

SHORT NOTE

A Remark on "An Efficient Real Space Method for Orbital-Free Density-Functional Theory"

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Abstract. In this short note we clarify some issues regarding the existence of minimizers for the Thomas-Fermi-von Weiszacker energy functional in orbital-free density functional theory, when the Wang-Teter corrections are included.

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In [1] it was claimed that there always exists a minimizer; however, the statement of Theorem 2.1 is incomplete. In this note we present the full statement, with a detailed proof.

The theorem stated in [1] holds as long as the number of electrons is below a certain critical value. The correct statement for the theorem in [1] is:

Theorem 1 (Existence of minimizers). *Given $v \in C^\infty(\bar{\Omega})$, and $K_{WT} \in L^2_{loc}(\mathbb{R}^3)$, consider the problem*

$$\inf_{u \in \mathcal{B}} F[u], \quad (1)$$

where F and \mathcal{B} are

$$\begin{aligned} F[u] = & \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \frac{7C_{TF}N^{2/3}}{25} \int_{\Omega} u^{10/3} + \frac{4C_{TF}N^{2/3}}{5} \int_{\Omega} |u|^{5/3} (K_{WT} * |u|^{5/3}) \\ & + \frac{N}{2} \int_{\Omega} u^2 \left(\frac{1}{|\mathbf{x}|} * u^2 \right) - \frac{3}{4} \left(\frac{3N}{\pi} \right)^{1/3} \int_{\Omega} u^{8/3} \\ & + \int_{\Omega} u^2 \varepsilon(Nu^2) + \int_{\Omega} v(\mathbf{x}) u^2(\mathbf{x}) d\mathbf{x}, \end{aligned} \quad (2)$$

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and

$$\mathcal{B} = \left\{ u \in H_0^1(\Omega) \mid u \geq 0, \int_{\Omega} u^2 = 1 \right\}. \tag{3}$$

In (2), the set Ω is open and bounded, and star-shaped with respect to 0; ε is defined as

$$\varepsilon(Nu^2) = \begin{cases} \frac{\gamma}{1 + \beta_1\sqrt{r_s} + \beta_2r_s}, & r_s \geq 1, \\ A\ln(r_s) + B + Cr_s\ln(r_s) + Dr_s, & r_s \leq 1, \end{cases} \tag{4}$$

where $r_s = (4\pi Nu^2/3)^{-\frac{1}{3}}$; the parameters used are $\gamma = -0.1423$, $\beta_1 = 1.0529$, $\beta_2 = 0.3334$, $A = 0.0311$, $B = -0.048$, and $C = 2.019151940622 \times 10^{-3}$ and $D = -1.163206637891 \times 10^{-2}$ are chosen so that $\varepsilon(r)$ and $\varepsilon'(r)$ are continuous at $r=1$ [6].

Then, there exists $N_0 > 0$ such that:

1. If $N < N_0$ then $\exists u^* \in \mathcal{B}$ such that

$$F[u^*] = \min_{u \in \mathcal{B}} F[u]. \tag{5}$$

2. If $N > N_0$ then

$$\inf_{u \in \mathcal{B}} F[u] = -\infty. \tag{6}$$

Proof. The second part of the theorem was proved in [2, 3]. We outline the proof here for completeness. Since $0 \in \Omega$, $\exists \delta_0 > 0$ such that $B(0, \delta_0) \subset \Omega$. Consider a compactly supported function $u_0 \in C_0^\infty(B(0,1))$, such that

$$\int_{\mathbb{R}^3} u_0^2 = 1, \tag{7}$$

and consider the rescaling

$$u_\delta(\mathbf{x}) = \frac{1}{\delta^{3/2}} u_0\left(\frac{\mathbf{x}}{\delta}\right), \quad 0 < \delta < \delta_0. \tag{8}$$

Then $u_\delta \in \mathcal{B}$, and

$$F[u_\delta] = \frac{1}{\delta^2} \left(\frac{1}{2} \int_{\Omega} |\nabla u_0|^2 - \frac{7C_{TF}N^{2/3}}{25} \int_{\Omega} u_0^{10/3} \right) + \mathcal{O}\left(\frac{1}{\delta}\right). \tag{9}$$

Define

$$A_0 = \inf_{u \in H_0^1(\Omega), \|u\|_2=1} \frac{\int_{\Omega} |\nabla u|^2}{\int_{\Omega} u^{10/3}} > 0. \tag{10}$$

Then if $A_0/2 < 7C_{TF}N^{2/3}/25$, we can choose u_0 so that the leading term in (9) is negative, and when $\delta \rightarrow 0$, the desired result follows.

For the existence of minimizers, assume that N is such that $A_0/2 > 7C_{TF}N^{2/3}/25$. By Lemma 1, there exist $C > 0, \delta > 0$ such that

$$\begin{aligned} F[u] &\geq \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \left(\frac{7C_{TF}N^{2/3}}{25} + \delta \right) \int_{\Omega} u^{10/3} - C \\ &\geq \left(\frac{1}{2} - \frac{1}{A_0} \left(\frac{7C_{TF}N^{2/3}}{25} + \delta \right) \right) \int_{\Omega} |\nabla u|^2 \geq \tau \int_{\Omega} |\nabla u|^2 - C, \end{aligned} \quad (11)$$

where $\tau > 0$. Therefore the functional is coercive, and the result follows from now from standard arguments in the Calculus of Variations [4], involving the Sobolev Embedding, and the Rellich-Kondrachov compactness theorem. \square

Remark 1. Note that given $\Omega \subset \mathbf{R}^3$, then

$$0 < A_0 = \inf_{u \in \mathcal{A}} \frac{\int_{\Omega} |\nabla u|^2}{\int_{\Omega} u^{10/3}}, \quad (12)$$

where

$$\mathcal{A} = \left\{ u \in H_0^1(\Omega) \mid u \geq 0, \int_{\Omega} u^2 = 1 \right\}. \quad (13)$$

By the Gagliardo-Nirenberg inequality, $\exists C_1 > 0$ such that

$$\left(\int_{\Omega} u^6 \right)^{1/3} \leq C_1 \int_{\Omega} |\nabla u|^2. \quad (14)$$

By the Riesz-Thorin theorem, since $u \in L^2(\Omega) \cap L^6(\Omega)$, and

$$\frac{3}{10} = \frac{\theta}{2} + \frac{1-\theta}{6}, \quad (15)$$

with $\theta = 2/5$, we get

$$\left(\int_{\Omega} u^{10/3} \right)^{3/10} \leq \left(\int_{\Omega} u^2 \right)^{\theta/2} \left(\int_{\Omega} u^6 \right)^{(1-\theta)/6}, \quad (16)$$

and therefore, since $\|u\|_2 = 1$,

$$\int_{\Omega} u^{10/3} \leq \left(\int_{\Omega} u^6 \right)^{5(1-\theta)/9} = \left(\int_{\Omega} u^6 \right)^{1/3} \leq C_1 \int_{\Omega} |\nabla u|^2. \quad (17)$$

Therefore,

$$\inf_{u \in \mathcal{A}} \frac{\int_{\Omega} |\nabla u|^2}{\int_{\Omega} u^{10/3}} \geq C_1^{-1} > 0. \quad (18)$$

In [1] it was proved that $K_{WT} \in L^2(\mathbb{R}^3)$. In the following lemma we establish the necessary inequalities to prove the coercivity of energy functional (2).

Lemma 1. *Assume $K_{WT} \in L^2(\mathbb{R}^3)$, $v \in L^\infty(\Omega)$, and ϵ is defined as in (4). Then, there exist constants $C_i, i = 1, \dots, 5$, dependent only on the domain Ω and on N , such that for all $u \in H_0^1(\Omega)$ satisfying $\|u\|_2 = 1$,*

$$\left| \int_{\Omega} |u|^{5/3} (K_{WT} * |u|^{5/3}) \right| \leq C_1 \|u^{5/3}\|_2 \|u^{5/3}\|_1 \|K_{WT}\|_2; \tag{19}$$

$$\left| \int_{\Omega} \left(u^2 * \frac{1}{|\mathbf{x}|} \right) u^2 \right| \leq C_2 \|u^2\|_{5/3} \|u\|_2^{7/3}; \tag{20}$$

$$\left| \int_{\Omega} u^{8/3} \right| \leq C_3 \|u^{5/3}\|_2 \|u\|_2; \tag{21}$$

$$\left| \int_{\Omega} u^2 \epsilon(Nu^2) \right| \leq C_4 + C_5 \left(\int_{\Omega} |u|^{10/3} \right)^{3/4}. \tag{22}$$

Proof. Since $K_{WT} \in L^2$, by the Cauchy-Schwarz inequality, followed by Young's inequality:

$$\begin{aligned} \left| \int_{\Omega} |u|^{5/3} (K_{WT} * |u|^{5/3}) \right| &\leq \|u^{5/3}\|_2 \|K_{WT} * |u|^{5/3}\|_2 \\ &\leq C_1 \|u^{5/3}\|_2 \|K_{WT}\|_2 \|u^{5/3}\|_1. \end{aligned} \tag{23}$$

Note that since $\|u\|_2 = 1$, by Hölder's inequality, $\|u^{5/3}\|_1 \leq |\Omega|^{1/6}$. This gives (19). The inequality (20) was proved in [5] (Theorem IV.1, page 75). The estimate (21) follows from the Cauchy-Schwarz inequality:

$$\left| \int_{\Omega} u^{8/3} \right| = \left| \int_{\Omega} u^{5/3} u \right| \leq C \|u^{5/3}\|_2 \|u\|_2. \tag{24}$$

From the definition of ϵ , we get that

$$\begin{aligned} \left| \int_{\Omega} u^2 \epsilon(Nu^2) \right| &\leq C_1 + \tilde{C}_2 \left| \int_{|u| \geq \frac{3}{4\pi N}} u^2 \log |u| \right| \\ &\leq C_1 + \hat{C}_2 \left| \int_{\Omega} |u|^{5/2} \right| \leq C_1 + C_2 \left(\int_{\Omega} |u|^{10/3} \right)^{3/4}. \end{aligned} \tag{25}$$

This concludes the proof. □

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References

- [1] C. J. García-Cervera, An efficient real space method for orbital-free density-functional theory, *Commun. Comput. Phys.*, 2(2) (2007), 334-357.
- [2] X. Blanc and E. Cancès, Nonlinear instability of density-independent orbital-free kinetic-energy functionals, *J. Chem. Phys.*, 122 (2005), 214106.
- [3] X. Blanc and E. Cancès, Technical report, <http://www.ann.jussieu.fr/publications/2005/R05014.pdf>, 2005.
- [4] B. Dacorogna, *Direct Methods in the Calculus of Variations*, Applied Mathematical Sciences, 78, Springer-Verlag, Berlin-New York, 1989.
- [5] E. H. Lieb and B. Simon, The Thomas-Fermi theory of atoms, molecules and solids, *Adv. Math.*, 23(1) (1977), 22-116.
- [6] J. P. Perdew and A. Zunger, Self interaction correction to density functional approximations for many electron systems, *Phys. Rev. B*, 23(10) (1981), 5048-5079.