

GLOBAL ATTRACTORS OF REACTION-DIFFUSION SYSTEMS AND THEIR HOMOGENIZATION*

Zhang Xingyou and Hu Xiaohong

(College of Math. and Physics, Chongqing University, 400030, China)

(E-mail: zhangxy@cqu.edu.cn; xiaohongcq@hotmail.com)

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Abstract In this paper, we study the existence of the global attractor \mathcal{A}^ε of reaction-diffusion equation

$$\partial_t u^\varepsilon(x, t) = A_\varepsilon u^\varepsilon(x, t) - f(x, \varepsilon^{-1}x, u^\varepsilon(x, t)),$$

and the homogenized attractor \mathcal{A}^0 of the corresponding homogenized equation, then give explicit estimates for the distance between the attractor \mathcal{A}^ε and the homogenized attractor \mathcal{A}^0 .

Key Words Homogenization; global attractor; reaction-diffusion systems; almost-periodic function; Diophantine conditions.

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1. Introduction and Main Results

We consider the reaction-diffusion system

$$\begin{cases} \partial_t u^\varepsilon(x, t) = A_\varepsilon u^\varepsilon(x, t) - f(x, \varepsilon^{-1}x, u^\varepsilon(x, t)), & (x, t) \in \Omega \times \mathbf{R}^+, \\ u^\varepsilon(x, t)|_{\partial\Omega} = 0, \quad u^\varepsilon(x, t)|_{t=0} = u_0, \end{cases} \quad (1.1)$$

where Ω is a bounded domain in \mathbf{R}^3 and $0 < \varepsilon \leq \varepsilon_0 < 1$. Here $u^\varepsilon = u^\varepsilon(x, t) = (u_\varepsilon^1, \dots, u_\varepsilon^k)$ is an unknown vector-valued function. The second order elliptic differential operators A_ε have the form as follows:

$$A_\varepsilon u := \text{diag}(A_\varepsilon^1 u^1, \dots, A_\varepsilon^k u^k), \quad (1.2)$$

with

$$A_\varepsilon^l u^l = \sum_{i,j=1}^3 \partial_{x_i} (a_{ij}^l(\varepsilon^{-1}x) \partial_{x_j} u^l(x)), \quad (1.3)$$

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where the functions $a_{ij}^l(y)$, $l = 1, \dots, k$, $y \in \mathbf{R}^3$, are assumed to be symmetric, smooth and \mathbf{Y} -periodic with respect to $y \in \mathbf{R}^3$, where $\mathbf{Y} \subset \mathbf{R}^3$ is a fixed cube. The uniform ellipticity condition

$$\sum_{i,j=1}^3 a_{ij}^l(y) \zeta_i \zeta_j \geq \nu |\zeta|^2, \quad \forall y, \zeta \in \mathbf{R}^3, \quad (1.4)$$

is also assumed (with an appropriate $\nu > 0$) to be valid for operators A_ε^l . We impose that $f(x, y, u)$ is almost-periodic ([1]) with respect to $y \in \mathbf{R}^3$ and satisfies the conditions as follows:

$$f \in C^1(\mathbf{R}^k, \mathbf{R}^k), \quad \partial_z f(x, y, z) \zeta \zeta \geq -C_2 \zeta \zeta, \quad \forall \zeta \in \mathbf{R}^k, \quad (1.5)$$

$$|f(x, y, u)| \leq C(1 + |u|^p), \quad \forall (x, y) \in \Omega \times \mathbf{R}^3, \quad (1.6)$$

$$\sum_{l=1}^k f^l u^l |u^l|^{p_l} \geq C \sum_{l=1}^k |u^l|^{p_l+2} - C_1, \quad \forall u \in \mathbf{R}^k, \quad (1.7)$$

where $p \geq 1, p_i \geq 2(p-1)$, $i = 1, \dots, k$. It is assumed also that the initial data $u_0 \in (L^2(\Omega))^k$.

Efendiev and Zelik (see [2]) studied the problem (1.1) when $f(x, y, u)$ is independent of y . Fiedler and Vishik (see [3]) studied the case when the $A_\varepsilon u$ in (1.1) is replaced by $a\Delta u$. In fact, one can obtain the existence of solutions and attractors for (1.1) with $f(x, y, u)$ depending on y by the standard method as those in [4]. However, when estimate the distance between the attractors for (1.1) and the attractors of the homogenized equation, the arguments in [2] or [3] don't work. We have to overcome these difficulties by combining the ideas in [3], [2] and analyzing carefully the properties of periodic and almost-periodic functions.

In order to simplify our expression, we denote $H = (L^2(\Omega))^k$, $V = (W_0^{1,2}(\Omega))^k$, $F = (L^\infty(\Omega))^k$, $\|\cdot\|_{(W^{l,p}(\Omega))^k} = \|\cdot\|_{l,p}$.

Theorem 1.1 *If the assumptions (1.2) – (1.7) hold, and the initial data $u_0 \in H$, then for any $T > 0$, $\varepsilon > 0$, the problem (1.1) possesses a unique solution $u^\varepsilon(x, t) \in L^\infty([0, T]; H) \cap L^2([0, T]; V)$, $u^\varepsilon \in C(R^+; H)$. The mapping $S_t^\varepsilon: u_0 \rightarrow u^\varepsilon(x, t)$ defines a continuous semigroup $S_t^\varepsilon: H \rightarrow H$. If, furthermore, $u_0 \in V$, then $u^\varepsilon(x, t) \in L^\infty([0, T]; V) \cap L^2([0, T]; W^{2,2}(\Omega))$, $u^\varepsilon \in C(R^+; V)$.*

Theorem 1.2 *If the assumptions (1.2) – (1.7) hold, and $u_0 \in H$, then for every $\varepsilon > 0$, the semigroup S_t^ε generated by the equation (1.1) possesses a global compact attractor \mathcal{A}^ε in H .*

Theorem 1.1 can be proved by the Faedo-Galerkin method with the help of R.Temam [4], and the details of the proof are omitted. Similar arguments as in [4] for the problem (1.1) yield the a priori estimates needed about $u^\varepsilon(x, t)$ in H and V , and we omit the details. Then Theorem 1.2, whose proof is also omitted, can be easily proved by the standard arguments [4, Theorem 1.1.1].

By the standard homogenization theory, one can obtain the homogenized problem (2.11), for which one can prove the similar results to Theorems 1.1 and 1.2. In order to

estimate the L^2 -distance between the attractors for (1.1) and the attractors of the homogenized equation (2.11), we can obtain the a priori estimates required, whose proofs are also omitted, by the similar arguments as those in [2] under the better initial data condition (see Section 2). Under some additional assumptions (mainly the so-called Diophantine conditions (2.21)), we have

Theorem 1.3 *Let the assumptions of Theorem 1.2, (2.1), (2.2) and the assumptions of Proposition 2.2 (see Section 2) hold. Let $u_0 \in F \cap V$ and let $u^\varepsilon(x, t)$ be the solution, defined in Theorem 1.1, of the problem (1.1), $u^0(x, t) \in L^\infty([0, T]; H) \cap L^2([0, T]; V)$ be the solution of the problem (2.11), then $\forall t > 0$, we have*

$$\|u^\varepsilon(x, t) - u^0(x, t)\|_H \leq C\varepsilon^{\frac{2}{3}}e^{\beta t},$$

where the constant $C > 0$ depends only on $\|u_0\|_{F \cap V}$ and $\beta > 0$ is a constant independent of u^ε and u^0 .

Theorem 1.4 *Let the assumptions of Theorem 1.3 and (2.39) hold. Let \mathcal{A}^ε be the global attractor of the equation (1.1) and \mathcal{A}^0 be the global attractor of the homogenized equation (2.11), and define the fractional convergence rate $k = \frac{2\rho}{3\rho + 3\beta}$, then there exists a constant $C > 0$ such that*

$$d(\mathcal{A}^\varepsilon, \mathcal{A}^0) := \text{dist}_H(\mathcal{A}^\varepsilon, \mathcal{A}^0) \leq C\varepsilon^k, \quad 0 < \varepsilon \leq \varepsilon_0.$$

2. The Homogenization and the Estimates of Errors

First, we study the homogenization of the problem (1.1). In addition to the assumptions (1.2)–(1.7), we assume the initial data $u_0 \in F \cap V$ and the $f(x, y, z)$ satisfies the conditions as follows:

$$f^l(x, y, z) = \sum_{j=1}^q b_l^j(x, y) f_{jl}(z), \quad |b_l^j(x, y)| \leq C, \quad (2.1)$$

where $f^l(x, y, z)$, $l = 1, \dots, k$, are the components of $f(x, y, z)$. Let

$$\sum_{l=1}^k |\partial_z f^l(x, y, z)| \leq C_1(|z|^4 + 1). \quad (2.2)$$

Recall that $w \in AP(\mathbf{R}^3)$ (the set of almost-periodic functions) possesses the mean value which can be calculated by :

$$\langle w \rangle = \langle w \rangle_x := \lim_{T \rightarrow \infty} \frac{1}{2^3 T^3} \int_{[-T, T]^3} w(x) dx, \quad (2.3)$$

and the Fourier expansion as follows (see [5])

$$w(x) = \sum_{\hat{w}(\xi) \neq 0} \hat{w}(\xi) e^{i(x, \xi)}, \quad (2.4)$$

where the amplitudes $\hat{w}(\xi) \in \mathbb{C}$, $\xi \in \mathbf{R}^3$, defined by $\hat{w}(\xi) = \langle w(x)e^{-i(x,\xi)} \rangle$. We denote by $\text{Trig}(\mathbf{R}^3)$ the space of all finite trigonometric polynomials of the form (2.4)

$$\text{Trig}(\mathbf{R}^3) := \left\{ w(x) = \sum_{k=1}^K w_k e^{i(x,\xi_k)} : K \in \mathbf{N}, \xi_k \in \mathbf{R}^3, w_k \in \mathbb{C} \right\}. \tag{2.5}$$

We state a classical result in the homogenization theory:

Proposition 2.1 ([6, 7]) *Let $g \in W^{-1,2}(\Omega)$ and $v^\varepsilon \in V$ be the solution of the equation $A_\varepsilon v^\varepsilon = g$, where the operator A_ε is defined by (1.3). Then,*

$$\begin{cases} v^\varepsilon \rightharpoonup v^0 & \text{weakly in } V, \\ A_\varepsilon v^\varepsilon \rightharpoonup A_0 v^0 & \text{weakly in } H, \end{cases} \tag{2.6}$$

where $v^0 \in V$ is a unique solution of the homogenized problem $A_0 v^0 = g$. The operator A_0 is defined by the form as follows:

$$A_0^l v^{0l} = \sum_{i,j=1}^3 \partial_{x_i} (a_{ij}^{0l} \partial_{x_j} v^{0l}), \quad A_0 v := \text{diag}(A_0^1 v^1, \dots, A_0^k v^k), \tag{2.7}$$

and the so-called homogenized coefficients $a_{ij}^{0l} = \langle a_{ij}^l(y) \rangle + \sum_{m=1}^3 \langle a_{im}^l(y) \partial_{y_m} N_m^l(y) \rangle$ are constants, where the \mathbf{Y} -periodic correctors $N_m^l(y)$, $m = 1, 2, 3$, $l = 1, \dots, k$, are the solutions of the auxiliary periodic problem as follows:

$$\sum_{i,j=1}^3 \partial_{y_i} (a_{ij}^l(y) \partial_{y_j} N_m^l(y)) = - \sum_{i=1}^3 \partial_{y_i} (a_{im}^l(y)), \quad y \in \mathbf{R}^3. \tag{2.8}$$

And the homogenized matrix A_0 satisfies the coerciveness condition (1.4).

The following lemma, whose proof is easy and so omitted, will be used in the sequel.

Lemma 2.1 *Let Assumptions (1.6), (2.1) hold and $f(x, y, u^\varepsilon)$ be almost-periodic in y , assume $u^\varepsilon \rightarrow u^0$ in H ($\varepsilon \rightarrow 0$), and denote $f_0(x, u^0) := \langle f(x, y, u^0) \rangle_y$, then we have the result as follows:*

$$f(x, \varepsilon^{-1}x, u^\varepsilon) \rightharpoonup f_0(x, u^0) \quad \text{weakly in } H. \tag{2.9}$$

$$f^l(x, u^0) = \sum_{j=1}^q b_l^{0j}(x) f_{jl}(u^0). \tag{2.10}$$

Now by the standard homogenization theory we obtain the homogenized problem

$$\begin{cases} \partial_t u^0 = A_0 u^0 - f_0(x, u^0), & (x, t) \in \Omega \times \mathbf{R}^+, \\ u^0|_{\partial\Omega} = 0, \quad u^0|_{t=0} = u_0. \end{cases} \tag{2.11}$$

Note that this equation satisfies all assumptions of the equation (1.1), consequently, it admits a unique solution $u^0(x, t) \in L^\infty([0, T]; H) \cap L^2([0, T]; V)$ and (2.11) possesses a global attractor \mathcal{A}^0 in H .

We now specify additional conditions which enable us to estimate the distance between the solutions $u^\varepsilon(x, t)$ and $u^0(x, t)$ in the norm of H . In order to give the distance estimate of $u^\varepsilon(x, t)$ and $u^0(x, t)$ in H , we need three propositions (see [2, 3]).

First, we introduce some results about divergence representations. Let $h(x, y) = h(x_1, \dots, x_3, y_1, \dots, y_3)$ be a sufficiently smooth function which is almost-periodic in $y = (y_1, \dots, y_3)$, ie :

(i) there exists a function $H(x, w_1, \dots, w_3) = H(x_1, \dots, x_3, w_{11}, \dots, w_{1k_1}, \dots, w_{31}, \dots, w_{3k_3})$ which is 2π -periodic with respect to each w_{ij} . Here $w_i = (w_{i1}, \dots, w_{ik_i}) \in \mathbf{R}^{k_i}$. ($i = 1, \dots, 3$)

(ii) there exists rationally independent frequency $\alpha_{11}, \dots, \alpha_{1k_1}, \dots, \alpha_{3k_3}$ such that

$$h(x, y) = H(x_1, \dots, x_3, \alpha_1 y, \dots, \alpha_3 y), \quad (2.12)$$

where $\alpha_l = (\alpha_{l1}, \dots, \alpha_{lk_l})$. Let $\tilde{H}(x, w) = H(x, w) - H_0(x)$, where

$$H_0(x) = |T^k|^{-1} \int_{T^k} H(x, w_1, \dots, w_3) dw_1 \cdots dw_3, \quad (2.13)$$

where $T^k = T^{k_1} \times \cdots \times T^{k_3}$, and $T^{k_i} = \mathbf{R}^{k_i} / (\mathbf{Z} \cdot 2\pi)^{k_i}$ is the k_i -dimensional torus. Assume that the Fourier series

$$H(x, w) = \sum_m H_m(x) e^{im \cdot w} \quad (2.14)$$

is convergent. Let

$$\tilde{h}(x, y) = h(x, y) - H_0(x) = \sum_{m \neq 0} H_m(x) \exp \left(i \sum_{j=1}^3 m_j \alpha_j y_j \right), \quad (2.15)$$

where $m_j = (m_{j1}, \dots, m_{jk_j}) \in \mathbf{Z}^{k_j}$, $\alpha_j \in \mathbf{R}^{k_j}$ and $y_j \in \mathbf{R}$. For any such almost periodic function $h(x, y)$, we construct a corresponding divergence representation by function $S_\sigma(x, y)$, $\sigma = 1, \dots, 3$.

$$\tilde{h}(x, y) = \sum_{\sigma=1}^3 \partial_{y_\sigma} S_\sigma(x, y). \quad (2.16)$$

We shall find $S_\sigma(x, y)$ of the form

$$S_\sigma(x, y) = \sum_{m \in \mathbf{Z}^k \setminus \{0\}} \eta_m^\sigma(x) \exp \left(i \sum_{j=1}^3 m_j \alpha_j y_j \right). \quad (2.17)$$

From (2.15) – (2.17) we derive:

$$\sum_{m \neq 0} H_m(x) \exp \left(i \sum_{j=1}^3 m_j \alpha_j y_j \right) = \tilde{h}(x, y) = \sum_{m \neq 0} \sum_{\sigma=1}^3 m_\sigma \cdot \alpha_\sigma \eta_m(x) \exp \left(i \sum_{j=1}^3 m_j \alpha_j y_j \right). \quad (2.18)$$

So (2.16) will hold if

$$\sum_{\sigma=1}^3 m_{\sigma} \cdot \alpha_{\sigma} \eta_m^{\sigma}(x) = -iH_m(x), \tag{2.19}$$

for all $m \in \mathbf{Z}^k \setminus \{0\}$. Let the following assumptions be satisfied for some positive δ and δ' :

$$\tilde{b}_l^j = \tilde{h}_l^j = \sum_{m \neq 0} H_{lm}^j(x) \exp\left(i \sum_j m_j \alpha_j y_j\right). \tag{2.20}$$

$$|m_{\sigma} \cdot \alpha_{\sigma}| \geq c|m_{\sigma}|^{-(k_{\sigma}-1+\delta)}, \quad \forall m_{\sigma} \in \mathbf{Z} \setminus \{0\}. \tag{2.21}$$

$$\|H_{lm}^j(x)\|_{C^0(\bar{\Omega})} \leq c(1 + |m_{\sigma}|)^{-(k_{\sigma}-1+\delta)}(1 + |m|)^{-(k+\delta')}. \tag{2.22}$$

$$\|\partial_{x_{\sigma}} H_{lm}^j(x)\|_{L^3(\Omega)} \leq c(1 + |m_{\sigma}|)^{-(k_{\sigma}-1+\delta)}(1 + |m|)^{-(k+\delta')}. \tag{2.23}$$

Now we can state the propositions as follows:

Proposition 2.2([3]) *Let the coefficients $b_l^j(x, y)$ of (2.1) satisfy the conditions as follows:*

- (i) $b_l^j(x, y)$ are almost-periodic in y , $j = 1, \dots, q$;
 - (ii) the corresponding frequencies α_{ij} satisfy Diophantine condition (2.21);
 - (iii) the coefficients $H_{lm}^j(x)$ in the series (2.20) of $\tilde{b}_l^j(x, y) = b_l^j(x, y) - b_l^{0j}(x)$ satisfy the decay conditions (2.22), (2.23),
- then we can represent $\tilde{b}_l^j(x, y)$ in the form

$$\tilde{b}_l^j(x, y) = \sum_{\sigma=1}^3 \partial_{y_{\sigma}} S_{l\sigma}^j(x, y), \tag{2.24}$$

which satisfies

$$|S_{l\sigma}^j(x, y)| \leq C_0, \quad \|\partial_{x_{\sigma}}^1 S_{l\sigma}^j(x, y)\|_{L^3(\Omega)} \leq C_0, \tag{2.25}$$

here $\partial_{x_{\sigma}}^1$ indicates partial derivatives with respect to the first argument x of the function $S_{l\sigma}^j(x, y)$.

Proposition 2.3([3]) *Let the assumptions (1.2)-(1.7), (2.1), (2.2) and Proposition 2.2 hold. Then*

$$\left| (f(x, \varepsilon^{-1}x, u^{\varepsilon}) - f_0(x, u^{\varepsilon}), u^{\varepsilon} - u^0) \right| \leq \varepsilon C \|u^{\varepsilon} - u^0\|_V, \tag{2.26}$$

where the constant $C > 0$ depends only on $\|u_0\|_{F \cap V}$.

Denote (see [5]):

$$u_1^{\varepsilon}(t) = u^0(t) + \varepsilon \sum_{k=1}^3 N_k(\varepsilon^{-1}x) \partial_{x_k} u^0(t), \tag{2.27}$$

where $N_k(\varepsilon^{-1}x)$, $k = 1, 2, 3$, are the solutions of the problem (2.8). Note that the function $u_1^{\varepsilon}(t)$ doesn't satisfy the 0-Dirichlet boundary condition. In order to avoid this

difficulty, we introduce a family of cut-off functions $\tau^\varepsilon(x)$ satisfying two conditions as follows (see [5]): (1) $\tau^\varepsilon(x) \in C_0^\infty(\Omega)$, $0 \leq \tau^\varepsilon \leq 1$, $\tau^\varepsilon(x) \equiv 1$ off the ε -neighborhood of the boundary of Ω ; (2) $\varepsilon |\nabla_x \tau^\varepsilon(x)| \leq C$ in Ω , where the constant C is independent of ε . Thus we take

$$\begin{aligned} w^\varepsilon(t) &= u_1^\varepsilon(t) - \varepsilon(1 - \tau^\varepsilon(x)) \sum_{k=1}^3 N_k(\varepsilon^{-1}x) \partial_{x_k} u^0(t) \\ &= u^0(t) + \varepsilon \tau^\varepsilon(x) \sum_{k=1}^3 N_k(\varepsilon^{-1}x) \partial_{x_k} u^0(t). \end{aligned} \quad (2.28)$$

Then, obviously, $w^\varepsilon(t) \in V$. we need the proposition as follows:

Proposition 2.4([2]) *Let the assumption (1.4) hold, and let $w^\varepsilon(t)$, $A_\varepsilon u^\varepsilon$, $A_0 u^0$ be defined by (2.28), (1.3), (2.7) respectively, $u^\varepsilon(t)$, $u^0(t)$ be the solution of the equation (1.1), (2.11) respectively. Then*

$$(A_\varepsilon u^\varepsilon(t) - A_0 u^0(t), u^\varepsilon(t) - w^\varepsilon(t)) \leq C \varepsilon^{\frac{2}{3}} \|u^0(t)\|_{2,2}^2, \quad (2.29)$$

where the constant $C > 0$ is independent of ε .

Proof of Theorem 1.3 Denote $v(x, t) = u^\varepsilon(x, t) - u^0(x, t)$. Subtracting (2.11) from (1.1), we get

$$\partial_t v = A_\varepsilon u^\varepsilon - A_0 u^0 - (f(x, \varepsilon^{-1}x, u^\varepsilon) - f_0(x, u^\varepsilon)) - (f_0(x, u^\varepsilon) - f_0(x, u^0)). \quad (2.30)$$

Multiplying both sides of (2.30) by v and integrating over Ω , we obtain

$$\begin{aligned} (\partial_t v, v) &= (A_\varepsilon u^\varepsilon - A_0 u^0, v) - (f(x, \varepsilon^{-1}x, u^\varepsilon) - f_0(x, u^\varepsilon), v) \\ &\quad - (f_0(x, u^\varepsilon) - f_0(x, u^0), v). \end{aligned} \quad (2.31)$$

To prove the theorem, we estimate each term of the right-hand side of (2.31) respectively. Using Proposition 2.4, we derive

$$\begin{aligned} \sum_{l=1}^k (A_\varepsilon^l u_\varepsilon^l(t) - A_0^l u^{0l}(t), v^l(t)) &= \sum_{l=1}^k (A_\varepsilon^l u_\varepsilon^l(t) - A_0^l u^{0l}(t), u_\varepsilon^l(t) - w^{\varepsilon l}(t)) \\ &\quad + \sum_{l=1}^k (A_\varepsilon^l u_\varepsilon^l(t) - A_0^l u^{0l}(t), v^l(t) - u_\varepsilon^l(t) + w^{\varepsilon l}(t)) \\ &\leq C \varepsilon^{\frac{2}{3}} \|u^0\|_{2,2}^2 + \|A_\varepsilon u^\varepsilon - A_0 u^0\|_H \cdot \|v - u^\varepsilon + w^\varepsilon\|_H. \end{aligned} \quad (2.32)$$

Note that the definitions (2.27), (2.28) and (2.5) imply the estimate

$$\|v(t) - u^\varepsilon(t) + w^\varepsilon(t)\|_H \leq C \varepsilon \|u^0(t)\|_V. \quad (2.33)$$

Similar methods as in [2] for the equation (1.1) and (2.11) yield

$$\begin{aligned} \int_T^{T+1} \|A_\varepsilon u^\varepsilon(t)\|_H^2 dt + \int_T^{T+1} \|A_0 u^0(t)\|_H^2 dt + \int_T^{T+1} \|u^0(t)\|_V^2 dt + \int_T^{T+1} \|u^0(t)\|_{2,2}^2 dt \\ \leq Q(\|u_0\|_{F \cap V}), \end{aligned} \quad (2.34)$$

for the appropriate function Q independent of $T \geq 0$ (here we have implicitly used the elliptic regularity estimate $\|u^0\|_{2,2} \leq C\|A_0 u^0\|_H$). Inserting the estimate (2.33) to (2.32) and integrating over $t \in [0, T]$ then taking the estimate (2.34) into account, we have

$$\sum_{l=1}^k \int_0^T \left(A_\varepsilon^l u_\varepsilon^l(t) - A_0^l u_0^l(t), v^l \right) dt \leq \varepsilon^{\frac{2}{3}} Q(\|u_0\|_{F \cap V}) T. \quad (2.35)$$

Applying (2.26) to the second term of the right-hand side of (2.31), integrating over $t \in [0, T]$, using Minkowski-inequality and (2.34), we obtain

$$\int_0^T |f(x, \varepsilon^{-1}x, u^\varepsilon(t)) - f_0(x, u^\varepsilon(t)), v(t)| dt \leq \varepsilon Q_1(\|u_0\|_{F \cap V}) T. \quad (2.36)$$

Assumption (2.2) implies

$$\begin{aligned} \int_0^T |(f_0(x, u^\varepsilon) - f_0(x, u^0), v)| dt &= \int_0^T \left| \left(\int_0^1 f'(su^\varepsilon + (1-s)u^0) ds \cdot v, v \right) \right| dt \\ &\leq C_2 \int_0^T \|v\|_H^2 dt. \end{aligned} \quad (2.37)$$

Integrating (2.31) over $t \in [0, T]$ and taking account of (2.35)-(2.37), we get

$$\|v(T)\|_{0,2}^2 \leq \varepsilon^{\frac{2}{3}} Q(\|u_0\|_{F \cap V}) T + 2\varepsilon Q_1(\|u_0\|_{F \cap V}) T + 2C_2 \int_0^T \|v(T)\|_H^2 dt. \quad (2.38)$$

Applying Gronwall's inequality to (2.38) proves Theorem 1.3 .

Now we are ready to derive the error's estimates for the global attractors \mathcal{A}^ε and \mathcal{A}^0 . To this end, we need some additional information about \mathcal{A}^0 which we in fact require to be exponentially attracting with exponential rate $\rho > 0$. We assume there exists a constant $C = C(\varepsilon_0)$ such that for all $t \geq 0$

$$d := \text{dist}_H(u^0, \mathcal{A}^0) \leq C e^{-\rho t}, \quad (2.39)$$

holds, uniformly for all $u_0 \in \bigcup_{0 < \varepsilon \leq \varepsilon_0} \mathcal{A}^\varepsilon$, where dist_H means the nonsymmetric Hausdorff distance (see [4]), i.e.

$$\text{dist}_H(A, B) := \sup_{x \in A} \inf_{y \in B} \|x - y\|_H. \quad (2.40)$$

Proof of Theorem 1.4 Let

$$\mathcal{B} := \bigcup_{0 < \varepsilon \leq \varepsilon_0} \mathcal{A}^\varepsilon. \quad (2.41)$$

Pick $0 < \varepsilon \leq \varepsilon_0$ and $u^\varepsilon \in \mathcal{A}^\varepsilon \subset \mathcal{B}$, arbitrarily. For $t \geq 0$ chosen below consider $u_0 \in \mathcal{A}^\varepsilon$ such that

$$S_t^\varepsilon u_0 = u^\varepsilon. \quad (2.42)$$

Then Theorem 1.3 and (2.39) imply

$$d(u^\varepsilon, \mathcal{A}^0) \leq d(u^\varepsilon, u^0) + d(u^0, \mathcal{A}^0) \leq C\varepsilon^{\frac{2}{3}}e^{\beta t} + Ce^{-\rho t}. \quad (2.43)$$

Choose $t \geq 0$, such that $\varepsilon^{\frac{2}{3}}e^{\beta t} = e^{-\rho t}$, thus $t = -\frac{\ln \varepsilon}{\beta + \rho}$. Substituting this choice of t back into (2.43), because of the arbitrariness of u^ε , we prove Theorem 1.4.

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