# AN H-BASED A – $\phi$ METHOD WITH A NONMATCHING GRID FOR EDDY CURRENT PROBLEM WITH DISCONTINUOUS COEFFICIENTS \*1)

Tong Kang

(Department of Applied Mathematics, School of Sciences, Beijing Broadcasting Institute, Beijing 100024, China)

Zheng-peng Wu

(Department of Applied Mathematics, School of Sciences, Beijing Broadcasting Institute, Beijing 100024, China)

(ICMSEC, Academy of Mathematics and System Sciences, Chinese Academy of Sciences, Beijing 100080, China)

De-hao Yu

(LSEC, ICMSEC, Academy of Mathematics and System Sciences, Chinese Academy of Sciences, Beijing 100080, China)

#### Abstract

In this paper, we investigate the finite element  $\mathbf{A} - \phi$  method to approximate the eddy current equations with discontinuous coefficients in general three-dimensional Lipschitz polyhedral eddy current region. Nonmatching finite element meshes on the interface are considered and optimal error estimates are obtained.

Mathematics subject classification: 65M60, 65N30, 35Q60. Key words: Eddy current problem, Finite element  $\mathbf{A} - \phi$  method, Nonmatching meshes, Error estimates.

### 1. Introduction

The eddy current model emerges from Maxwell's equations by formally dropping the displacement currents, which amounts to neglecting capacitive effects (space charges) and provides a reasonable approximation in the low frequency range and in the presence of high conductivity. Various formulations of the eddy current problem have been suggested in [1], which differ in their choice of the primary unknown. Instead of finding magnetic and electric fields directly, the  $\mathbf{A} - \phi$  approach is to seek vector and scalar potentials. Although this method increases the number of scalar unknowns and equations, this apparent complication is justified by a better way of dealing with the possible discontinuities in process of the numerical approximations.

It is well-known that the  $\mathbf{A} - \phi$  method has been applied to the eddy current model extensively in practice, but further theoretical research in this aspect has rarely shown so far. For some recent relative work, we refer readers to [2, 8-12] for eddy current problem. In [4], Ciarlet and Zou first studied both nodal finite element methods and edge finite element methods for Maxwell equations. Chen et al. in [3] also discussed a fully discrete finite element method for Maxwell equations with discontinuous coefficients by introducing Lagrangian multipliers. In

<sup>\*</sup> Received April 8, 2003; final revised September 8, 2003.

<sup>&</sup>lt;sup>1)</sup> Supported by the Special Fund for State Major Basic Research Project of China (G19990328), the Scientific Research Fund of the State Administration of Radio, Film & TV of China and the Com<sup>2</sup>MaC-KOSEF, POSTECH Research Fund 2003-2004 of Republic of Korea.

this paper, we will study eddy current equations with discontinuous coefficients by using the above methods and give their error estimates in the meanwhile.

This paper is organized as follows. In section 2, the eddy current model in eddy current region is first presented. Second, we give its  $\mathbf{A} - \phi$  variational form based on an optimal control formulation of the interface problem and study the feather of its solutions in section 3. Finally, the fully-discrete coupled and decoupled  $\mathbf{A} - \phi$  schemes are proposed and their optimal error estimates are obtained in section 4 and 5 respectively.

### 2. Eddy Current Problem

For simplification, this paper is only concerned with the following eddy current equations in the eddy current region (high conductivity) neglecting the effect of outside source current:

$$\operatorname{curl} \mathbf{H} = \sigma \mathbf{E}, \quad \text{in } \Omega \times (0, T),$$

$$(2.1)$$

$$\operatorname{curl} \mathbf{E} = -\frac{\partial(\mu \mathbf{H})}{\partial t}, \quad \text{in } \Omega \times (0, T).$$
 (2.2)

Here  $\Omega \subset \mathbb{R}^3$  is a simply-connected Lipschitz polyhedral domain with connected boundary which is occupied by the dielectric material. We assume that the magnetic permeability parameter  $\mu$ and the conductivity  $\sigma$  of the medium are discontinuous across an interface  $\Gamma \subset \Omega$  respectively, where  $\Gamma$  is the boundary of a simply-connected Lipschitz polyhedral domain  $\Omega_1$  with  $\overline{\Omega}_1 \subset \Omega$ and  $\Omega_2 = \Omega \setminus \overline{\Omega}_1$ .  $\Omega_2$  should be multiply-connected as  $\Omega_1$  is simply-connected and lies strictly in  $\Omega$ . Without loss of generality we consider only the case with  $\mu$  and  $\sigma$  being two piecewise constant function in the domain  $\Omega$ , namely,

$$\mu = \begin{cases} \mu_1 & \text{in } \Omega_1, \\ \mu_2 & \text{in } \Omega_2, \end{cases} \qquad \sigma = \begin{cases} \sigma_1 & \text{in } \Omega_1, \\ \sigma_2 & \text{in } \Omega_2, \end{cases}$$

and  $\mu_i$ ,  $\sigma_i$  (i = 1, 2) are positive constants. It is known that magnetic and electric fields must satisfy the following jump conditions across the interface  $\Gamma$ :

$$[\mathbf{H} \times \mathbf{n}] = \mathbf{0},\tag{2.3}$$

$$[\mathbf{E} \times \mathbf{n}] = \mathbf{0},\tag{2.4}$$

where **n** is the unit outward normal to  $\partial \Omega_1$ . Throughout the paper, the jump of any function A across the interface  $\Gamma$  is defined as

$$[A] := A_2|_{\Gamma} - A_1|_{\Gamma}$$

with  $A_i = A|_{\Omega_i}$ , i = 1, 2. From (2.1) and (2.4) we can see that,

$$\left[\frac{1}{\sigma}\operatorname{\mathbf{curl}}\mathbf{H}\times\mathbf{n}\right] = \mathbf{0}, \quad \text{on } \Gamma\times(0,T).$$
(2.5)

We supplement the equation (2.1)-(2.2) with the boundary condition

$$\mathbf{H} \times \mathbf{n} = \mathbf{h}(\mathbf{x}, t), \tag{2.6}$$

and the initial condition

$$\mathbf{H}(\mathbf{x},0) = \mathbf{H}_0(\mathbf{x}), \quad \text{in } \Omega, \tag{2.7}$$

with

$$\operatorname{div}(\mu \mathbf{H}_0) = 0.$$

By taking divergence to both hand sides of (2.2), we easily see that,

$$\operatorname{div}(\mu \mathbf{H}) = 0, \quad \text{in } \Omega \times (0, T).$$

From the equation (2.1), we can suggest the introduction of a vector **A**, defined by

$$\sigma \mathbf{E} = \mathbf{curl} \, \mathbf{A}, \quad \text{in } \Omega_1 \cup \Omega_2 \times (0, T). \tag{2.8}$$

We then have

$$\mathbf{H} = \mathbf{A} + \nabla \phi, \quad \text{in } \Omega_1 \cup \Omega_2 \times (0, T), \tag{2.9}$$

where  $\phi(t)$  is an arbitrary scalar function. We assume that  $\nabla \phi_1 \times \mathbf{n} = \nabla \phi_2 \times \mathbf{n} = \mathbf{0}$  on  $\Gamma$  with  $\phi_i = \phi|_{\Omega_i}, i = 1, 2$ . So the system (2.1)-(2.2) is taken as the following  $\mathbf{A} - \phi$  form:

$$\mu \frac{\partial (\mathbf{A} + \nabla \phi)}{\partial t} + \mathbf{curl} \left( \frac{1}{\sigma} \mathbf{curl} \, \mathbf{A} \right) = 0, \quad \text{in } \Omega_1 \cup \Omega_2 \times (0, T), \tag{2.10}$$

$$\operatorname{div}(\mu(\mathbf{A} + \nabla\phi)) = 0, \quad \text{in } \Omega_1 \cup \Omega_2 \times (0, T)$$
(2.11)

with the following interface and boundary conditions

$$[\mathbf{A} \times \mathbf{n}] = \mathbf{0}, \quad \nabla \phi_1 \times \mathbf{n} = \nabla \phi_2 \times \mathbf{n} = \mathbf{0}, \quad \text{on } \Gamma \times (0, T), \tag{2.12}$$

$$\left[\frac{1}{\sigma}\operatorname{\mathbf{curl}}\mathbf{A}\times\mathbf{n}\right] = \mathbf{0}, \quad \text{ on } \Gamma\times(0,T), \tag{2.13}$$

$$\mathbf{A} \times \mathbf{n} = \mathbf{0}, \quad \nabla \phi \times \mathbf{n} = \mathbf{h}(\mathbf{x}, t), \quad \text{on } \partial \Omega \times (0, T), \tag{2.14}$$

and the initial conditions

$$\phi(\mathbf{x},0) = \phi_0(\mathbf{x}) \quad \text{and} \quad \mathbf{A}(\mathbf{x},0) = \mathbf{H}_0(\mathbf{x}) - \nabla\phi_0, \quad \text{in } \Omega.$$
(2.15)

where  $\phi_0$  is a given function with  $\nabla \phi_0 \times \mathbf{n} = \mathbf{h}(\mathbf{x}, 0)$  on  $\partial \Omega$  and  $\nabla \phi_0 \times \mathbf{n} = \mathbf{0}$  on  $\Gamma$ . For the sake of simplicity, we assume that  $\mathbf{h}(\mathbf{x}, t) = \mathbf{0}$  in the following theoretical analysis.

#### 3. The Variational Formulation

For solving the system (2.10)-(2.15), the finite element method with a matching finite element mesh on the interface  $\Gamma$  need impose a serious restriction: both must match with each other on  $\Gamma$ . We are now going to relax this restriction and consider a nonmatching mesh on the interface  $\Gamma$  that allows the triangulations in  $\Omega_1$  and  $\Omega_2$  to be generated independently. This advantage, however, brings some difficulty to the convergence analysis since the resulting finite element spaces will be nonconforming for **A**. So we will deal with the constraint  $[\mathbf{A} \times \mathbf{n}] = [\frac{1}{2} \operatorname{\mathbf{curl}} \mathbf{A} \times \mathbf{n}] = \mathbf{0}$  on  $\Gamma$  by a Lagrangian multiplier approach.

First, we introduce some notations that will be used throughout the paper.

Let  $L^p(0,T;X)$  denote the set of all strongly measurable functions  $u(t,\cdot)$  from [0,T] into the Hilbert space X such that

$$\int_0^T \|u(t)\|_X^p \, dt < \infty, \qquad 1 \le p < \infty.$$

We say  $u \in H^s(0,T;X)$  (s is a positive integer number) if and only if  $u, \frac{\partial u}{\partial t}, \cdots, \frac{\partial^s u}{\partial^s t}$  are in  $L^2(0,T;X)$ . Let

$$H(\operatorname{\mathbf{curl}};\Omega) = \{\mathbf{v} \in L^2(\Omega)^3; \operatorname{\mathbf{curl}} \mathbf{v} \in L^2(\Omega)^3\},\$$
$$H^{\alpha}(\operatorname{\mathbf{curl}};\Omega) = \{\mathbf{v} \in H^{\alpha}(\Omega)^3; \operatorname{\mathbf{curl}} \mathbf{v} \in H^{\alpha}(\Omega)^3\}, \ (\alpha > 0),\$$
$$H_0(\operatorname{\mathbf{curl}};\Omega) = \{\mathbf{v} \in H(\operatorname{\mathbf{curl}};\Omega); \mathbf{v} \times \mathbf{n} = \mathbf{0} \text{ on } \partial\Omega\}$$

with the norms

$$\begin{split} \|\mathbf{v}\|_{\mathbf{curl},\Omega} &= \left(\|\mathbf{v}\|_{0,\Omega}^2 + \|\mathbf{curl}\,\mathbf{v}\|_{0,\Omega}^2\right)^{1/2},\\ \|\mathbf{v}\|_{\alpha,\mathbf{curl},\Omega} &= \left(\|\mathbf{v}\|_{\alpha,\Omega}^2 + \|\mathbf{curl}\,\mathbf{v}\|_{\alpha,\Omega}^2\right)^{1/2}. \end{split}$$

Here and in what follows,  $\|\cdot\|_{0,\Omega}$  denotes the  $L^2(\Omega)^3$ -norm (or the  $L^2(\Omega)$ -norm for scalar functions) and for s > 0,  $\|\cdot\|_{s,\Omega}$  denotes the norm of the Sobolev space  $H^s(\Omega)^3$  (or  $H^s(\Omega)$  for scalar functions). Similar definitions are adopted for  $\Omega_1$  and  $\Omega_2$ . The constant C will always represent a generic constant independent of the time step and the mesh size.

For the convenience of presentation, let us introduce the following spaces:

$$X_1 = H(\mathbf{curl}; \Omega_1), \quad X_2 = \{\mathbf{v} \in H(\mathbf{curl}; \Omega_2); \mathbf{v} \times \mathbf{n} = \mathbf{0} \text{ on } \partial \Omega \},\$$

$$Y_1 = \{ \psi \in H^1(\Omega_1); \, \nabla \psi \times \mathbf{n} = 0 \text{ on } \partial \Omega_1 \}, \ Y_2 = \{ \psi \in H^1(\Omega_2); \, \nabla \psi \times \mathbf{n} = 0 \text{ on } \partial \Omega_2 \}.$$

Set  $X = X_1 \times X_2$  and  $Y = Y_1 \times Y_2$ .

Second, to establish an appropriate variational formulation for the system (2.10)-(2.15), we need use a few important mathematical analysis tools borrowed from [3].

Let

$$T(\Gamma) = \{ \mathbf{s} \in H^{-1/2}(\Gamma)^3; \exists \mathbf{v} \in H_0(\mathbf{curl}; \Omega) \text{ such that } \mathbf{v} \times \mathbf{n} = \mathbf{s} \text{ on } \Gamma \}.$$

It is not difficult to see that  $T(\Gamma)$  is a Banach space under the norm:

$$\|\mathbf{s}\|_{T(\Gamma)} = \inf\{\|\mathbf{v}\|_{\mathbf{curl},\Omega}; \mathbf{v} \in H_0(\mathbf{curl};\Omega) \text{ and } \mathbf{v} \times \mathbf{n} = \mathbf{s} \text{ on } \Gamma\}.$$

For any  $\mathbf{s} \in T(\Gamma)$ , we define

$$\ll \mathbf{s}, \mathbf{w} \gg_{1,\Gamma} = \int_{\Omega_1} \mathbf{v} \cdot \mathbf{curl} \, \mathbf{w} \, dx - \int_{\Omega_1} \mathbf{curl} \, \mathbf{v} \cdot \mathbf{w} \, dx, \quad \forall \, \mathbf{w} \in X_1,$$
(3.1)

$$\ll \mathbf{s}, \mathbf{w} \gg_{2,\Gamma} = \int_{\Omega_2} \operatorname{\mathbf{curl}} \mathbf{v} \cdot \mathbf{w} \, dx - \int_{\Omega_2} \mathbf{v} \cdot \operatorname{\mathbf{curl}} \mathbf{w} \, dx, \quad \forall \, \mathbf{w} \in X_2,$$
(3.2)

where  $\mathbf{v} \in H_0(\mathbf{curl}; \Omega)$  such that  $\mathbf{v} \times \mathbf{n} = \mathbf{s}$  on  $\Gamma$ . We know that (3.1)-(3.2) are independent of the choice of  $\mathbf{v} \in H(\mathbf{curl}; \Omega)$  such that  $\mathbf{v} \times \mathbf{n} = \mathbf{s}$  on  $\Gamma$ .

Then, we have Lemma 3.1-3.2 which are proved in [3]. Lemma 3.1. For any  $\mathbf{s} \in T(\Gamma)$ , we have the equality

$$\|\mathbf{s}\|_{T(\Gamma)} = \sup_{\mathbf{w}\in X_1\times X_2} \frac{\ll \mathbf{s}, \mathbf{w} \gg_{1,\Gamma} - \ll \mathbf{s}, \mathbf{w} \gg_{2,\Gamma}}{\|\mathbf{w}\|_{X_1\times X_2}}$$

A direct consequence of this lemma is that  $T(\Gamma)$  is indeed a Hilbert space.

In practice, Lemma 3.1 is rather inconvenient as it uses information from both  $\Omega_1$  and  $\Omega_2$  to define the norm on  $T(\Gamma)$ . To overcome the inconvenience, we note that

$$\|\mathbf{s}\|_{1,\Gamma} = \sup_{\mathbf{w}\in X_1} \frac{\ll \mathbf{s}, \mathbf{w} \gg_{1,\Gamma}}{\|\mathbf{w}\|_{X_1}}, \quad \|\mathbf{s}\|_{2,\Gamma} = \sup_{\mathbf{w}\in X_2} \frac{\ll \mathbf{s}, \mathbf{w} \gg_{1,\Gamma}}{\|\mathbf{w}\|_{X_2}}$$
(3.3)

are also norms of  $T(\Gamma)$ . So we have

**Lemma 3.2.** The norms  $\|\cdot\|_{1,\Gamma}$  and  $\|\cdot\|_{2,\Gamma}$  are equivalent to  $\|\cdot\|_{T(\Gamma)}$ .

Taking any  $\overline{\mathbf{A}} \in X$ ,  $\nabla \overline{\phi}$  (any  $\overline{\phi} \in Y$ ) in (2.10) as test functions respectively and applying the standard technique of integration by parts lead immediately to the following weak formulations of (2.10)-(2.15):

Problem (I). Find  $(\mathbf{A}, \phi, \mathbf{p}) \in H^1(0, T; X) \times H^1(0, T; Y) \times L^2(0, T; T(\Gamma))$  such that

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\partial (\mathbf{A} + \nabla \phi)}{\partial t}, \overline{\mathbf{A}})_{\Omega_{i}} + (\frac{1}{\sigma} \operatorname{\mathbf{curl}} \mathbf{A}, \operatorname{\mathbf{curl}} \overline{\mathbf{A}})_{\Omega_{i}} \right\} + \ll \mathbf{p}, \overline{\mathbf{A}} \gg_{2,\Gamma} - \ll \mathbf{p}, \overline{\mathbf{A}} \gg_{1,\Gamma} = 0, \quad \forall \overline{\mathbf{A}} \in X,$$

$$(3.4)$$

$$\sum_{i=1}^{2} \left(\mu \frac{\partial (\mathbf{A} + \nabla \phi)}{\partial t}, \nabla \overline{\phi}\right)_{\Omega_{i}} = 0, \quad \forall \overline{\phi} \in Y,$$
(3.5)

$$\ll \mathbf{A}, \, \overline{\mathbf{p}} \gg_{2,\Gamma} - \ll \mathbf{A}, \, \overline{\mathbf{p}} \gg_{1,\Gamma} = 0, \quad \forall \, \overline{\mathbf{p}} \in T(\Gamma).$$
 (3.6)

The system (3.4)-(3.6) is consistent with the finite element discretization on a nonmatching grid on the interface  $\Gamma$  and can be derived based on an optimal control formulation of the interface problem (2.10)-(2.15).

In order to analyze the feature of the solution of Problem (I), we first study the existence and uniqueness of the solution to the following problem:

Problem (II). Find  $(\mathbf{H}, \mathbf{p}) \in H^1(0, T; X) \times L^2(0, T; T(\Gamma))$  such that

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\partial \mathbf{H}}{\partial t}, \overline{\mathbf{H}})_{\Omega_{i}} + (\frac{1}{\sigma} \operatorname{\mathbf{curl}} \mathbf{H}, \operatorname{\mathbf{curl}} \overline{\mathbf{H}})_{\Omega_{i}} \right\} + \ll \mathbf{p}, \overline{\mathbf{H}} \gg_{2,\Gamma} - \ll \mathbf{p}, \overline{\mathbf{H}} \gg_{1,\Gamma} = 0, \quad \forall \overline{\mathbf{H}} \in X,$$

$$(3.7)$$

$$\ll \mathbf{H}, \, \overline{\mathbf{p}} \gg_{2,\Gamma} - \ll \mathbf{H}, \, \overline{\mathbf{p}} \gg_{1,\Gamma} = 0, \quad \forall \, \overline{\mathbf{p}} \in T(\Gamma).$$
 (3.8)

Furthermore, we only need to discuss the following stationary variational problem, whose result can be extended to Problem (II) by the standard analytic method.

Problem (III). Given  $\mathbf{f} \in L^2(\Omega)^3$ , find  $(\mathbf{Q}, \mathbf{p}) \in X \times T(\Gamma)$  such that

$$\sum_{i=1}^{2} \left\{ (\alpha_{i} \operatorname{\mathbf{curl}} \mathbf{Q}, \operatorname{\mathbf{curl}} \overline{\mathbf{Q}})_{\Omega_{i}} + (\beta_{i} \mathbf{Q}, \overline{\mathbf{Q}})_{\Omega_{i}} \right\} + \ll \mathbf{p}, \overline{\mathbf{Q}} \gg_{2,\Gamma} - \ll \mathbf{p}, \overline{\mathbf{Q}} \gg_{1,\Gamma} = \sum_{i=1}^{2} (\mathbf{f}, \overline{\mathbf{Q}})_{\Omega_{i}}, \quad \forall \overline{\mathbf{Q}} \in X,$$

$$(3.9)$$

$$\ll \mathbf{Q}, \, \overline{\mathbf{p}} \gg_{2,\Gamma} - \ll \mathbf{Q}, \, \overline{\mathbf{p}} \gg_{1,\Gamma} = 0, \quad \forall \, \overline{\mathbf{p}} \in T(\Gamma).$$
 (3.10)

Here  $\alpha_i$  and  $\beta_i$  are piecewise positive constants in  $\Omega_i$  for i = 1, 2. **Theorem 3.1.** There exists a unique solution  $(\mathbf{Q}, \mathbf{p}) \in X \times T(\Gamma)$  to Problem (III). *Proof.* First, define a bilinear form  $a: X \times X \to R$ :

$$a(\mathbf{u}, \mathbf{v}) = \sum_{i=1}^{2} \left\{ (\alpha_i \operatorname{\mathbf{curl}} \mathbf{u}, \operatorname{\mathbf{curl}} \mathbf{v})_{\Omega_i} + (\beta_i \mathbf{u}, \mathbf{v})_{\Omega_i} \right\}, \quad \mathbf{u}, \mathbf{v} \in X.$$

It is obvious that

$$a(\mathbf{u}, \mathbf{u}) \ge a_0 \|\mathbf{u}\|_X^2 \tag{3.11}$$

for some constant  $a_0 > 0$ .

Next, we verify the inf-sup condition: there exists a constant C > 0 such that

$$\sup_{\mathbf{B}\in X} \frac{\ll \mathbf{s}, \mathbf{B}_2 \gg_{2,\Gamma} - \ll \mathbf{s}, \mathbf{B}_1 \gg_{1,\Gamma}}{\|\mathbf{B}\|_X} \ge C \|\mathbf{s}\|_{T(\Gamma)}, \quad \forall \, \mathbf{s} \in T(\Gamma),$$
(3.12)

where  $\mathbf{B} = \mathbf{B}_i$  in  $\Omega_i$  for i = 1, 2. Let  $\mathbf{B} \in H(\mathbf{curl}; \Omega_1)$  be the solution of the following problem:

$$(\operatorname{\mathbf{curl}} \mathbf{B}, \operatorname{\mathbf{curl}} \overline{\mathbf{B}})_{\Omega_1} + (\mathbf{B}, \overline{\mathbf{B}})_{\Omega_1} = \ll \mathbf{s}, \overline{\mathbf{B}} \gg_{1,\Gamma}, \quad \forall \, \overline{\mathbf{B}} \in H(\operatorname{\mathbf{curl}}; \, \Omega_1).$$
(3.13)

We define

$$\widetilde{\mathbf{B}} = \begin{cases} -\mathbf{B} & \text{in } \Omega_1, \\ \mathbf{0} & \text{in } \Omega_2. \end{cases}$$
(3.14)

It is obvious that  $\widetilde{\mathbf{B}} \in X$  and

$$\|\mathbf{\tilde{B}}\|_X = \|\mathbf{B}\|_{\mathbf{curl},\Omega_1}.$$
(3.15)

Thus, by (3.13), we are able to obtain

$$\ll \mathbf{s}, \widetilde{\mathbf{B}}_2 \gg_{2,\Gamma} - \ll \mathbf{s}, \widetilde{\mathbf{B}}_1 \gg_{1,\Gamma} = \ll \mathbf{s}, \mathbf{B} \gg_{1,\Gamma} = \|\mathbf{B}\|^2_{\mathbf{curl},\Omega_1}$$

which yields, together with Lemma 3.2,

$$\frac{\ll \mathbf{s}, \mathbf{B}_2 \gg_{2,\Gamma} - \ll \mathbf{s}, \mathbf{B}_1 \gg_{1,\Gamma}}{\|\mathbf{\widetilde{B}}\|_X} = \|\mathbf{B}\|_{\mathbf{curl},\Omega_1} = \|\mathbf{s}\|_{1,\Gamma} \ge C \|\mathbf{s}\|_{T(\Gamma)}.$$

From (3.11)-(3.12), we then have finished the proof of the theorem.

Thus, from the result of Theorem 3.1, we conclude that the solution  $(\mathbf{H}, \mathbf{p})$  of Problem (II) is existing and unique. Meanwhile, by using an appropriate application of the Green's formula, the Lagrange multiplier  $\mathbf{p}$  in Problem (II) satisfies the following relation:

$$\mathbf{p} = \frac{1}{\sigma} \operatorname{\mathbf{curl}} \mathbf{H} \times \mathbf{n}, \quad \text{ in } T(\Gamma) \times (0, T).$$

For the solution **H** of Problem (II) and a given  $\phi \in Y$ , let  $\mathbf{A} = \mathbf{H} - \nabla \phi$ . Especially, if we append the divergence-free property of **A**, **A** is unique. Taking  $\overline{\mathbf{H}} = \overline{\mathbf{A}}$  and  $\overline{\mathbf{H}} = \nabla \overline{\phi}$  for any  $\overline{\mathbf{A}} \in X$  and  $\overline{\phi} \in Y$  respectively in (3.7)-(3.8), we conclude that **A**,  $\phi$ , **p** satisfy Problem (I); that is:

**Theorem 3.2.** The solution  $(\mathbf{A}, \phi, \mathbf{p})$  of Problem (I) is existing, but only  $\mathbf{p}$  is unique. Furthermore, the Lagrange multiplier  $\mathbf{p}$  in Problem (I) satisfies:

$$\mathbf{p} = \frac{1}{\sigma} \operatorname{\mathbf{curl}} \mathbf{A} \times \mathbf{n}, \quad \text{in } T(\Gamma) \times (0, T).$$
(3.16)

886

## 4. A Fully-discrete Coupled $A - \phi$ Scheme with a Nonmatching Grid for Eddy Current Problem

In this section we propose a finite element method for solving Problem (I), which allows a nonmatching finite element grid on the interface  $\Gamma$ .

Let  $\mathcal{T}^{h_1}$  and  $\mathcal{T}^{h_2}$  be a shape regular triangulation of  $\Omega_1$  and  $\Omega_2$  respectively. They induce naturally two finite element triangulations  $\Gamma_{h_1}$  and  $\Gamma_{h_2}$  on the interface  $\Gamma$ . Let  $\Gamma_{h_0}$  be an another shape regular triangulation over  $\Gamma$ . Note that  $\Gamma_{h_i}$ , i = 0, 1, 2, are allowed to be different from each other. However, we make the following reasonable assumption:

(H1) Each triangle in  $\Gamma_{h_1}$  and  $\Gamma_{h_2}$  must be contained in some triangles of  $\Gamma_{h_0}$ .

We introduce the Nédélec  $H(\operatorname{curl}, \Omega_i)$ -conforming edge element space defined by

$$X_{h_i} = \{ \mathbf{v}_h \in X_i; \, \mathbf{v}_h = \mathbf{a}_K + \mathbf{b}_K \times \mathbf{x} \text{ on } K, \, \forall K \in \mathcal{T}^{h_i} \}, \quad i = 1, 2$$

where  $\mathbf{a}_K$  and  $\mathbf{b}_K$  are two constant vectors. It is known that any function  $\mathbf{v}_h \in X_{h_i}$  is uniquely determined by the degrees of freedom in the set  $M_E(\mathbf{v})$  of the moments on each element  $K \in \Gamma_{h_i}$ , which is given by

$$M_E(\mathbf{v}) = \bigg\{ \int_e \mathbf{v} \cdot \tau \, ds; \ e \text{ is an edge of } K \bigg\}.$$

Here  $\tau$  is the unit vector along the edge. For i = 1, 2 and any  $\mathbf{v} \in H^s(\Omega_i)^3$  with  $\operatorname{curl} \mathbf{v} \in L^p(\Omega_i)^3$ , where s > 1/2 and p > 2, we can define an interpolation  $\pi_h \mathbf{v} \in X_{h_i}$ , and  $\pi_h \mathbf{v}$  has the same degrees of freedom (defined by  $M_E(\mathbf{v})$ ) as  $\mathbf{v}$  on each K in  $\Gamma_{h_i}$ .

Let  $T_{h_0}(\Gamma)$  be the Nédélec  $T(\Gamma)$ -conforming edge element space defined by

$$T_{h_0}(\Gamma) = \{ \mathbf{s}_h \in T(\Gamma); \, \mathbf{s}_h = (\alpha_\tau + \beta_\tau \times \mathbf{n}) \times \mathbf{n}, \text{ on any } \tau \in \Gamma_{h_0}, \, \alpha_\tau, \, \beta_\tau \in \mathbb{R}^3 \}.$$

We also define the following finite element spaces

$$Y_{h_i} = \{ \varphi_h \in Y_i; \, \varphi_h |_K \in \mathcal{P}_1, \, \forall K \in \mathcal{T}^{h_i} \}, \quad i = 1, 2,$$

where  $\mathcal{P}_1$  is the space of linear polynomials. Let  $\Pi_h$  be the standard interpolating operator. Now set

$$X_h = X_{h_1} \times X_{h_2}, \quad Y_h = Y_{h_1} \times Y_{h_2}.$$

We will assume the inf-sup condition:

(H2) There exists a constant  $C^* > 0$  independent of  $h_0$ ,  $h_1$ ,  $h_2$  such that

$$\sup_{\mathbf{w}_{h_i} \in X_{h_i}} \frac{\ll \mathbf{s}_h, \mathbf{w}_{h_i} \gg_{i,\Gamma}}{\|\mathbf{w}_{h_i}\|_{\mathbf{curl},\Omega_i}} \ge C^* \|\mathbf{s}_h\|_{T(\Gamma)}, \quad \forall \mathbf{s}_h \in T_{h_0}(\Gamma), \quad i = 1 \text{ or } 2.$$
(4.1)

The assumption (H2) indicate that the mesh  $\Gamma_{h_0}$  should be coarse enough compared with the meshes  $\mathcal{T}^{h_1}$  or  $\mathcal{T}^{h_2}$  in order to stabilize the effect of the introduced Lagrangian multiplier. In subsection 4.4 of [3], by using a general compactness argument, (4.1) is verified to be valid at least when the mesh  $h_1$  or  $h_2$  is suitably small compared with  $h_0$ .

Let us divide the time interval (0,T) into M equally-spaced subintervals by using nodal points

$$0 = t_0 < t_1 < \dots < t_M = T$$

with  $t_n = n\tau$  and  $\tau = T/M$ , and denote n-th subinterval by  $I^n = (t_{n-1}, t_n]$ . For a continuous mapping  $u : [0,T] \to L^2(\Omega)$  or  $L^2(\Omega)^3$ , we define  $u^n(\cdot) = u(\cdot, t_n)$  for  $1 \le n \le M$ .

Now we are in a position to introduce the discrete version of Problem (I).

Problem (VI). For  $n = 0, 1, \dots, M-1$ , find  $(\mathbf{A}_h^{n+1}, \phi_h^{n+1}, \mathbf{p}_h^{n+1}) \in X_h \times Y_h \times T_{h_0}(\Gamma)$  such that

$$\mathbf{A}_h^0 = \pi_h \mathbf{A}_0, \quad \phi_h^0 = \Pi_h \phi_0 \tag{4.2}$$

and

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\mathbf{A}_{h}^{n+1} - \mathbf{A}_{h}^{n}}{\tau}, \overline{\mathbf{A}}_{h})_{\Omega_{i}} + (\frac{1}{\sigma} \operatorname{curl} \mathbf{A}_{h}^{n+1}, \operatorname{curl} \overline{\mathbf{A}}_{h})_{\Omega_{i}} + (\mu \nabla \frac{\phi_{h}^{n+1} - \phi_{h}^{n}}{\tau}, \overline{\mathbf{A}}_{h})_{\Omega_{i}} \right\} + \ll \mathbf{p}_{h}^{n+1}, \overline{\mathbf{A}}_{h} \gg_{2,\Gamma} - \ll \mathbf{p}_{h}^{n+1}, \overline{\mathbf{A}}_{h} \gg_{1,\Gamma} = 0, \quad \forall \overline{\mathbf{A}}_{h} \in X_{h},$$

$$(4.3)$$

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\mathbf{A}_{h}^{n+1} - \mathbf{A}_{h}^{n}}{\tau}, \nabla \overline{\phi}_{h})_{\Omega_{i}} + (\mu \nabla \frac{\phi_{h}^{n+1} - \phi_{h}^{n}}{\tau}, \nabla \overline{\phi}_{h})_{\Omega_{i}} \right\} = 0, \quad \forall \overline{\phi}_{h} \in Y_{h},$$
(4.4)

$$\ll \mathbf{A}_{h}^{n+1}, \, \overline{\mathbf{p}}_{h} \gg_{2,\Gamma} - \ll \mathbf{A}_{h}^{n+1}, \, \overline{\mathbf{p}}_{h} \gg_{1,\Gamma} = 0, \quad \forall \, \overline{\mathbf{p}}_{h} \in T_{h_{0}}(\Gamma).$$
(4.5)

We then have the following result:

**Theorem 4.1.** Under the assumptions (H1)-(H2) and  $\mathbf{A}_{h}^{n+1} \times \mathbf{n} = \mathbf{0}$  on  $\partial\Omega$ , the solution  $(\mathbf{A}_{h}^{n+1}, \phi_{h}^{n+1}, \mathbf{p}_{h}^{n+1}) \in X_{h} \times Y_{h} \times T_{h_{0}}(\Gamma)$  of Problem (VI) is existing.

*Proof.* We first introduce the following discrete version of Problem (III) and prove that it has a unique solution.

Problem (V). Find  $(\mathbf{Q}_h, \mathbf{p}_h) \in X_h \times T_{h_0}(\Gamma)$  such that

$$\sum_{i=1}^{2} \left\{ (\alpha_{i} \operatorname{\mathbf{curl}} \mathbf{Q}_{h}, \operatorname{\mathbf{curl}} \overline{\mathbf{Q}}_{h})_{\Omega_{i}} + (\beta_{i} \mathbf{Q}_{h}, \overline{\mathbf{Q}}_{h})_{\Omega_{i}} \right\}$$

$$+ \ll \mathbf{p}_{h}, \overline{\mathbf{Q}}_{h} \gg_{2,\Gamma} - \ll \mathbf{p}_{h}, \overline{\mathbf{Q}}_{h} \gg_{1,\Gamma} = \sum_{i=1}^{2} (\mathbf{f}, \overline{\mathbf{Q}}_{h})_{\Omega_{i}}, \ \forall \, \overline{\mathbf{Q}}_{h} \in X_{h}.$$

$$(4.6)$$

$$\ll \mathbf{Q}_h, \, \overline{\mathbf{p}}_h \gg_{2,\Gamma} - \ll \mathbf{Q}_h, \, \overline{\mathbf{p}}_h \gg_{1,\Gamma} = 0, \quad \forall \, \overline{\mathbf{p}}_h \in Q_{h_0}(\Gamma).$$
(4.7)

Furthermore, we only need verify the inf-sup condition: for any  $\mathbf{s}_h \in T_{h_0}(\Gamma)$ , there exists a constant C > 0 such that

$$\sup_{\mathbf{B}_{h}\in X_{h}} \frac{\ll \mathbf{s}_{h}, \mathbf{B}_{h_{2}} \gg_{2,\Gamma} - \ll \mathbf{s}_{h}, \mathbf{B}_{h_{1}} \gg_{1,\Gamma}}{\|\mathbf{B}_{h}\|_{1,\Omega}} \ge C \|\mathbf{s}_{h}\|_{T(\Gamma)},$$
(4.8)

where  $\mathbf{B}_h = \mathbf{B}_{h_i}$  in  $\Omega_i$  for i = 1, 2. Without loss of generality we assume that (4.1) is valid for i = 1. Thus there exists a  $\mathbf{w}_{h_1} \in X_{h_1}$  such that

$$\frac{\ll \mathbf{s}_h, \mathbf{w}_{h_1} \gg_{1,\Gamma}}{\|\mathbf{w}_{h_1}\|_{\mathbf{curl},\Omega_1}} \ge C^* \|\mathbf{s}_h\|_{T(\Gamma)}.$$
(4.9)

Let  $\mathbf{B}_{h_1} \in X_{h_1}$  be the solution of the following problem:

$$(\operatorname{\mathbf{curl}} \mathbf{B}_{h_1}, \operatorname{\mathbf{curl}} \overline{\mathbf{B}}_{h_1})_{\Omega_1} + (\mathbf{B}_{h_1}, \overline{\mathbf{B}}_{h_1})_{\Omega_1} = \ll \mathbf{s}_h, \overline{\mathbf{B}}_{h_1} \gg_{1,\Gamma}, \quad \forall \overline{\mathbf{B}}_{h_1} \in X_{h_1}.$$
(4.10)

Taking  $\overline{\mathbf{B}}_{h_1} = \mathbf{B}_{h_1}$  and  $\overline{\mathbf{B}}_{h_1} = \mathbf{w}_{h_1}$  as test functions respectively and using (4.9), we obtain

$$C^* \|\mathbf{s}_h\|_{T(\Gamma)} \le \|\mathbf{B}_{h_1}\|_{\mathbf{curl},\Omega_1} \le \|\mathbf{s}_h\|_{1,\Gamma} \le C \|\mathbf{s}_h\|_{T(\Gamma)}.$$
(4.11)

An H-based  $\mathbf{A} - \phi$  Method with a Nonmatching Grid for Eddy Current Problem with ...

We define

$$\widetilde{\mathbf{B}}_h = \begin{cases} -\mathbf{B}_{h_1} & \text{in } \Omega_1, \\ \mathbf{0} & \text{in } \Omega_2. \end{cases}$$

Then,

$$\|\mathbf{B}_{h}\|_{X} = \|\mathbf{B}_{h_{1}}\|_{\mathbf{curl},\Omega_{1}}.$$
(4.12)

Thus, by (4.11), we are able to obtain

$$\ll \mathbf{s}_h, \widetilde{\mathbf{B}}_{h_2} \gg_{2,\Gamma} - \ll \mathbf{s}_h, \widetilde{\mathbf{B}}_{h_1} \gg_{1,\Gamma} = \ll \mathbf{s}_h, \mathbf{B}_{h_1} \gg_{1,\Gamma} = \|\mathbf{B}_{h_1}\|_{\mathbf{curl},\Omega_1}^2$$

which yields,

$$\frac{\ll \mathbf{s}_h, \mathbf{\ddot{B}}_{h_2} \gg_{2,\Gamma} - \ll \mathbf{s}_h, \mathbf{\ddot{B}}_{h_1} \gg_{1,\Gamma}}{\|\mathbf{\widetilde{B}}_h\|_X} = \frac{\ll \mathbf{s}_h, \mathbf{B}_{h_1} \gg_{1,\Gamma}}{\|\mathbf{B}_{h_1}\|_{\mathbf{curl},\Omega_1}} = \|\mathbf{B}_{h_1}\|_{\mathbf{curl},\Omega_1} \ge C \|\mathbf{s}_h\|_{T(\Gamma)}.$$

Therefore, Problem (V) has a unique solution  $(\mathbf{Q}_h, \mathbf{p}_h)$ .

From the definition of the edge element space,  $\nabla \varphi_h \in X_h$  for any  $\varphi_h \in Y_h$ . Thus, for the solution  $\mathbf{Q}_h$  of Problem (V) and a given  $\phi_h^{n+1} \in Y_h$ ,  $\mathbf{A}_h^{n+1} = \mathbf{Q}_h - \nabla \phi_h^{n+1} \in X_h$ . Noting that div  $\mathbf{A}_h^{n+1} = 0$ , we have that  $\mathbf{A}_h^{n+1}$  is unique; while  $\phi_h^{n+1}$  depending on its boundary condition is not unique.  $\alpha = \frac{1}{\sigma}, \beta = \frac{\mu}{\tau}$  and  $\mathbf{f} = \frac{\mu}{\tau} (\mathbf{A}_h^n + \nabla \phi_h^n)$ . Taking  $\overline{\mathbf{Q}}_h = \overline{\mathbf{A}}_h$  and  $\overline{\mathbf{Q}}_h = \nabla \overline{\phi}_h$  for any  $\overline{\mathbf{A}}_h \in X_h$  and  $\overline{\phi}_h \in Y_h$  in (4.6)-(4.7) respectively and noting that

$$\mathbf{p}_h = \frac{1}{\sigma} \operatorname{\mathbf{curl}} \mathbf{Q}_h \times \mathbf{n} = \mathbf{p}_h^{n+1}, \quad \text{ in } T(\Gamma),$$

we see that  $(\mathbf{A}_h^{n+1}, \phi_h^{n+1}, \mathbf{p}_h^{n+1})$  satisfies Problem (VI). Thus, we finish the proof of Theorem 4.1.

Now we can state the following theorem on the relevant error estimate. **Theorem 4.2.** Under the condition of Theorem 4.1, let  $(\mathbf{A}^{n+1}, \phi^{n+1}, \mathbf{p}^{n+1})$  and  $(\mathbf{A}^{n+1}_h, \phi^{n+1}_h, \mathbf{p}^{n+1}_h)$  be the solutions of Problem (I) and Problem (VI) at time  $t = t_{n+1}$  respectively. Supposing that for some  $\alpha > 1/2$ ,

$$\mathbf{A} \in H^2(0,T; H^{\alpha}(\mathbf{curl};\Omega_1) \times H^{\alpha}(\mathbf{curl};\Omega_2)), \quad \phi \in H^2(0,T; Y \cap H^{1+\alpha}(\Omega_1) \times H^{1+\alpha}(\Omega_2)).$$

Then, we have the following error estimate:

$$\max_{0 \le n \le M-1} \left( \sum_{i=1}^{2} \| (\mathbf{A}_{h}^{n+1} + \nabla \phi_{h}^{n+1}) - (\mathbf{A}^{n+1} + \nabla \phi^{n+1}) \|_{0,\Omega_{i}}^{2} \right) \le C\tau^{2} + \sum_{i=1}^{2} C_{i} h_{i}^{2\alpha} + C_{0} h_{0}^{2\alpha}.$$

*Proof.* For the convenience of presentation, let  $b: X \times T(\Gamma) \to R$  the bilinear form as follows:

 $b(\mathbf{Q},\mathbf{p}) = \ll \mathbf{Q}, \mathbf{p} \gg_{2,\Gamma} - \ll \mathbf{Q}, \mathbf{p} \gg_{1,\Gamma}, \quad \forall (\mathbf{Q},\mathbf{p}) \in X \times T(\Gamma).$ 

Define the elliptic projection operator  $P_h: X \times T(\Gamma) \to X_h \times T_{h_0}(\Gamma)$ . For any  $(\mathbf{Q}, \mathbf{p}) \in X \times T(\Gamma)$ , we have

$$\sum_{i=1}^{2} \left\{ (P_h \mathbf{Q} - \mathbf{Q}, \overline{\mathbf{Q}}_h)_{\Omega_i} + (\frac{1}{\sigma} \operatorname{\mathbf{curl}} (P_h \mathbf{Q} - \mathbf{Q}), \operatorname{\mathbf{curl}} \overline{\mathbf{Q}}_h)_{\Omega_i} \right\} + b \left( \overline{\mathbf{Q}}_h, P_h \mathbf{p} - \mathbf{p} \right) = 0, \quad \forall \, \overline{\mathbf{Q}}_h \in X_h,$$

$$(4.13)$$

$$b(P_h \mathbf{Q} - \mathbf{Q}, \overline{\mathbf{p}}_h) = 0, \quad \forall \, \overline{\mathbf{p}}_h \in T_{h_0}(\Gamma).$$
 (4.14)

Let  $\overline{\mathbf{Q}}_h = \nabla \psi_h \in X_h$  for any  $\psi_h \in Y_h$ . Then, for any  $\mathbf{Q} \in X$ , we have by (3.16) and the definition of the spaces  $Y_h$  and  $T_{h_0}(\Gamma)$ ,

$$\sum_{i=1}^{2} (P_h \mathbf{Q} - \mathbf{Q}, \nabla \psi_h)_{\Omega_i} = 0.$$
(4.15)

We use the backward difference at  $t = t_{n+1}$  for (3.4)-(3.5) and obtain

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\mathbf{A}^{n+1} - \mathbf{A}^{n}}{\tau}, \overline{\mathbf{A}})_{\Omega_{i}} + (\frac{1}{\sigma} \mathbf{curl} \, \mathbf{A}^{n+1}, \mathbf{curl} \, \overline{\mathbf{A}})_{\Omega_{i}} + (\mu \nabla \frac{\phi^{n+1} - \phi^{n}}{\tau}, \overline{\mathbf{A}})_{\Omega_{i}} \right\} + b \, (\overline{\mathbf{A}}, \mathbf{p}^{n+1}) = -\sum_{i=1}^{2} \left\{ (\mu \mathbf{R}_{1}^{n+1}, \overline{\mathbf{A}})_{\Omega_{i}} + (\mu \mathbf{R}_{2}^{n+1}, \overline{\mathbf{A}})_{\Omega_{i}} \right\}, \quad \forall \overline{\mathbf{A}} \in X,$$

$$(4.16)$$

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\mathbf{A}^{n+1} - \mathbf{A}^{n}}{\tau}, \nabla \overline{\phi})_{\Omega_{i}} + (\mu \nabla \frac{\phi^{n+1} - \phi^{n}}{\tau}, \nabla \overline{\phi})_{\Omega_{i}} \right\}$$

$$= -\sum_{i=1}^{2} \left\{ (\mu \mathbf{R}_{1}^{n+1}, \nabla \overline{\phi})_{\Omega_{i}} + (\mu \mathbf{R}_{2}^{n+1}, \nabla \overline{\phi})_{\Omega_{i}} \right\}, \quad \forall \overline{\phi} \in Y,$$

$$(4.17)$$

$$b(\mathbf{A}^{n+1}, \overline{\mathbf{p}}) = 0, \quad \forall \, \overline{\mathbf{p}} \in T(\Gamma),$$
(4.18)

where

$$\mathbf{R}_{1}^{n+1} = \left(\frac{\partial \mathbf{A}}{\partial t}\right)_{t_{n+1}} - \frac{\mathbf{A}^{n+1} - \mathbf{A}^{n}}{\tau}, \quad \|\mathbf{R}_{1}^{n+1}\|_{0;\Omega_{i}} = \mathcal{O}(\tau), \tag{4.19}$$

$$\mathbf{R}_{2}^{n+1} = \left(\frac{\partial \nabla \phi}{\partial t}\right)_{t_{n+1}} - \nabla \frac{\phi^{n+1} - \phi^{n}}{\tau}, \quad \|\mathbf{R}_{2}^{n+1}\|_{0;\Omega_{i}} = \mathcal{O}(\tau).$$
(4.20)

Let  $\overline{\mathbf{A}} = \overline{\mathbf{A}}_h$ ,  $\overline{\phi} = \overline{\phi}_h$  and  $\overline{\mathbf{p}} = \overline{\mathbf{p}}_h$ . Subtracting (4.16)-(4.18) from (4.3)-(4.5), we have

$$\sum_{i=1}^{2} \left\{ \left( \mu \frac{(\mathbf{A}_{h}^{n+1} - \mathbf{A}^{n+1}) - (\mathbf{A}_{h}^{n} - \mathbf{A}^{n})}{\tau}, \overline{\mathbf{A}}_{h} \right)_{\Omega_{i}} + \left( \frac{1}{\sigma} \mathbf{curl} \left( \mathbf{A}_{h}^{n+1} - \mathbf{A}^{n+1} \right), \mathbf{curl} \overline{\mathbf{A}}_{h} \right)_{\Omega_{i}} + \left( \mu \nabla \frac{(\phi_{h}^{n+1} - \phi^{n+1}) - (\phi_{h}^{n} - \phi^{n})}{\tau}, \overline{\mathbf{A}}_{h} \right)_{\Omega_{i}} \right\} + b \left( \overline{\mathbf{A}}_{h}, \mathbf{p}_{h}^{n+1} - \mathbf{p}^{n+1} \right)$$

$$= \sum_{i=1}^{2} \left\{ \left( \mu \mathbf{R}_{1}^{n+1}, \overline{\mathbf{A}}_{h} \right)_{\Omega_{i}} + \left( \mu \mathbf{R}_{2}^{n+1}, \overline{\mathbf{A}}_{h} \right)_{\Omega_{i}} \right\}, \quad \forall \overline{\mathbf{A}}_{h} \in X_{h},$$

$$(4.21)$$

$$\sum_{i=1}^{2} \left\{ \left( \mu \frac{(\mathbf{A}_{h}^{n+1} - \mathbf{A}^{n+1}) - (\mathbf{A}_{h}^{n} - \mathbf{A}^{n})}{\tau}, \nabla \overline{\phi}_{h} \right)_{\Omega_{i}} + \left( \mu \nabla \frac{(\phi_{h}^{n+1} - \phi^{n+1}) - (\phi_{h}^{n} - \phi^{n})}{\tau}, \nabla \overline{\phi}_{h} \right)_{\Omega_{i}} \right\}$$

$$(4.22)$$

$$=\sum_{i=1}^{2}\left\{(\mu\mathbf{R}_{1}^{n+1},\nabla\overline{\phi}_{h})_{\Omega_{i}}+(\mu\mathbf{R}_{2}^{n+1},\nabla\overline{\phi}_{h})_{\Omega_{i}}\right\},\quad\forall\overline{\phi}_{h}\in Y_{h},$$
$$b\left(\mathbf{A}_{h}^{n+1}-\mathbf{A}^{n+1},\overline{\mathbf{p}}_{h}\right)=0,\quad\forall\overline{\mathbf{p}}_{h}\in T_{h_{0}}(\Gamma).$$
(4.23)

Set  $\Theta_h^{n+1} = \mathbf{A}_h^{n+1} - P_h \mathbf{A}^{n+1}$  and  $\eta_h^{n+1} = \phi_h^{n+1} - \Pi_h \phi^{n+1}$ . Taking  $\overline{\mathbf{A}}_h = \Theta_h^{n+1}$  and  $\overline{\phi}_h = \eta_h^{n+1}$  in (4.21)-(4.22), together with the definition of the projection operator  $P_h$ , we come to

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\Theta_{h}^{n+1} - \Theta_{h}^{n}}{\tau}, \Theta_{h}^{n+1})_{\Omega_{i}} + (\frac{1}{\sigma} \operatorname{curl} \Theta_{h}^{n+1}, \operatorname{curl} \Theta_{h}^{n+1})_{\Omega_{i}} + (\mu \nabla \frac{\eta_{h}^{n+1} - \eta_{h}^{n}}{\tau}, \Theta_{h}^{n+1})_{\Omega_{i}} \right\}$$

$$= \sum_{i=1}^{2} \left\{ (\mu \frac{(\mathbf{A}^{n+1} - P_{h} \mathbf{A}^{n+1}) - (\mathbf{A}^{n} - P_{h} \mathbf{A}^{n})}{\tau}, \Theta_{h}^{n+1})_{\Omega_{i}} + (\mathbf{A}^{n+1} - P_{h} \mathbf{A}^{n+1}, \Theta_{h}^{n+1})_{\Omega_{i}} + (\mu \nabla \frac{(\phi^{n+1} - \Pi_{h} \phi^{n+1}) - (\phi^{n} - \Pi_{h} \phi^{n})}{\tau}, \Theta_{h}^{n+1})_{\Omega_{i}} + (\mu \mathbf{R}_{1}^{n+1}, \Theta_{h}^{n+1})_{\Omega_{i}} + (\mu \mathbf{R}_{2}^{n+1}, \Theta_{h}^{n+1})_{\Omega_{i}} \right\},$$

$$(4.24)$$

$$\sum_{i=1}^{2} \left\{ \left( \mu \frac{\mathbf{\Theta}_{h}^{n+1} - \mathbf{\Theta}_{h}^{n}}{\tau}, \nabla \eta_{h}^{n+1} \right)_{\Omega_{i}} + \left( \mu \nabla \frac{\eta_{h}^{n+1} - \eta_{h}^{n}}{\tau}, \nabla \eta_{h}^{n+1} \right)_{\Omega_{i}} \right\}$$

$$= \sum_{i=1}^{2} \left\{ \left( \mu \frac{(\mathbf{A}^{n+1} - P_{h} \mathbf{A}^{n+1}) - (\mathbf{A}^{n} - P_{h} \mathbf{A}^{n})}{\tau}, \nabla \eta_{h}^{n+1} \right)_{\Omega_{i}}$$

$$+ \left( \mu \nabla \frac{(\phi^{n+1} - \Pi_{h} \phi^{n+1}) - (\phi^{n} - \Pi_{h} \phi^{n})}{\tau}, \nabla \eta_{h}^{n+1} \right)_{\Omega_{i}}$$

$$+ \left( \mu \mathbf{R}_{1}^{n+1}, \nabla \eta_{h}^{n+1} \right)_{\Omega_{i}} + \left( \mu \mathbf{R}_{2}^{n+1}, \nabla \eta_{h}^{n+1} \right)_{\Omega_{i}} \right\}.$$

$$(4.25)$$

Now adding up (4.24) and (4.25), multiplying both sides by  $\tau$  and using (4.15) and  $a(a-b) \ge a^2/2 - b^2/2$ , for any real numbers a and b, we have

$$\sum_{i=1}^{2} \left\{ \frac{1}{2} \| \sqrt{\mu} (\boldsymbol{\Theta}_{h}^{n+1} + \nabla \eta_{h}^{n+1}) \|_{0;\Omega_{i}}^{2} - \frac{1}{2} \| \sqrt{\mu} (\boldsymbol{\Theta}_{h}^{n} + \nabla \eta_{h}^{n}) \|_{0;\Omega_{i}}^{2} + \tau \| \frac{1}{\sqrt{\sigma}} \operatorname{curl} \boldsymbol{\Theta}_{h}^{n+1} \|_{0;\Omega_{i}}^{2} \right\} \\
\leq \sum_{i=1}^{2} \left\{ \tau (\mu \frac{(\mathbf{A}^{n+1} - P_{h} \mathbf{A}^{n+1}) - (\mathbf{A}^{n} - P_{h} \mathbf{A}^{n})}{\tau}, \boldsymbol{\Theta}_{h}^{n+1} + \nabla \eta_{h}^{n+1})_{\Omega_{i}} + \tau (\mathbf{A}^{n+1} - P_{h} \mathbf{A}^{n+1}, \boldsymbol{\Theta}_{h}^{n+1} + \nabla \eta_{h}^{n+1})_{\Omega_{i}} + \tau (\mu \nabla \frac{(\phi^{n+1} - \Pi_{h} \phi^{n+1}) - (\phi^{n} - \Pi_{h} \phi^{n})}{\tau}, \boldsymbol{\Theta}_{h}^{n+1} + \nabla \eta_{h}^{n+1})_{\Omega_{i}} + \tau (\mu \mathbf{R}_{1}^{n+1}, \boldsymbol{\Theta}_{h}^{n+1} + \nabla \eta_{h}^{n+1})_{\Omega_{i}} + \tau (\mu \mathbf{R}_{2}^{n+1}, \boldsymbol{\Theta}_{h}^{n+1} + \nabla \eta_{h}^{n+1})_{\Omega_{i}} \right\}.$$

$$(4.26)$$

On the other hand, for the projection operator  $P_h$ , we have

$$\sum_{i=1}^{2} \|\mathbf{A}^{n+1} - P_{h}\mathbf{A}^{n+1}\|_{\mathbf{curl},\Omega_{i}} + \|\mathbf{p}^{n+1} - P_{h}\mathbf{p}^{n+1}\|_{T(\Gamma)}$$

$$\leq C \inf_{\mathbf{Q}_{h}\in X_{h}} \|\mathbf{A}^{n+1} - \mathbf{Q}_{h}\|_{\mathbf{curl},\Omega_{i}} + C \inf_{\mathbf{q}_{h}\in T_{h_{0}}(\Gamma)} \|\mathbf{q}_{h} - \mathbf{p}^{n+1}\|_{T(\Gamma)}.$$
(4.27)

Now using the edge element interpolation estimate in [5], we get

$$\inf_{\mathbf{Q}_h \in X_{h_i}} \|\mathbf{A}^{n+1} - \mathbf{Q}_h\|_{\mathbf{curl},\Omega_i} \le Ch_i^{\alpha} \|\mathbf{A}^{n+1}\|_{\alpha,\mathbf{curl},\Omega_i}.$$
(4.28)

Next we introduce a triangulation  $\mathcal{T}^{h_0}$  in  $\Omega_1$  whose restriction on  $\Gamma$  coincides with  $\Gamma_{h_0}$  and let  $X_{h_0}$  be the Nédélec  $H(\operatorname{curl}, \Omega_1)$ -conforming edge element over the mesh  $\mathcal{T}^{h_0}$ . Then from the definition of  $T_{h_0}(\Gamma)$  we know that

$$T_{h_0}(\Gamma) = \{ \mathbf{v}_h \times \mathbf{n}; \ \mathbf{v}_h \in X_{h_0} \}.$$

Now using the fact that  $\mathbf{p}^{n+1} = \frac{1}{\sigma} \operatorname{curl} \mathbf{A}^{n+1} \times \mathbf{n} := \mathbf{T}^{n+1} \times \mathbf{n}$ , we can easily show that, by (2.10),  $\mathbf{T}^{n+1} \in H^{\alpha}(\operatorname{curl}, \Omega_1)$  and

$$\|\mathbf{T}^{n+1}\|_{\alpha,\mathbf{curl},\Omega_1} \le C(\|\mathbf{A}^{n+1}\|_{\alpha,\mathbf{curl},\Omega_1} + \|\phi_t^{n+1}\|_{1+\alpha,\Omega_1} + \|\mathbf{A}_t^{n+1}\|_{\alpha,\Omega_1}).$$

Thus we have by Lemma 3.2 and the standard edge element error estimate in [5] that

$$\inf_{\mathbf{q}_h \in T_{h_0}(\Gamma)} \|\mathbf{q}_h - \mathbf{p}^{n+1}\|_{T(\Gamma)} \leq C \inf_{\mathbf{v}_h \in X_{h_0}} \|\mathbf{v}_h - \mathbf{T}^{n+1}\|_{\mathbf{curl},\Omega_1} \leq C h_0^{\alpha} \|\mathbf{T}^{n+1}\|_{\alpha,\mathbf{curl},\Omega_1}.$$
(4.29)

From (4.28)-(4.29), we derive

$$\sum_{i=1}^{2} \|\mathbf{A}^{n+1} - P_h \mathbf{A}^{n+1}\|_{\mathbf{curl},\Omega_i} \le \sum_{i=1}^{2} C_i h_i^{\alpha} + C_0 h_0^{\alpha}.$$
(4.30)

Similar to the proof of Theorem 3.2 in [11], using the discrete Gronwall's inequality, the finite element interpolation element estimate and (4.30) to (4.26), we easily complete the proof of the theorem with the help of the triangle inequality. We omit the details.

## 5. A Fully-discrete Decoupled $A - \phi$ Scheme with a Nonmatching Grid for Eddy Current Problem

To avoid increasing the number of freedoms and equations by solving Problem (VI) directly, we present a new decoupled  $\mathbf{A} - \phi$  scheme in this part.

First, we need to extend **A** and  $\phi$  with some regularity from the time interval [0, T] to the interval  $[-\tau, T]$ . Let

$$A^{-1} = 0$$
 and  $\phi^{-1} = 0$ .

Then, the decoupled  $\mathbf{A} - \phi$  scheme is:

$$\mathbf{A}_{h}^{0} = \pi_{h} \mathbf{A}_{0}, \quad \phi_{h}^{0} = \Pi_{h} \phi_{0}, \quad \phi_{h}^{-1} = 0$$
(5.1)

and for  $n = 0, 1, \dots, M - 1$ , find  $(\mathbf{A}_h^{n+1}, \mathbf{p}_h^{n+1}) \in X_h \times T_{h_0}(\Gamma)$  such that

$$\sum_{i=1}^{2} \left\{ (\mu \frac{\mathbf{A}_{h}^{n+1} - \mathbf{A}_{h}^{n}}{\tau}, \overline{\mathbf{A}}_{h})_{\Omega_{i}} + (\frac{1}{\sigma} \operatorname{curl} \mathbf{A}_{h}^{n+1}, \operatorname{curl} \overline{\mathbf{A}}_{h})_{\Omega_{i}} \right\} + \ll \mathbf{p}_{h}^{n+1}, \overline{\mathbf{A}}_{h} \gg_{2,\Gamma}$$

$$- \ll \mathbf{p}_{h}^{n+1}, \overline{\mathbf{A}}_{h} \gg_{1,\Gamma} = -\sum_{i=1}^{2} (\mu \nabla \frac{\phi_{h}^{n} - \phi_{h}^{n-1}}{\tau}, \overline{\mathbf{A}}_{h})_{\Omega_{i}}, \quad \forall \overline{\mathbf{A}}_{h} \in X_{h},$$

$$\ll \mathbf{A}_{h}^{n+1}, \overline{\mathbf{p}}_{h} \gg_{2,\Gamma} - \ll \mathbf{A}_{h}^{n+1}, \overline{\mathbf{p}}_{h} \gg_{1,\Gamma} = 0, \quad \forall \overline{\mathbf{p}}_{h} \in T_{h_{0}}(\Gamma),$$

$$(5.2)$$

and find  $\phi_h^{n+1} \in Y_h$  such that

$$\sum_{i=1}^{2} (\mu \nabla \frac{\phi_h^{n+1} - \phi_h^n}{\tau}, \nabla \overline{\phi}_h)_{\Omega_i} = -\sum_{i=1}^{2} (\mu \frac{\mathbf{A}_h^{n+1} - \mathbf{A}_h^n}{\tau}, \nabla \overline{\phi}_h)_{\Omega_i}, \quad \forall \overline{\phi}_h \in Y_h.$$
(5.4)

From the discussion to Problem (V) in Theorem 4.1 we know that under the same assumptions of Theorem 4.1, the system (5.2)-(5.3) has a unique solution  $(\mathbf{A}_{h}^{n+1}, \mathbf{p}_{h}^{n+1})$  at each time step. Moreover, by introducing the elliptic projection operator ( referring to the proof of Theorem 4.2) and imitating the proof of Theorem 3.3 in [11], we have

**Theorem 5.1.** Under the same assumptions of Theorem 4.1, let  $(\mathbf{A}^{n+1}, \phi^{n+1}, \mathbf{p}^{n+1})$  and  $(\mathbf{A}_h^{n+1}, \phi_h^{n+1}, \mathbf{p}_h^{n+1})$  be the solutions of Problem (I) and the decoupled  $\mathbf{A} - \phi$  approximation (5.1)-(5.4) at time  $t = t_{n+1}$  respectively. Supposing that for some  $\alpha > 1/2$ ,

 $\mathbf{A} \in H^3(0,T; H^{\alpha}(\mathbf{curl};\Omega_1) \times H^{\alpha}(\mathbf{curl};\Omega_2)), \quad \phi \in H^2(0,T; Y \cap H^{1+\alpha}(\Omega_1) \times H^{1+\alpha}(\Omega_2)).$ 

Then, we have the following error estimate:

$$\max_{0 \le n \le M-1} \left( \sum_{i=1}^{2} \| (\mathbf{A}_{h}^{n+1} + \nabla \phi_{h}^{n+1}) - (\mathbf{A}^{n+1} + \nabla \phi^{n+1}) \|_{0,\Omega_{i}}^{2} \right) \le C\tau^{2} + \sum_{i=1}^{2} C_{i} h_{i}^{2\alpha} + C_{0} h_{0}^{2\alpha}.$$

Acknowledgement. The authors thank greatly two referees for many constructive comments which lead to a great improvement of the results and the presentation of the paper.

#### References

- R. Albanese and G. Rubinacci, Formulation of the eddy-current problem, *IEE Proceedings*, 137 (1990), 16-22.
- [2] H. Ammari, A. Buffa and J.-C. Nédélec, Ajustification of eddy currents model for the Maxwell equations, SIAM J. Appl. Math., 60 (2000), 1805-1823.
- [3] Z. Chen, Q. Du and J. Zou, Finite element methods with matching and nonmatching meshes for Maxwell equations with discontinuous coefficients, SIAM J. Numer. Anal., 37 (2000), 1542-1570.
- [4] P. Ciarlet, J. Zou, Finite element convergence for the Darwin model to Maxwell's equations, RAIRO Math. Modelling Numer. Anal., 31 (1997), 213-249.
- [5] P. Ciarlet, J. Zou, Fully discrete finite element approaches for time-dependent Maxwell's equations, Numerische Mathematik, 82 (1999), 193-219.
- [6] J.L. Guermond and L. Quartapelle, On the approximation of the unsteady Navier-Stokes equations by finite element projection method, *Numerische Mathematik*, 80 (1998), 207-238.
- [7] R. Hiptmair, Symmetric coupling for eddy current problems, SIAM J. Numer. Anal., 40 (2002), 41-65.
- [8] Q. Hu, G. Liang and J. Liu, Construction of a preconditioner for three-dimensional domain decomposition methods with Lagrangian Multipliers, J. Comput. Math., 19 (2001), 213-224.
- [9] Q. Hu and J. Zou, A non-overlapping domain decomposition method for Maxwell's equations in three dimensions, SIAM J. Numer. Anal., 42 (2003), 1682-1708.
- [10] T. Kang, C. Ma and G. Liang, H-based  $\mathbf{A} \phi$  approaches for eddy current problem, to appear.
- [11] T. Kang, C. Ma and G. Liang, H-based  $\mathbf{A} \phi$  approaches of approximating eddy current problem by way of solving different systems inside and outside the conductor, to appear in Appl. Math. Comp.
- [12] C. Ma, Studies on  $\mathbf{A} \phi$  methods for time-dependent electromagnetic fields (in Chinese), Ph.D. thesis, Chinese Academy of Scienses, 2003.

- [13] P. Monk, Analysis of a finite element method for Maxwell's equations, SIAM J. Numer. Anal., 29 (1992), 714-729.
- [14] J.-C. Nédélec, Mixed finite elements in  $\mathbb{R}^3$ . Numer. Math., **35** (1980), 315-341.
- [15] D. Yu, Natural boundary integral method and its applications, Kluwer Academic Publisher, 2002.