha I+S)^{-1}\right) \\
& \leq\left\|(\alpha I-H)(\alpha I+H)^{-1}\right\|_{2}\left\|(\alpha I-S)(\alpha I+S)^{-1}\right\|_{2}
\end{aligned}
$$

Since $S$ is skew-Hermitian, there exists a unitary matrix $U$, such that $S=U \Lambda U^{H}$, where $\Lambda=\operatorname{diag}\left(i \tau_{1}, \cdots, i \tau_{n}\right)$. It follows that

$$
\begin{aligned}
\left\|(\alpha I-S)(\alpha I+S)^{-1}\right\|_{2} & =\left\|U(\alpha I-\Lambda)(\alpha I+\Lambda)^{-1} U^{H}\right\|_{2} \\
& =\left\|(\alpha I-\Lambda)(\alpha I+\Lambda)^{-1}\right\|_{2} \\
& =\max _{i \tau_{j} \in \Lambda(S)}\left|\frac{\alpha-i \tau_{j}}{\alpha+i \tau_{j}}\right|
\end{aligned}
$$

Similarly, it holds that $\left\|(\alpha I-H)(\alpha I+H)^{-1}\right\|_{2}=\max _{\lambda_{j} \in \Lambda(H)}\left|\frac{\alpha-\lambda_{j}}{\alpha+\lambda_{j}}\right|$. Thus we obtain

$$
\rho(T(\alpha)) \leq \max _{\lambda_{j} \in \Lambda(H)}\left|\frac{\alpha-\lambda_{j}}{\alpha+\lambda_{j}}\right| \max _{i \tau_{j} \in \Lambda(S)}\left|\frac{\alpha-i \tau_{j}}{\alpha+i \tau_{j}}\right|
$$

Under the assumption that $\lambda_{j}>0$ nd $\operatorname{Re}(\alpha)>0, \operatorname{Im}(\alpha) \tau_{j} \geq 0$, we have

$$
\rho(T(\alpha)) \leq \omega(\alpha)<1
$$

i.e., the HSS iteration converges.

In particular, if the parameter $\alpha=a>0$ is chosen to be a real number (i.e., $b=0$ ), then when the matrix $A$ is positive definite, the HSS iteration (1.2) converges; and it holds that

$$
\rho(T(\alpha)) \leq \max _{\lambda_{j} \in \Lambda(H)}\left|\frac{a-\lambda_{j}}{a+\lambda_{j}}\right|<1
$$

This is the conclusion of Theorem 1.1.
The upper bound (2.1) shows that $\omega(\alpha)$ includes the spectral information of the matrix $S$, and if a suitable complex parameter $\alpha$ is chosen, it is possible that $\omega(\alpha)<\sigma(\alpha)$.

Remark 2.1. For a complex parameter $\alpha=a+i b$, the HS splitting can be viewed as the NS splitting with a real parameter $a$ :

$$
A=\alpha I+H-(\alpha I-S)=a I+(i b I+H)-(a I-(-i b I+S))
$$

and

$$
A=\alpha I+S-(\alpha I-H)=a I+(i b I+S)-(a I-(-i b I+H))
$$

where $H_{1} \equiv i b I+H$ and $H_{2} \equiv-i b I+H$ are normal matrices but $H_{1} \neq H_{2}$, and $S_{1} \equiv$ $-i b I+S, S_{2} \equiv i b I+S$ are skew-Hermitian matrices but $S_{1} \neq S_{2}$ too. Thus the HSS method with a complex parameter cannot be viewed as the NSS method with a real parameter [7].

Next we discuss the estimation to the optimal parameter of the upper bound $\omega(\alpha)$.
Let

$$
\omega_{1}(\alpha)=\max _{\lambda_{j} \in \Lambda(H)}\left|\frac{\alpha-\lambda_{j}}{\alpha+\lambda_{j}}\right|, \quad \omega_{2}(\alpha)=\max _{\tau_{j} \in \Lambda(S)}\left|\frac{\alpha-i \tau_{j}}{\alpha+i \tau_{j}}\right| .
$$

According to the condition for the convergence of HSS with a complex parameter in Theorem 2.1, we assume the eigenvalues $i \tau_{j}$ of the matrix $S$ satisfying

$$
\tau_{1} \geq \cdots \geq \tau_{n} \geq 0
$$

and choose the parameter $\alpha=a+i b$ such that $a>0$ and $b \geq 0$.

Lemma 2.1. Let $\lambda_{1} \geq \cdots \geq \lambda_{n}>0, \quad \tau_{1} \geq \cdots \geq \tau_{n} \geq 0$, and $a>0, \quad b \geq 0$. Then it holds that

$$
\begin{equation*}
\omega_{1}(\alpha)=\max \left(\frac{\left|a-\lambda_{1}+i b\right|}{\left|a+\lambda_{1}+i b\right|}, \frac{\left|a-\lambda_{n}+i b\right|}{\left|a+\lambda_{n}+i b\right|}\right) \tag{2.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\omega_{2}(\alpha)=\max \left(\frac{\left|a+i\left(b-\tau_{1}\right)\right|}{\left|a+i\left(b+\tau_{1}\right)\right|}, \frac{\left|a+i\left(b-\tau_{n}\right)\right|}{\left|a+i\left(b+\tau_{n}\right)\right|}\right) . \tag{2.3}
\end{equation*}
$$

Proof. We consider $\omega_{1}^{2}(\alpha)=\max _{\lambda_{j} \in \Lambda(H)} \frac{\left(a-\lambda_{j}\right)^{2}+b^{2}}{\left(a+\lambda_{j}\right)^{2}+b^{2}}$. Let

$$
g(\lambda) \equiv \frac{(a-\lambda)^{2}+b^{2}}{(a+\lambda)^{2}+b^{2}}, \quad 0<\lambda_{n} \leq \lambda \leq \lambda_{1}
$$

It is clear that $\lambda=\sqrt{a^{2}+b^{2}}$ is the unique positive minimizer of $g(\lambda)$. Therefore the maximum of $g(\lambda)$ reaches at the point $\lambda_{1}$ or $\lambda_{n}$. This proves the result (2.2). Similarly, the result (2.3) can be shown in the same way.

Lemma 2.2. Under the assumption of Lemma 2.1, for a fixed $b \geq 0$,
(1) if $b^{2} \leq \frac{\lambda_{n}}{2}\left(\lambda_{1}-\lambda_{n}\right)$, there is a minimizer of $\omega_{1}: a^{*} \equiv \operatorname{argmin}_{a>0} \omega_{1}(a+i b)=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}$, and it holds that

$$
\omega_{1}\left(a^{*}+i b\right)=\sqrt{\frac{\lambda_{1}+\lambda_{n}-2 a^{*}}{\lambda_{1}+\lambda_{n}+2 a^{*}}}=\sqrt{\frac{\lambda_{1}+\lambda_{n}-2 \sqrt{\lambda_{1} \lambda_{n}-b^{2}}}{\lambda_{1}+\lambda_{n}+2 \sqrt{\lambda_{1} \lambda_{n}-b^{2}}}}
$$

(2) if $b^{2}>\frac{\lambda_{n}}{2}\left(\lambda_{1}-\lambda_{n}\right)$, there is a minimizer of $\omega_{1}: a_{*} \equiv \operatorname{argmin}_{a>0} \omega_{1}(a+i b)=\sqrt{\lambda_{n}^{2}+b^{2}}$, and it holds that

$$
\omega_{1}\left(a_{*}+i b\right)=\sqrt{\frac{a_{*}-\lambda_{n}}{a_{*}+\lambda_{n}}}=\sqrt{\frac{\sqrt{\lambda_{n}^{2}+b^{2}}-\lambda_{n}}{\sqrt{\lambda_{n}^{2}+b^{2}}+\lambda_{n}}}
$$

Similarly, for a fixed $a>0$,
(3) if $a^{2} \leq \frac{\tau_{n}}{2}\left(\tau_{1}-\tau_{n}\right)$, there is a minimizer of $\omega_{2}: b^{*} \equiv \operatorname{argmin}_{b \geq 0} \omega_{2}(a+i b)=\sqrt{\tau_{1} \tau_{n}-a^{2}}$, and it holds that

$$
\omega_{2}\left(a+i b^{*}\right)=\sqrt{\frac{\tau_{1}+\tau_{n}-2 b^{*}}{\tau_{1}+\tau_{n}+2 b^{*}}}=\sqrt{\frac{\tau_{1}+\tau_{n}-2 \sqrt{\tau_{1} \tau_{n}-a^{2}}}{\tau_{1}+\tau_{n}+2 \sqrt{\tau_{1} \tau_{n}-a^{2}}}} ;
$$

(4) if $a^{2}>\frac{\tau_{n}}{2}\left(\tau_{1}-\tau_{n}\right)$, there is a minimizer of $\omega_{2}: b_{*} \equiv \operatorname{argmin}_{b \geq 0} \omega_{2}(a+i b)=\sqrt{\tau_{n}^{2}+a^{2}}$, and it holds that

$$
\omega_{2}\left(a+i b_{*}\right)=\sqrt{\frac{b_{*}-\tau_{n}}{b_{*}+\tau_{n}}}=\sqrt{\frac{\sqrt{\tau_{n}^{2}+a^{2}}-\tau_{n}}{\sqrt{\tau_{n}^{2}+a^{2}}+\tau_{n}}}
$$

Proof. We only prove the result (2.4). The conclusion (2.5) can be proven in the same way. We consider the minimum of $\omega_{1}^{2}(a+i b)$ for a fixed $b$. For $j=1$ and $j=n$, let

$$
\xi_{j}(a) \equiv \frac{\left(a-\lambda_{j}\right)^{2}+b^{2}}{\left(a+\lambda_{j}\right)^{2}+b^{2}}
$$

Then

$$
\xi_{j}^{\prime}(a)=\frac{4 \lambda_{j}\left(a^{2}-\lambda_{j}^{2}-b^{2}\right)}{\left[\left(a+\lambda_{j}\right)^{2}+b^{2}\right]^{2}}
$$

It is clare that $a=\sqrt{\lambda_{j}^{2}+b^{2}}$ is the unique positive minimizer of $\xi_{j}(a)$.
(i) In the case of $b^{2} \geq \lambda_{1} \lambda_{n} \geq \frac{\lambda_{n}}{2}\left(\lambda_{1}-\lambda_{n}\right)$, since for any $a>0$,

$$
\begin{equation*}
\xi_{1}(a)-\xi_{n}(a)=\frac{4 a}{\left(\left(a+\lambda_{1}\right)^{2}+b^{2}\right)\left(\left(a+\lambda_{n}\right)^{2}+b^{2}\right)}\left(\lambda_{1} \lambda_{n}-b^{2}-a^{2}\right)\left(\lambda_{1}-\lambda_{n}\right)<0 \tag{2.6}
\end{equation*}
$$

it holds that $\omega_{1}^{2}(a+i b)=\xi_{n}(a)$ for a fixed $b \geq \sqrt{\lambda_{1} \lambda_{n}}$. Therefore, the minimizer of $\omega_{1}(a+i b)$ with respect to $a$ is the minimizer of $\xi_{n}(a): a_{*}=\sqrt{\lambda_{n}^{2}+b^{2}}$.
(ii) In the case of $b^{2}<\lambda_{1} \lambda_{n}$, when $a>0, \xi_{1}$ and $\xi_{n}$ have the unique intersection point $a^{*}=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}$, and $a^{*}<\lambda_{1}$.

When $a^{*}=\sqrt{\lambda_{1} \lambda_{n}-b^{2}} \geq \sqrt{\lambda_{n}^{2}+b^{2}}$, or equivalently, $b^{2} \leq \frac{\lambda_{n}}{2}\left(\lambda_{1}-\lambda_{n}\right)$, for $a \in\left(\sqrt{\lambda_{n}^{2}+b^{2}}, \sqrt{\lambda_{1}^{2}+b^{2}}\right), \xi_{1}(a)$ decreases, while $\xi_{n}$ increases, so the minimum of $\omega_{1}(a+i b)$ with respect to $a$ is reached at the intersection point $a^{*}=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}$, which is the result of (2.4').

When $a^{*}=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}<\sqrt{\lambda_{n}^{2}+b^{2}}$, or equivalently, $b^{2}>\frac{\lambda_{n}}{2}\left(\lambda_{1}-\lambda_{n}\right)$, for $a \in\left(0, \sqrt{\lambda_{n}^{2}+b^{2}}\right)$, both $\xi_{1}$ and $\xi_{n}$ decrease, and from (2.6), we have

$$
\xi_{1}\left(a_{*}\right)-\xi_{n}\left(a_{*}\right)=\frac{4 a_{*}}{\left(\left(a_{*}+\lambda_{1}\right)^{2}+b^{2}\right)\left(\left(a_{*}+\lambda_{n}\right)^{2}+b^{2}\right)}\left(\lambda_{1} \lambda_{n}-b^{2}-\left(\lambda_{n}^{2}+b^{2}\right)\right)\left(\lambda_{1}-\lambda_{n}\right)<0
$$

so the minimum value of $\omega_{1}(a+i b)$ with respect to $a$ is reached at $a_{*}=\sqrt{\lambda_{n}^{2}+b^{2}}$ of $\xi_{n}$, which is the result $\left(2.4^{\prime \prime}\right)$.

According to the expression (2.1) of $\omega(a+b i)$, we consider the minimization of $\omega(a+i b)$ in the area $\Omega=\left[0<a \leq \lambda_{1}, 0 \leq b \leq \tau_{1}\right]$ :

$$
\begin{equation*}
\min _{(a, b) \in \Omega} \omega(a+i b)=\min _{(a, b) \in \Omega} \omega_{1}(a+i b) \omega_{2}(a+i b) \tag{2.7}
\end{equation*}
$$

Clearly, the following lemma holds.
Lemma 2.3. Let $f_{1}(a, b) \geq 0, f_{2}(a, b) \geq 0 \quad$ for $(a, b) \in \Omega$. Then for any non-negative function $b=s(a) \leq \tau_{1}$, it holds that

$$
\min _{(a, b) \in \Omega} f_{1}(a, b) f_{2}(a, b) \leq \min _{0<a \leq \lambda_{1}}\left[f_{1}(a, s(a)) \min _{0 \leq b \leq \tau_{1}} f_{2}(a, b)\right] .
$$

Similarly, for any positive function $a=t(b) \leq \lambda_{1}$, it holds that

$$
\min _{(a, b) \in \Omega} f_{1}(a, b) f_{2}(a, b) \leq \min _{0 \leq b \leq \tau_{1}}\left[f_{2}(t(b), b) \min _{0<a \leq \lambda_{1}} f_{1}(a, b)\right] .
$$

Set

$$
a_{0}=\min \left(\sqrt{\frac{\tau_{n}}{2}\left(\tau_{1}-\tau_{n}\right)}, \lambda_{1}\right), \quad b_{0}=\min \left(\sqrt{\frac{\lambda_{n}}{2}\left(\lambda_{1}-\lambda_{n}\right)}, \tau_{1}\right)
$$

According to Lemma 2.2, when $0<a<a_{0}, \omega_{2}(a, b)$ reaches its minimum at the point $b^{*}=$ $\sqrt{\tau_{1} \tau_{n}-a^{2}}$, so we set $s(a)=\sqrt{\tau_{1} \tau_{n}-a^{2}}$ in Lemma 2.3; when $a_{0}<a \leq \lambda_{1}, \omega_{2}(a, b)$ reaches its minimum at the point $b_{*}=\sqrt{\tau_{n}^{2}+a^{2}}$, so we set $s(a)=\sqrt{\tau_{1} \tau_{n}-a^{2}}$ in Lemma 2.3. Thus we have the following result (2.9). Similarly, by setting $t(b)=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}$ and $t(b)=\sqrt{\lambda_{n}^{2}+b^{2}}$ in Lemma 2.3, we can obtain the result (2.10). We summarize these in the following theorem.

Theorem 2.2. Under the assumption of Lemma 2.1, the following results hold true

$$
\begin{equation*}
\min _{(a, b) \in \Omega} \omega(a+i b) \leq \min \left(\min _{0<a \leq a_{0}} \phi_{1}(a), \min _{a_{0} \leq a \leq \lambda_{1}} \phi_{2}(a)\right) \tag{2.9}
\end{equation*}
$$

and

$$
\begin{equation*}
\min _{(a, b) \in \Omega} \omega(a+i b) \leq \min \left(\min _{0 \leq b \leq b_{0}} \psi_{1}(b), \min _{b_{0} \leq b \leq \tau_{1}} \psi_{2}(b)\right), \tag{2.10}
\end{equation*}
$$

where

$$
\begin{aligned}
& \phi_{1}(a)=\omega_{1}\left(a+i \sqrt{\tau_{1} \tau_{n}-a^{2}}\right) \omega_{2}\left(a+i \sqrt{\tau_{1} \tau_{n}-a^{2}}\right), \\
& \phi_{2}(a)=\omega_{1}\left(a+i \sqrt{\tau_{n}^{2}+a^{2}}\right) \omega_{2}\left(a+i \sqrt{\tau_{n}^{2}+a^{2}}\right), \\
& \psi_{1}(b)=\omega_{1}\left(\sqrt{\lambda_{1} \lambda_{n}-b^{2}}+i b\right) \omega_{2}\left(\sqrt{\lambda_{1} \lambda_{n}-b^{2}}+i b\right), \\
& \psi_{2}(b)=\omega_{1}\left(\sqrt{\lambda_{n}^{2}+b^{2}}+i b\right) \omega_{2}\left(\sqrt{\lambda_{n}^{2}+b^{2}}+i b\right) .
\end{aligned}
$$

Next, we discuss the problem $\min _{0 \leq b \leq b_{0}} \psi_{1}(b)$ by considering the minimization of $\psi_{1}^{2}(b)$. The problem $\min _{0<a \leq a_{0}} \phi_{1}(a)$ can be discussed in the same way.

From (2.4') in Lemma 2.2 and by setting $a^{*}(b)=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}$, we have

$$
\omega_{1}^{2}\left(a^{*}(b)+i b\right)=1-\frac{4 a^{*}(b)}{\lambda_{1}+\lambda_{n}+2 a^{*}(b)}
$$

and

$$
\omega_{2}^{2}\left(a^{*}(b)+i b\right)=\max \left(\frac{\left(a^{*}(b)\right)^{2}+\left(b-\tau_{1}\right)^{2}}{\left(a^{*}(b)\right)^{2}+\left(b+\tau_{1}\right)^{2}}, \frac{\left(a^{*}(b)\right)^{2}+\left(b-\tau_{n}\right)^{2}}{\left(a^{*}(b)\right)^{2}+\left(b+\tau_{n}\right)^{2}}\right) .
$$

For $j=1$ and $j=n$, let

$$
h_{j}(b) \equiv \frac{\left(a^{*}(b)\right)^{2}+\left(b-\tau_{j}\right)^{2}}{\left(a^{*}(b)\right)^{2}+\left(b+\tau_{j}\right)^{2}}=1-\frac{4 b \tau_{j}}{\tau_{j}^{2}+\lambda_{1} \lambda_{n}+2 b \tau_{j}} .
$$

Then

$$
h_{1}(b)-h_{n}(b)=\frac{4 b\left(\tau_{1}-\tau_{n}\right)\left(\tau_{1} \tau_{n}-\lambda_{1} \lambda_{n}\right)}{\left(\lambda_{1} \lambda_{n}+\tau_{1}^{2}+2 b \tau_{1}\right)\left(\lambda_{1} \lambda_{n}+\tau_{n}^{2}+2 b \tau_{n}\right)} .
$$

If $\tau_{1} \tau_{n} \geq \lambda_{1} \lambda_{n}$, then $h_{1}(b) \geq h_{n}(b)$. Therefore,

$$
\omega_{2}^{2}\left(a^{*}(b)+i b\right)=h_{1}(b)=1-\frac{4 b \tau_{1}}{\tau_{1}^{2}+\lambda_{1} \lambda_{n}+2 b \tau_{1}}
$$

If $\tau_{1} \tau_{n} \leq \lambda_{1} \lambda_{n}$, then $h_{1}(b) \leq h_{n}(b)$. Therefore,

$$
\omega_{2}^{2}\left(a^{*}(b)+i b\right)=h_{n}(b)=1-\frac{4 b \tau_{n}}{\tau_{n}^{2}+\lambda_{1} \lambda_{n}+2 b \tau_{n}} .
$$

Thus, $\psi_{1}^{2}(b)$ has the following expression in $\left[0, b_{0}\right]$ :

$$
\psi_{1}^{2}(b)= \begin{cases}\left(1-\frac{4 a^{*}(b)}{\lambda_{1}+\lambda_{n}+2 a^{*}(b)}\right)\left(1-\frac{4 b \tau_{1}}{\tau_{1}^{2}+\lambda_{1} \lambda_{n}+2 b \tau_{1}}\right), & \tau_{1} \tau_{n} \geq \lambda_{1} \lambda_{n}, \\ \left(1-\frac{4 a^{*}(b)}{\lambda_{1}+\lambda_{n}+2 a^{*}(b)}\right)\left(1-\frac{4 \tau_{n}}{\tau_{n}^{2}+\lambda_{1} \lambda_{n}+2 b \tau_{n}}\right), & \tau_{1} \tau_{n} \leq \lambda_{1} \lambda_{n} .\end{cases}
$$

Similarly, $\phi_{1}^{2}(a)$ has the following expression in ( $0, a_{0}$ ]:

$$
\phi_{1}^{2}(a)= \begin{cases}\left(1-\frac{4 b^{*}(a)}{\tau_{1}+\tau_{n}+2 b^{*}(a)}\right)\left(1-\frac{4 a \lambda_{1}}{\lambda_{1}^{2}+\tau_{1} \tau_{n}+2 a \lambda_{1}}\right), & \tau_{1} \tau_{n} \leq \lambda_{1} \lambda_{n}, \\ \left(1-\frac{4 b^{*}(a)}{\tau_{1}+\tau_{n}+2 b^{*}(a)}\right)\left(1-\frac{4 a \lambda_{n}}{\lambda_{n}^{2}+\tau_{1} \tau_{n}+2 a \lambda_{n}}\right), & \tau_{1} \tau_{n} \geq \lambda_{1} \lambda_{n},\end{cases}
$$

where $b^{*}(a)=\sqrt{\tau_{1} \tau_{n}-a^{2}}$.
The following theorem provides the minima of $\phi_{1}(a)$ and $\psi_{1}(b)$.

Theorem 2.3. Under the assumption of Lemma 2.1 and $\tau_{1} \neq \tau_{n}$, there is the minimizer $a_{\text {est }}^{\prime}=$ $\sqrt{\tilde{a}^{*}}$ of $\phi_{1}(a)$ in $\left(0, \sqrt{\tau_{1} \tau_{n}}\right]$, and $\tilde{a}^{*}$ is a positive root of either

$$
p_{1}(\tilde{a})=0, \quad \lambda_{1} \lambda_{n} \geq \tau_{1} \tau_{n}
$$

or

$$
p_{n}(\tilde{a})=0, \quad \lambda_{1} \lambda_{n} \leq \tau_{1} \tau_{n}
$$

where for $j=1$ and $j=n$,

$$
\begin{aligned}
p_{j}(\tilde{a}) \equiv 16 & \lambda_{j}^{2}\left(\lambda_{j}^{2} \tau_{c}^{2}+u_{j}^{2}\right) \tilde{a}^{3}-48 \tau_{1} \tau_{n} \lambda_{j}^{2} u_{j}^{2} \tilde{a}^{2} \\
& +u_{j}^{2}\left[\tau_{c}^{2} u_{j}^{2}+\lambda_{j}^{2}\left(\tau_{1}-\tau_{n}\right)^{2}\left(\tau_{1}^{2}+\tau_{n}^{2}-10 \tau_{1} \tau_{n}\right)\right] \tilde{a}-\lambda_{j}^{2} u_{j}^{2} \tau_{1} \tau_{n}\left(\tau_{1}-\tau_{n}\right)^{4}
\end{aligned}
$$

$u_{j}=\lambda_{j}^{2}+\tau_{1} \tau_{n}, \tau_{c}=\tau_{1}+\tau_{n}, \tilde{a}=a^{2}$.
Similarly, there is a minimizer $b_{\text {est }}^{\prime \prime}=\sqrt{\tilde{b}^{*}}$ of $\psi_{1}(b)$ in $\left[0, \sqrt{\lambda_{1} \lambda_{n}}\right]$, and $\tilde{b}^{*}$ is a non-negative root of either

$$
q_{1}(\tilde{b})=0, \quad \tau_{1} \tau_{n} \geq \lambda_{1} \lambda_{n}
$$

or

$$
q_{n}(\tilde{b})=0, \quad \tau_{1} \tau_{n} \leq \lambda_{1} \lambda_{n}
$$

where for $j=1$ and $j=n$,

$$
\begin{aligned}
q_{j}(\tilde{b}) & \equiv 16 \tau_{j}^{2}\left(\tau_{j}^{2} \lambda_{c}^{2}+v_{j}^{2}\right) \tilde{b}^{3}-48 \lambda_{1} \lambda_{n} \tau_{j}^{2} v_{j}^{2} \tilde{b}^{2} \\
& +v_{j}^{2}\left[\lambda_{c}^{2} v_{j}^{2}+\tau_{j}^{2}\left(\lambda_{1}-\lambda_{n}\right)^{2}\left(\lambda_{1}^{2}+\lambda_{n}^{2}-10 \lambda_{1} \lambda_{n}\right)\right] \tilde{b}-\tau_{j}^{2} v_{j}^{2} \lambda_{1} \lambda_{n}\left(\lambda_{1}-\lambda_{n}\right)^{4}
\end{aligned}
$$

$v_{j}=\tau_{j}^{2}+\lambda_{1} \lambda_{n}, \lambda_{c}=\lambda_{1}+\lambda_{n}, \tilde{b}=b^{2}$.
Proof. We prove the result $\left(2.12^{\prime}\right)$. The results $\left(2.12^{\prime \prime}\right)$ and (2.11) can be shown in the same way.

For simplicity, we now omit the superscript of $a^{*}(b)$ and let $a(b)=\sqrt{\lambda_{1} \lambda_{n}-b^{2}}$. Then

$$
\psi_{1}^{2}(b)=\frac{\lambda_{c}-2 a(b)}{\lambda_{c}+2 a(b)} \frac{v_{1}-2 b \tau_{1}}{v_{1}+2 b \tau_{1}}
$$

By setting $\frac{d \psi_{1}^{2}}{d b}=0$, we have

$$
\begin{equation*}
q(b) \equiv 4 \lambda_{c} \tau_{1}^{2} a^{\prime}(b) b^{2}-\lambda_{c} v_{1}^{2} a^{\prime}(b)+4 \tau_{1} v_{1} a(b)^{2}-\tau_{1} v_{1} \lambda_{c}^{2}=0 \tag{2.13}
\end{equation*}
$$

Consider the value of $q(b)$ at $b_{\epsilon}=\sqrt{\lambda_{1} \lambda_{n}-\epsilon^{2}}$, where $\epsilon>0$ is a small positive number. Note that $a\left(b_{\epsilon}\right)=\epsilon>0$ and $a^{\prime}\left(b_{\epsilon}\right)=-\frac{\sqrt{\lambda_{1} \lambda_{n}-\epsilon^{2}}}{\epsilon}$. Thus we have

$$
\begin{aligned}
q\left(b_{\epsilon}\right) & =4 \lambda_{c} \tau_{1}^{2}\left(\lambda_{1} \lambda_{n}-\epsilon^{2}\right)\left(-\frac{\sqrt{\lambda_{1} \lambda_{n}-\epsilon^{2}}}{\epsilon}\right)-\lambda_{c} v_{1}^{2}\left(-\frac{\sqrt{\lambda_{1} \lambda_{n}-\epsilon^{2}}}{\epsilon}\right)+4 \tau_{1} v_{1} \epsilon^{2}-\tau_{1} v_{1} \lambda_{c}^{2} \\
& =\frac{1}{\epsilon}\left[\lambda_{c} \sqrt{\lambda_{1} \lambda_{n}-\epsilon^{2}}\left(v_{1}^{2}-4 \tau_{1}^{2} \lambda_{1} \lambda_{n}\right)+\mathcal{O}(\epsilon)\right] \\
& =\frac{1}{\epsilon}\left[\lambda_{c} \sqrt{\lambda_{1} \lambda_{n}-\epsilon^{2}}\left(\tau_{1}^{2}-\lambda_{1} \lambda_{n}\right)^{2}+\mathcal{O}(\epsilon)\right] .
\end{aligned}
$$

For a suitable small $\epsilon>0, q\left(b_{\epsilon}\right) \geq 0$. Also for a suitable small $\epsilon>0$,

$$
q(\epsilon)=-\tau_{1} v_{1}\left(\lambda_{1}-\lambda_{n}\right)^{2}+\mathcal{O}(\epsilon) \leq 0
$$

This shows that as $b$ increases in $\left[0, \sqrt{\lambda_{1} \lambda_{n}}\right], \psi_{1}^{2}(b)$ firstly decreases and then increases, and therefore, $\psi_{1}^{2}(b)$ has the minimizer $b_{e s t}^{\prime \prime}$ in $\left[0, \sqrt{\lambda_{1} \lambda_{n}}\right]$, which satisfies $q\left(b_{e s t}^{\prime \prime}\right)=0$. If $\lambda_{1} \neq \lambda_{n}$, $b_{e s t}^{\prime \prime} \in\left(0, \sqrt{\lambda_{1} \lambda_{n}}\right]$.

Next, we simplify the expression (2.13). By substituting $a(b)^{2}=\lambda_{1} \lambda_{n}-b^{2}$ and $a^{\prime}(b)=$ $-\frac{b}{\sqrt{\lambda_{1} \lambda_{n}-b^{2}}}$ into (2.13), we have

$$
-4 \lambda_{c} \tau_{1}^{2} \frac{b^{3}}{\sqrt{\lambda_{1} \lambda_{n}-b^{2}}}+\lambda_{c} v_{1}^{2} \frac{b}{\sqrt{\lambda_{1} \lambda_{n}-b^{2}}}-4 \tau_{1} v_{1} b^{2}+\tau_{1} v_{1}\left(4 \lambda_{1} \lambda_{n}-\lambda_{c}^{2}\right)=0
$$

or equivalently,

$$
\lambda_{c}^{2}\left(v_{1}^{2}-4 \tau_{1}^{2} b^{2}\right)^{2} b^{2}=\tau_{1}^{2} v_{1}^{2}\left[4 b^{2}+\left(\lambda_{1}-\lambda_{n}\right)^{2}\right]^{2}\left(\lambda_{1} \lambda_{n}-b^{2}\right)
$$

Let $\tilde{b}=b^{2}$. Then we have

$$
\begin{aligned}
& \left(16 \lambda_{c}^{2} \tau_{1}^{4}+16 \tau_{1}^{2} v_{1}^{2}\right) \tilde{b}^{3}+\left[-8 \lambda_{c}^{2} \tau_{1}^{2} v_{1}^{2}+8 \tau_{1}^{2} v_{1}^{2}\left(\lambda_{1}-\lambda_{n}\right)^{2}-16 \tau_{1}^{2} v_{1}^{2} \lambda_{1} \lambda_{n}\right] \tilde{b}^{2} \\
& \quad+\left[\lambda_{c}^{2} v_{1}^{4}-8 \tau_{1}^{2} v_{1}^{2}\left(\lambda_{1}-\lambda_{n}\right)^{2} \lambda_{1} \lambda_{n}+\tau_{1}^{2} v_{1}^{2}\left(\lambda_{1}-\lambda_{n}\right)^{4}\right] \tilde{b}-\tau_{1}^{2} v_{1}^{2}\left(\lambda_{1}-\lambda_{n}\right)^{4} \lambda_{1} \lambda_{n}=0
\end{aligned}
$$

which leads to $\left(2.12^{\prime}\right)$.
The minimization problems of $\phi_{2}(a)$ and $\psi_{2}(b)$ can be discussed in a similar way, and hence are omitted. It is worthy mentioned that our numerical experiments indicate that the minimum value of $\phi_{1}(a)$ or $\psi_{1}(b)$ is smaller than the minimum value of $\phi_{2}(a)$ or $\psi_{2}(b)$ in (2.9) and (2.10).

Consequently by Theorem 2.2, we have the following upper bound for $\min \omega(a+i b)$,

$$
\begin{equation*}
\min _{0<a \leq \lambda_{1}, 0 \leq b \leq \tau_{1}} \omega(a+i b) \leq \min \left(\min _{0<a \leq \sqrt{\tau_{1} \tau_{2}}} \phi_{1}(a), \min _{0 \leq b \leq \sqrt{\lambda_{1} \lambda_{2}}} \psi_{1}(b)\right) \tag{2.14}
\end{equation*}
$$

This minimization problem, by Theorem 2.3, can be solved via solving equations of (2.11) and (2.12).

We now summarize our discussions and provide our parameter $\alpha_{\text {est }}$ which estimates the optimal parameter of the upper bound $\omega(\alpha)$ as follows.

Given $\lambda_{1} \geq \lambda_{n}>0$ and $\tau_{1}>\tau_{n} \geq 0$,
(1) for the case $\lambda_{1} \lambda_{n} \geq \tau_{1} \tau_{n}$ :

- find the minimizer $a_{\text {est }}^{\prime}$ of $\phi_{1}(a)$ by solving the positive root $p_{1}(\tilde{a})=0$ in (2.11'), and the minimizer $b_{\text {est }}^{\prime \prime}$ of $\psi_{1}(b)$ by solving the positive root $q_{n}(\tilde{b})=0$ in $\left(2.12^{\prime \prime}\right)$;
- if $\phi_{1}\left(a_{e s t}^{\prime}\right) \leq \psi_{1}\left(b_{e s t}^{\prime \prime}\right)$, let $a_{e s t}=a_{e s t}^{\prime}, b_{e s t}=\sqrt{\tau_{1} \tau_{n}-a_{e s t}^{2}}$ to form $\alpha_{e s t}=a_{e s t}+i b_{e s t}$; otherwise
- if $\phi_{1}\left(a_{\text {est }}^{\prime}\right)>\psi_{1}\left(b_{e s t}^{\prime \prime}\right)$, let $b_{\text {est }}=b_{\text {est }}^{\prime \prime}, a_{\text {est }}=\sqrt{\lambda_{1} \lambda_{n}-b_{\text {est }}^{2}}$ to form $\alpha_{\text {est }}=a_{\text {est }}+i b_{\text {est }}$.
(2) for the case $\lambda_{1} \lambda_{n}<\tau_{1} \tau_{n}$ :
- find the minimizer $a_{\text {est }}^{\prime}$ of $\phi_{1}(a)$ by solving the positive root $p_{n}(\tilde{a})=0$ in $\left(2.11^{\prime \prime}\right)$, and the minimizer $b_{\text {est }}^{\prime \prime}$ of $\psi_{1}(b)$ by solving the positive root $q_{1}(\tilde{b})=0$ in $\left(2.12^{\prime}\right)$;
- if $\phi_{1}\left(a_{e s t}^{\prime}\right) \leq \psi_{1}\left(b_{e s t}^{\prime \prime}\right)$, let $a_{e s t}=a_{e s t}^{\prime}, b_{e s t}=\sqrt{\tau_{1} \tau_{n}-a_{e s t}^{2}}$ to form $\alpha_{e s t}=a_{e s t}+i b_{e s t}$; otherwise
- if $\phi_{1}\left(a_{\text {est }}^{\prime}\right)>\psi_{1}\left(b_{e s t}^{\prime \prime}\right)$, let $b_{\text {est }}=b_{\text {est }}^{\prime \prime}, a_{\text {est }}=\sqrt{\lambda_{1} \lambda_{n}-b_{\text {est }}^{2}}$ to form $\alpha_{\text {est }}=a_{\text {est }}+i b_{\text {est }}$.


## 3. Implementation Aspects

In the HSS method, two shifted sub-systems with respect to $\alpha I+H$ and $\alpha I+S$ must be solved in each iteration.

For a real parameter $\alpha>0$ and a Hermitian positive definite matrix $H$, the shifted matrix $\alpha I+H$ remains a Hermitian positive definite, so the 'exact' solution could be solved by the (complex) Cholesky factorization method, while its 'inexact' solution could be solved by the (complex) CG method; however, for the real shifted skew-Hermitian matrix $\alpha I+S$, when $S$ is structured [3], the difficulty in solving the relevant system varies from case to case. But in general, this will not be the case.

For a complex parameter $\alpha=a+i b$, since the sub-system $(\alpha I+S) u=g$ is equivalent to $(i \alpha I+i S) u=i g$, both shifted sub-systems become systems with respect to a complex shifted Hermitian positive definite matrix $\alpha I+H$ or $i \alpha I+i S$ under the assumption of Lemma 2.1. For large sparse matrices, these shifted sub-systems could be solved by certain shifted Krylov subspace method like the Lanczos method, or the MINRES method et al. [13-16,18-20]. Since the Krylov method keeps shift invariance, the basis vector can be constructed by the matrix $H$ or $i S$ (using short recurrence), irrelative to the shift $\alpha$ or $i \alpha$. Thus the computational cost with a complex parameter could not be significantly higher than that with a real parameter. This deserves a further and detailed investigation.

The HSS preconditioner is nowadays mostly discussed. The matrix

$$
P=\frac{1}{2 \alpha}(H+\alpha I)(S+\alpha I)
$$

is called the HSS preconditioner, where $\alpha$ is referred to as the preconditioning parameter. It is known that in general, the smaller the spectral radius $\rho(T(\alpha))$ of the HSS iterative matrix is, the better the gathering of the spectrum the preconditioned matrix $P^{-1} A$. This will make the Krylov subspace method converge faster; see Experiment 2 in the next section. Also, two shifted systems with respect to $H+\alpha I$ and $S+\alpha I$ (or $i S+i \alpha I$ ) have to be solved in solving the preconditioning sub-system each iteration. We can still use the shifted Krylov subspace method to solve them as a polynomial preconditioner. In addition, the incomplete factorization [14] of these shifted matrix may be used as a proconditioner, instead of $H+\alpha I$ and $S+\alpha I$, which could probably keep the computational cost economic, even for the complex parameter.

How to solve the systems with respect to a shifted Hermitian or skew-Hermitian efficiently is a further problem. We will leave it for our further work.

## 4. Numerical Experiments

In this section, we present two numerical experiments in solving the complex linear systems (1.1) from two modified model problems by the HSS iteration with our estimated parameter $\alpha_{\text {est }}$ and with the real parameter $\tilde{\alpha}$ in (1.5); moreover, for comparison purpose, the HSS method with the experimental 'optimal' parameter $\alpha_{\text {exp }}=a_{\text {exp }}+i b_{\text {exp }}$ is also tested, which is obtained by computing the spectral radius $\rho(T(\alpha))$ of the iterative matrix $T(\alpha)$ with all parameter $\alpha$ on $100 \times 100$ mesh points in $\left[0, \lambda_{1} ; 0, \tau_{1}\right]$. In Experiment 2, we also report results from full GMRES method, full GMRES preconditioned by $\operatorname{HSS}(\alpha)$ preconditioner (see Table 4.5), and the sensitivity of the estimated parameter $\alpha_{\text {est }}$ and $\rho\left(T\left(\alpha_{\text {est }}\right)\right)$ with respect to the spectral information of $H$ and $S$ (see Table 4.6).

In our test, the eigenvalues of a matrix are solved by the function eig and the root of $\pi(x)=0$ is solved by the function root in Matlab( 7.4 ed ).

We report the results of our numerical experiments with a Fortran 77 implementation of the HSS method based on the iteration (1.2). Two shifted sub-systems are solved by exact factorization method. The right-hand side of the linear system (1.1) is formed by $f=A x$, where

$$
\begin{equation*}
x=(1-i, 1-i, \cdots, 1-i)^{T} \tag{4.1}
\end{equation*}
$$

The initial value is $x^{(0)}=0$, and the stopping criterion is based on the residual of system (1.1) $\left\|f-A x^{(k)}\right\|<10^{-6}$.

Experiment 1. Consider a differential equation with complex coefficient which forms a nonHermitian complex matrix,

$$
-u_{x x}+i \delta x u_{x}+\gamma u=g, \quad x \in[0,1],
$$

where $\delta \in R, \gamma \in C, i=\sqrt{-1}$. This equation is modified to the model problem

$$
-u_{x x}+\delta x u_{x}+\gamma u=g, \quad \delta, \gamma \in R .
$$

When the centered difference to $u_{x x}$ and centered differences or the forward difference to $u_{x}$ are applied to the above model, we get the linear system (1.1) with the tridiagonal complex coefficient matrices, denoted by $A_{c}$ and $A_{f}$ respectively. The $k$ th rows of $A_{c}$ and $A_{f}$ are:

$$
\left(A_{c}\right)_{k}=\operatorname{tridiag}\left(-1-i k P, 2+\gamma h^{2},-1+i k P\right)
$$

or

$$
\left(A_{f}\right)_{k}=\operatorname{tridiag}\left(-1,2+(\gamma-i k \delta) h^{2},-1+i k \delta h^{2}\right)
$$

where $h=\frac{1}{n+1}$ is the step-size and $P=\frac{\delta h^{2}}{2}$. In our numerical tests, we set $n=100$ (since we want to seek the experimental 'optimal' parameter $\alpha_{\text {exp }}$ on $100 \times 100$ mesh points, it will take too much time if $n$ is large) and $\gamma=2000+i 20000$. The numerical results with $A_{c}$ and $A_{f}$ are listed in Table 4.1 and Table 4.2 respectively.

In these tables, the experimental 'optimal' parameter $\alpha_{\text {exp }}$, the estimated parameter $\alpha_{\text {est }}$, the upper-bound minimizer $\tilde{\alpha}$ for real parameter, and the corresponding spectral radii $\rho(\alpha) \equiv$ $\rho(T(\alpha))$ of the HSS iteration matrix $T(\alpha)$ are presented; moreover the number of iterations (denoted by IT) for the convergence of the HSS iteration method are also listed with different parameter $\delta$.

Table 4.1: $\alpha, \rho(\alpha)$ and iteration number(IT) of HSS with $A_{c}$ for Ex. 1.

| $\delta$ | 5 | 10 | 100 |
| :---: | :---: | :---: | :---: |
| $\alpha_{\text {exp }}$ | $0.0000+1.9608 i$ | $0.0000+1.9611 i$ | $0.0000+1.9655 i$ |
| $\rho\left(\alpha_{\text {exp }}\right)$ | $1.1549 \times 10^{-4}$ | $2.5441 \times 10^{-4}$ | $2.4646 \times 10^{-3}$ |
| IT | 3 | 3 | 4 |
| $\alpha_{\text {est }}$ | $6.3362 \times 10^{-9}+1.9605 i$ | $2.5280 \times 10^{-8}+1.9606 i$ | $3.3174 \times 10^{-7}+1.9606 i$ |
| $\rho\left(\alpha_{\text {est }}\right)$ | $8.5953 \times 10^{-5}$ | $1.2693 \times 10^{-4}$ | $1.2431 \times 10^{-3}$ |
| IT | 3 | 3 | 4 |
| $\tilde{\alpha}$ | 0.9087 | 0.9075 | 0.3406 |
| $\rho(\tilde{\alpha})$ | 0.6439 | 0.6443 | 0.8553 |
| IT | 39 | 39 | $\times$ |

From Table 4.1 we see that the numerical results produced with $\alpha_{\text {est }}$ coincide with those using $\alpha_{\text {exp }}$, but it does not for $\tilde{\alpha}$.

Remark 4.1. (1) The situation $\rho\left(\alpha_{\text {exp }}\right)>\rho\left(\alpha_{\text {est }}\right)$ occurs. This is because the mesh $100 \times 100$ over $\left[0, \lambda_{1} ; 0, \tau_{1}\right]$ is rough, e.g., for $\delta=5, \lambda_{1}=4.1953, \tau_{1}=1.9608$, and the mesh size is $0.04195 \times 0.01961$, which cannot reach the precision $\mathcal{O}\left(10^{-3}\right)$. Therefore it is possible that $\alpha_{\exp }$ is not an exact optimal parameter.
(2) The symbol ' $\times$ ' in Table 4.1 (also in Table 4.2) means that the HSS method could not meet the stopping criterion after 200 iteration steps.

Table 4.2: $\alpha, \rho(\alpha)$ and iteration number(IT) of HSS with $A_{f}$ for Ex. 1.

| $\delta$ | 5 | 10 | 100 |
| :---: | :---: | :---: | :---: |
| $\alpha_{\text {exp }}$ | $0.0000+1.9211 i$ | $0.0000+1.8817 i$ | $1.9200+1.9556 i$ |
| $\rho\left(\alpha_{\text {exp }}\right)$ | 0.0139 | 0.0288 | 0.4373 |
| IT | 5 | 6 | 19 |
| $\alpha_{\text {est }}$ | $0.2246 \times 10^{-3}+1.9138 i$ | $0.9409 \times 10^{-3}+1.8659 i$ | $0.0957+0.4701 i$ |
| $\rho\left(\alpha_{\text {est }}\right)$ | 0.0120 | 0.0246 | 0.5855 |
| IT | 5 | 6 | 49 |
| $\tilde{\alpha}$ | 0.9087 | 0.9075 | 0.3528 |
| $\rho(\tilde{\alpha})$ | 0.6439 | 0.6444 | 0.8501 |
| IT | 39 | 39 | $\times$ |

The results in Table 4.2 indicate that that $\alpha_{e s t}$ is a good approximation to $\alpha_{\text {exp }}$ for small $\delta$, and the iteration number for convergence with $\alpha_{e s t}$ is the same as that with $\alpha_{\text {exp }}$. On the other hand for large $\delta, \alpha_{\text {est }}$ is not close to $\alpha_{\text {exp }}$, but is still better than $\tilde{\alpha}$.

Experiment 2. (See [1,3]) The linear systems (1.1) is of the form

$$
\begin{equation*}
(W+i Z) x=f \tag{4.2}
\end{equation*}
$$

where $W=\tilde{K}+w_{1} I, \quad Z=\tilde{K}+w_{2} I$ and $w_{1}=\frac{3+\sqrt{3}}{\tau}, w_{2}=\frac{3-\sqrt{3}}{\tau}, \quad \tau$ is the time step-size and $\tilde{K}$ is the five-point centered difference matrix approximating the negative Laplacian operator $L=-\Delta$ with homogeneous Dirichlet boundary conditions, on a uniform mesh in the unit square $[0,1] \times[0,1]$ with the mesh-size $h=\frac{1}{m+1}$. Thus, the systems is a complex symmetric system with the $n \times n$ coefficient matrix $A=W+i Z$ and $n=m \times m$. For the complex symmetric systems, the modified HSS method [3] has the considerable advantage.

In our experiment, in order to show the advantage of the HSS with a complex parameter, we modify the above so that the resulting system is a complex non-symmetric system, but satisfying with the convergence condition of Theorem 2.1.

First, the matrix $\tilde{K}$ is changed to $K$, which is the five-point centered difference matrix approximating the operator $L=-\Delta+\gamma\left(\partial_{x}+\partial_{y}\right), \gamma \in R$. Let

$$
\begin{equation*}
W=K+w_{1} I \quad \text { and } \quad Z=K+w_{2} I \tag{4.3}
\end{equation*}
$$

Thus, the matrix $A=W+i Z$ is complex non-symmetric. In our tests, we also take $\tau=h$ and normalize the coefficient matrix and the right-hand side by multiplying both by $h^{2}$; see [3].

Let $H=\frac{1}{2}\left(A+A^{H}\right)$ and $S=\frac{1}{2}\left(A-A^{H}\right)$. We use the function eig in Matlab to solve the eigenvalues of the matrices $H$ and $S: \lambda_{1}=\lambda_{\max }(H)=8.2119, \lambda_{n}=\lambda_{\min }(H)=0.3448 ; \tau_{1}=$
$\tau_{\max }(S)=8.0082, \tau_{n}=\tau_{\min }(S)=0.1410$. Our estimated parameter $\alpha_{\text {est }}$ can be derived by (2.11), (2.12). The experimental result for Ex. 2 with (4.2), (4.3) and $m=16, \gamma=1$ is listed in Table 4.3. Fig. 1 shows the spectral radii $\rho(\alpha)$ with all $\alpha=a+i b$ on the $100 \times 100$ mesh points in $\left[0, \lambda_{1} ; 0, \tau_{1}\right]$.


Fig. 4.1. $\rho(\alpha)$ with all $\alpha=a+i b$ on $100 \times 100$ mesh points in $\left[0, \lambda_{1} ; 0, \tau_{1}\right]$ for Ex. 2 with (4.2), (4.3) and $m=16, \gamma=1$.

Table 4.3: $\alpha, \rho(\alpha), \omega(\alpha)$ or $\sigma(\alpha)$ and iteration number(IT) of HSS for Ex. 2 with (4.2), (4.3) and $m=16, \gamma=1$.

| parameter | $\alpha_{\text {est }}=1.5799+0.5792 i$ | $\tilde{\alpha}=1.6827$ | $\alpha_{\text {exp }}=1.3139+0.7207 i$ |
| :---: | :---: | :---: | :---: |
| spectral radius | $\rho\left(\alpha_{\text {est }}\right)=0.6375$ | $\rho(\tilde{\alpha})=0.6598$ | $\rho\left(\alpha_{\text {exp }}\right)=0.6089$ |
| upper bound | $\omega\left(\alpha_{\text {est }}\right)=0.6409$ | $\sigma(\tilde{\alpha})=0.6599$ | - |
| IT | 37 | 39 | 33 |

From Table 4.3 we see that $\alpha_{\text {est }}$ yields a better approximation to $\alpha_{\text {exp }}$ than $\tilde{\alpha}$ does; however $\alpha_{\text {est }}$ makes improvement on the spectral radius a little bit on $\tilde{\alpha}$. We note that $\tilde{\alpha}$ is close to the real part of $\alpha_{\exp }\left(\right.$ or $\alpha_{e s t}$ ) and this real part is 'dominant' to the imaginary part.

Now we exchange $w_{1}$ with $w_{2}$ in (4.3), such that

$$
\begin{equation*}
W=K+w_{2} I \quad \text { and } \quad Z=K+w_{1} I . \tag{4.4}
\end{equation*}
$$

Note that $w_{1}>w_{2}$, which means that the imaginary part $Z$ of the matrix $A$ is 'dominant' to the real part $W$. In this case, the eigenvalues of the matrices $H$ and $S$ are $\lambda_{1}=8.0082, \lambda_{n}=$ $0.1410 ; \tau_{1}=8.2119, \tau_{n}=0.3448$. The experimental result for Ex. 2 with (4.2), (4.4) and $m=16, \gamma=1$ is listed in Table 4.4.

Fig. 2 shows that the optimal parameter is complex indeed, and $\rho(\alpha)$ can not reach its minimum on the real axis.

The results in Table 4.4 show that $\alpha_{\text {est }}$ yields a good approximation to $\alpha_{\text {exp }}$, and a considerable improvement on the spectral radius, as well as the iteration number for convergence on the parameter $\tilde{\alpha}$. Also, We note that the imaginary part of $\alpha_{\text {exp }}$ is 'dominant' to the real part.


Fig. 4.2. $\rho(\alpha)$ with all $\alpha=a+i b$ on $100 \times 100$ mesh points in $\left[0, \lambda_{1} ; 0, \tau_{1}\right]$ for Ex. 2 with (4.2), (4.4) and $m=16, \gamma=1$.

Next we make the imaginary part $Z$ more 'dominant' to the real part $W$ by setting

$$
\begin{equation*}
W=K+\frac{w_{2}}{2} I \quad \text { and } \quad Z=K+2 w_{1} I . \tag{4.5}
\end{equation*}
$$

The eigenvalues of the matrices $H$ and $S$ are $\lambda_{1}=7.9709, \lambda_{n}=0.1037 ; \tau_{1}=8.4903, \tau_{n}=$ 0.6231. The experimental results for Ex. 2 with (4.2), (4.5) and $m=16, \gamma=1$ are also shown in Table 4.4. Our estimated parameter $\alpha_{\text {est }}$ makes more improvement.

Other mesh-size $m$ and the parameter $\gamma$ are also tested; see the following Table 4.5. However we do not seek $\alpha_{\text {exp }}$, since it will take much time for large $m$. Our estimated parameter $\alpha_{\text {est }}$ makes the HSS iteration a considerable improvement on the parameter $\tilde{\alpha}$.

Table 4.4: $\alpha, \rho(\alpha), \omega(\alpha)$ or $\sigma(\alpha)$ and iteration number(IT) of HSS for Ex. 2 with (4.2), (4.4), (4.5) and $m=16, \gamma=1$.

| Ex.2 with (4.2), (4.4) |  |  |  |
| :---: | :---: | :---: | :---: |
| parameter | $\alpha_{\text {est }}=0.5792+1.5799 i$ | $\tilde{\alpha}=1.0626$ | $\alpha_{\text {exp }}=0.7207+1.3139 i$ |
| spectral radius | $\rho\left(\alpha_{\text {est }}\right)=0.6375$ | $\rho(\tilde{\alpha})=0.7656$ | $\rho\left(\alpha_{\text {exp }}\right)=0.6089$ |
| upper bound | $\omega\left(\alpha_{\text {est }}\right)=0.6409$ | $\sigma(\tilde{\alpha})=0.7657$ | - |
| IT | 37 | 61 | 33 |
| Ex.2 with $(4.2),(4.5)$ |  |  |  |
| parameter | $\alpha_{\text {est }}=0.2088+2.2906 i$ | $\tilde{\alpha}=0.9092$ | $\alpha_{\text {exp }}=0.8768+1.7830 i$ |
| spectral radius | $\rho\left(\alpha_{\text {est }}\right)=0.5683$ | $\rho(\tilde{\alpha})=0.7952$ | $\rho\left(\alpha_{\text {exp }}\right)=0.5395$ |
| upper bound | $\omega\left(\alpha_{\text {est }}\right)=0.5703$ | $\sigma(\tilde{\alpha})=0.7952$ | - |
| IT | 30 | 74 | 28 |

Table 4.5: $\alpha, \rho(\alpha), \omega(\alpha)$ or $\sigma(\alpha)$, and iteration number(IT) of HSS, GMRES, $\operatorname{PGMRES}(\alpha)$ for Ex. 2 with (4.2), (4.4) and different $m, \gamma$.

|  | $m=32$ |  | $m=48$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\gamma=2$ | $\gamma=8$ | $\gamma=3$ | $\gamma=12$ |
| $\alpha_{\text {est }}$ | $0.3520+1.0835 i$ | $0.2012+1.0194 i$ | $0.2640+0.8734 i$ | $0.0436+0.7791 i$ |
| $\rho\left(\alpha_{\text {est }}\right)$ | 0.7368 | 0.7389 | 0.7809 | 0.8148 |
| $\omega\left(\alpha_{\text {est }}\right)$ | 0.7428 | 0.7700 | 0.7891 | 0.8244 |
| IT | 55 | 47 | 68 | 59 |
| $\tilde{\alpha}$ | 0.6624 | 0.4696 | 0.5082 | 0.1860 |
| $\rho(\tilde{\alpha})$ | 0.8474 | 0.8890 | 0.8808 | 0.9545 |
| $\sigma(\tilde{\alpha})$ | 0.8474 | 0.8897 | 0.8808 | 0.9548 |
| IT | 97 | 100 | 123 | 192 |
| GMRES | 54 | 62 | 71 | 90 |
| PGMRES $\left(\alpha_{\text {est }}\right)$ | 14 | 17 | 17 | 23 |
| PGMRES $(\tilde{\alpha})$ | 21 | 23 | 26 | 30 |

With these parameters, in Table 4.5 we also report results for full GMRES method and full preconditioned GMRES with an $\operatorname{HSS}(\alpha)$ preconditioner (1.6), denoted by $\operatorname{PGMRES}(\alpha)$. The code of GMRES we use is from the function gmres in Mablab (7.4 ed) and the preconditioning systems with respect to $\alpha I+H$ and $\alpha I+S$ are solved by the exact factorization method. We test the real parameter $\tilde{\alpha}$ as well as the complex parameter $\alpha_{e s t}$. From these results, we observe that the preconditioner $H S S\left(\alpha_{\text {est }}\right)$ performs much better than the preconditioner $H S S(\tilde{\alpha})$.

Since our estimated parameter $\alpha_{\text {est }}$ depends on the extreme eigenvalues of $H$ and $S$, one may ask how sensitive is the performance of the HSS iteration with respect to the spectrum of $H$ and $S$. We last use an 'approximation', denoted by $\tilde{\alpha}_{\text {est }}$, to our estimated parameter $\alpha_{\text {est }}$ to test the HSS iteration. The 'approximation' $\tilde{\alpha}_{\text {est }}$ is derived by our approach from the approximating eigenvalues of $H$ and $S$ containing $10 \%$ noise of the exact eigenvalues. We test the HSS iteration and PGMRES $\left(\tilde{\alpha}_{e s t}\right)$ for Ex. 2 with (4.2), (4.4) and $m=32, \gamma=2 ; m=48, \gamma=3$.

For $m=32, \gamma=2$, the exact eigenvalues are $\lambda_{1}=8.0221, \lambda_{n}=0.0547, \tau_{1}=8.1271, \tau_{n}=$ 0.1597. The approximating eigenvalues are derived by putting $+10 \%$ error: $\tilde{\lambda}_{1}=8.8243, \tilde{\lambda}_{n}=$ $0.0602, \tilde{\tau}_{1}=8.9398, \tilde{\tau}_{n}=0.1757$. By these approximating eigenvalues, we can derive the approximation $\tilde{\alpha}_{\text {est }}=0.3545+1.2021 i$. The detailed numerical results are listed in Table 4.6. Also we put $-10 \%$ error of the exact eigenvalues to derive an approximation $\tilde{\alpha}_{\text {est }}=$ $0.3167+0.9751 i$.

The results for this example show that the estimated parameter $\alpha_{\text {est }}$ has a normal sensitivity (about $10 \%$ error) to the eigenvalues, while the spectral radius $\rho\left(\alpha_{\text {est }}\right)$ is less sensitive (about $2 \%$ ).

Table 4.6: $\tilde{\alpha}_{e s t}, \rho\left(\tilde{\alpha}_{e s t}\right)$ and iteration number(IT) of $\operatorname{HSS}\left(\tilde{\alpha}_{e s t}\right), \operatorname{PGMRES}\left(\tilde{\alpha}_{e s t}\right)$ for Ex. 2 with (4.2), (4.4) and $m=32, \gamma=2, m=48, \gamma=3$.

|  | $m=32, \gamma=2$ |  | $m=48, \gamma=3$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $+10 \%$ | $-10 \%$ | $+10 \%$ | $-10 \%$ |
| $\tilde{\alpha}_{\text {est }}$ | $0.3874+1.1919 i$ | $0.3167+0.9751 i$ | $0.2902+0.9608 i$ | $0.2377+0.7860 i$ |
| $\rho\left(\tilde{\alpha}_{e s t}\right)$ | 0.7577 | 0.7271 | 0.7988 | 0.7760 |
| IT | 60 | 50 | 75 | 62 |
| PGMRES $\left(\tilde{\alpha}_{e s t}\right)$ | 15 | 13 | 18 | 16 |

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